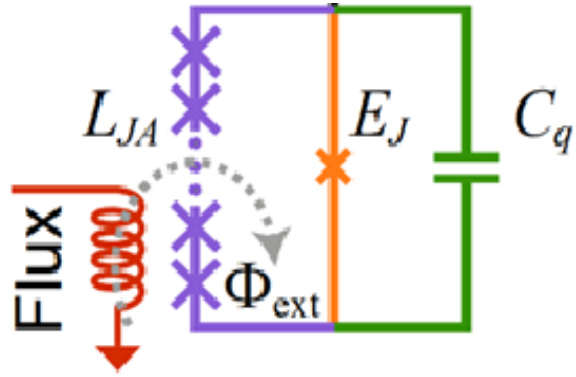


Reconsidering qubit control paradigms for high fidelity fluxonium gates

Chunyang Ding, Helin Zhang, Daniel Weiss, Sai Paivan Chitta, Yuwei Ma,
Jens Koch, David Schuster

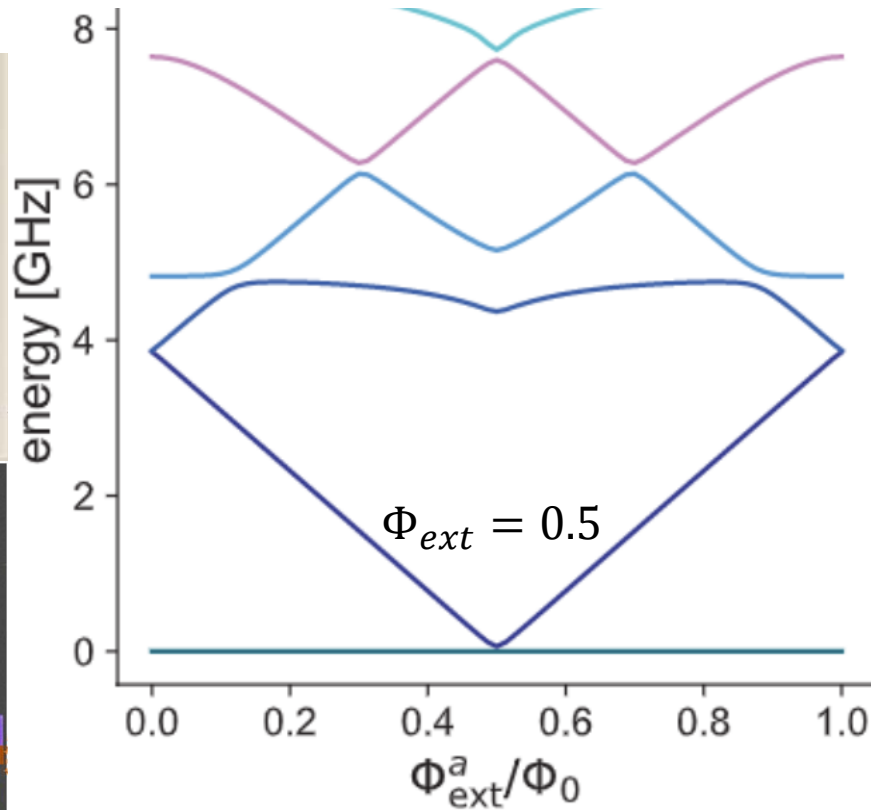


Fluxonium qubits

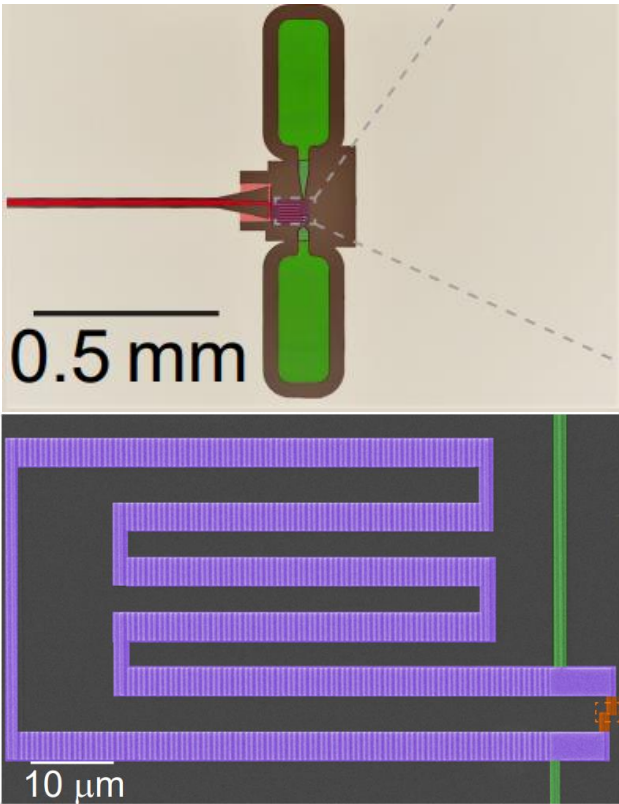
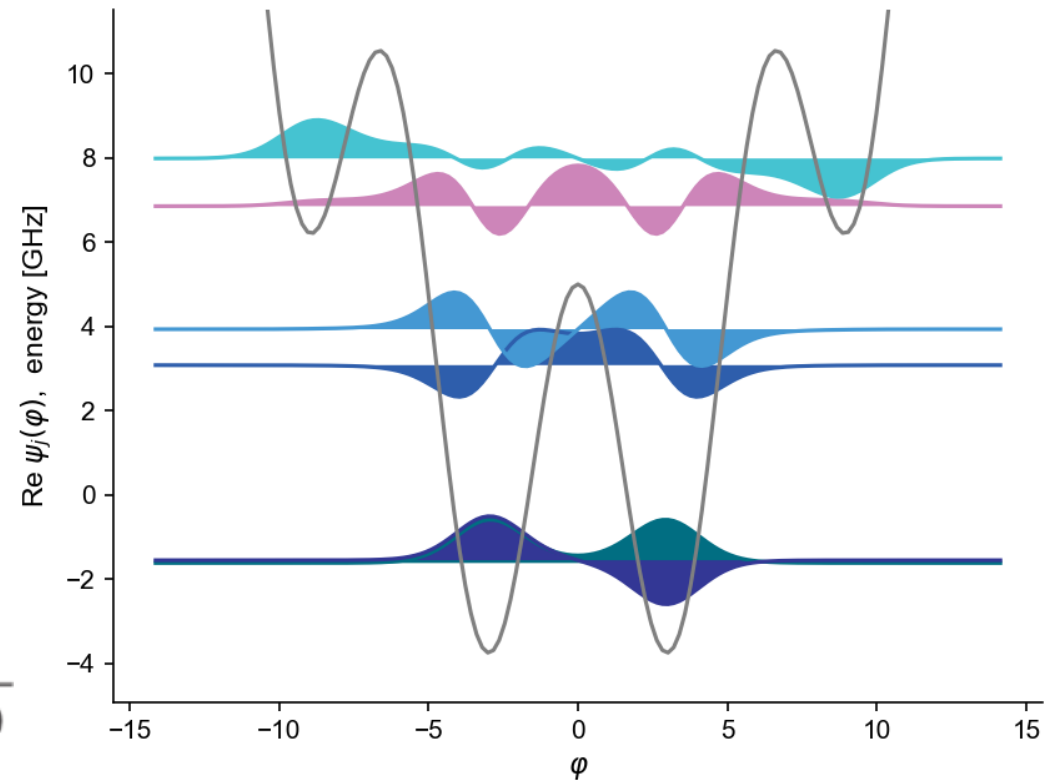


$$H_f = -4E_C \frac{d^2}{d\varphi^2} - E_J \cos(\varphi - 2\pi\Phi_{\text{ext}}/\Phi_0) + \frac{1}{2}E_L\varphi^2$$

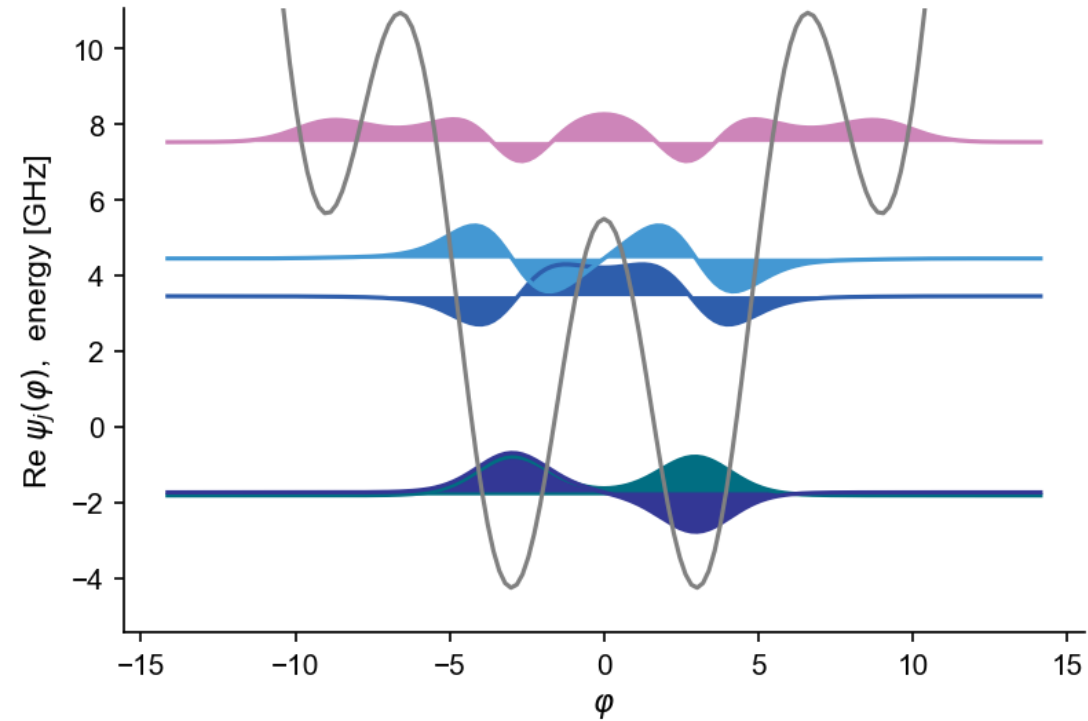
Energy spectrum (Φ_{ext})



$\Phi_{\text{ext}} = 0.5$



Why Fluxonium?



Important quantity for most SC qubit gates is:

Gate time / **Coherence time**

- Coherence (esp. T_1) scales inversely with qubit frequency
- Gate time scales like $\min(\omega_{01}, \alpha)$

Fluxonium decoherence should be slower, but gates can be faster

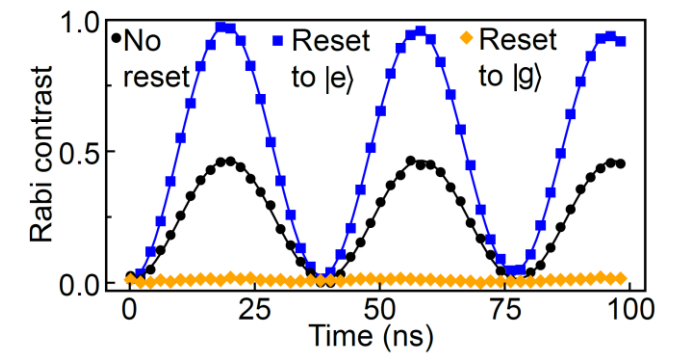
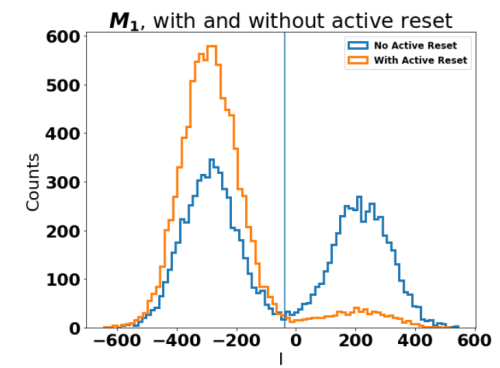
Low frequency is cheaper / easier

Less crowding of transitions

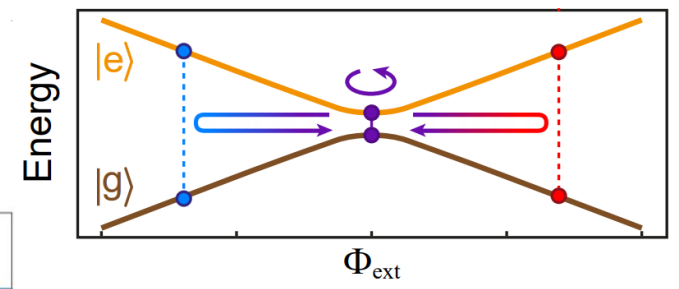
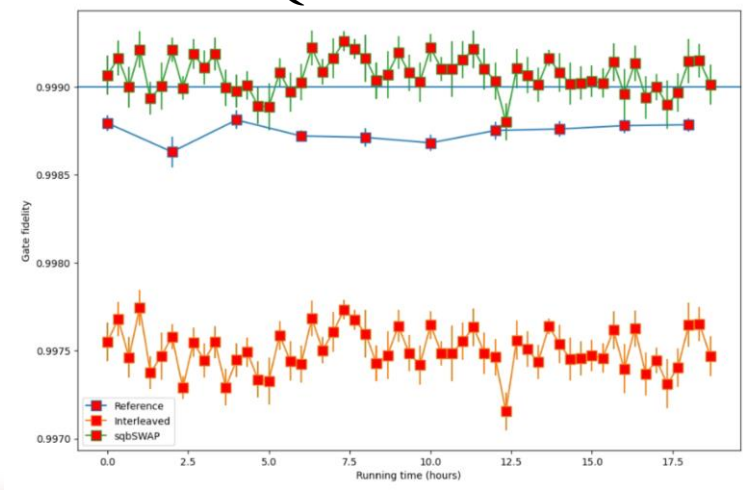
Rethinking paradigms



- Initialization via laser cooling
- Initialization via active reset
- Ultrafast single qubit gates
- High fidelity galvanically-coupled two qubit gates



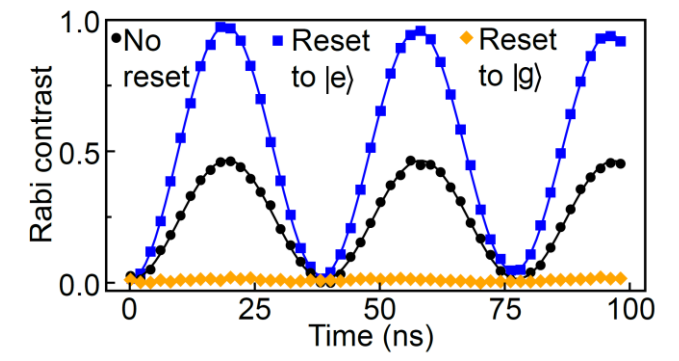
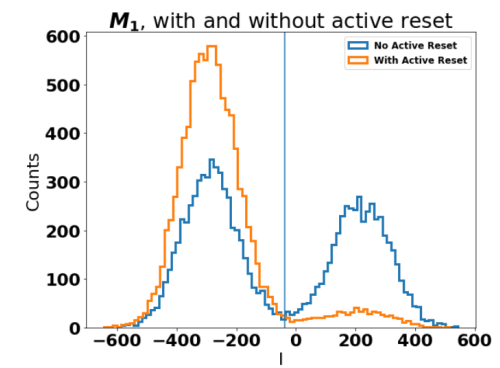
SQBSWAP XEB



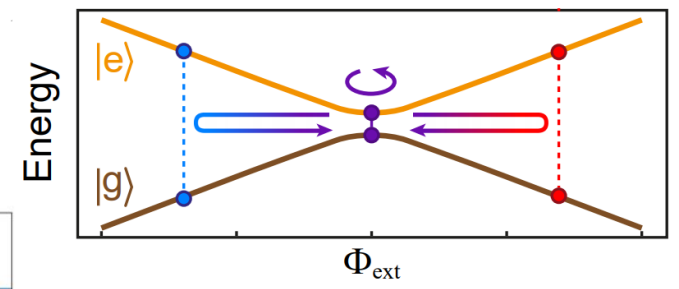
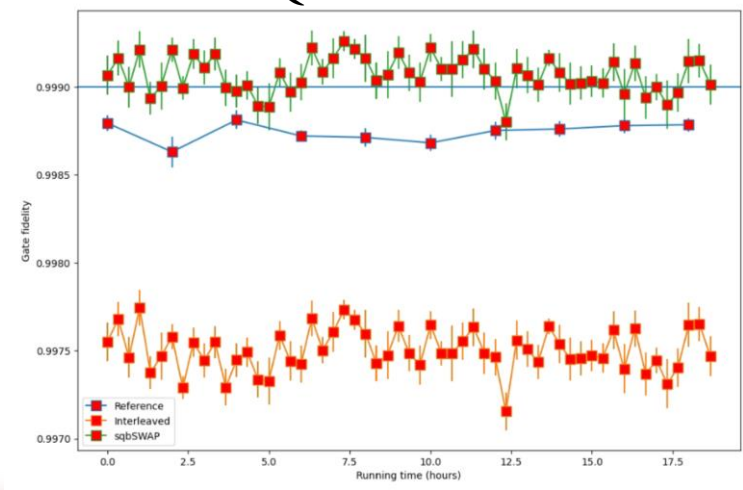
Rethinking paradigms



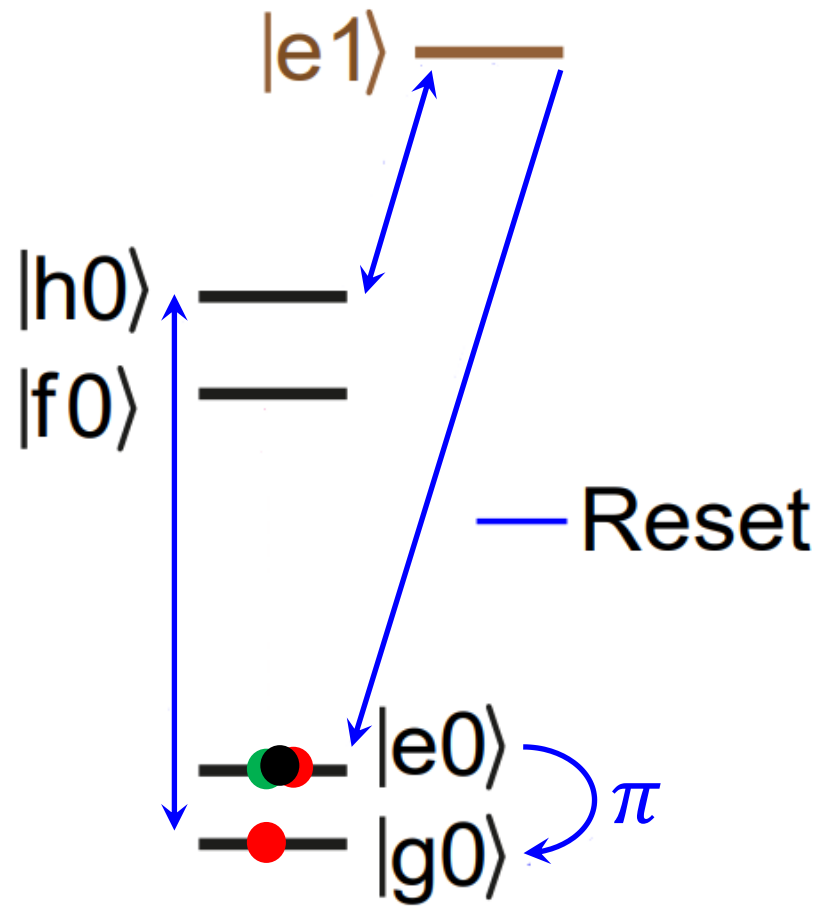
- Initialization via laser cooling
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SQBSWAP XEB

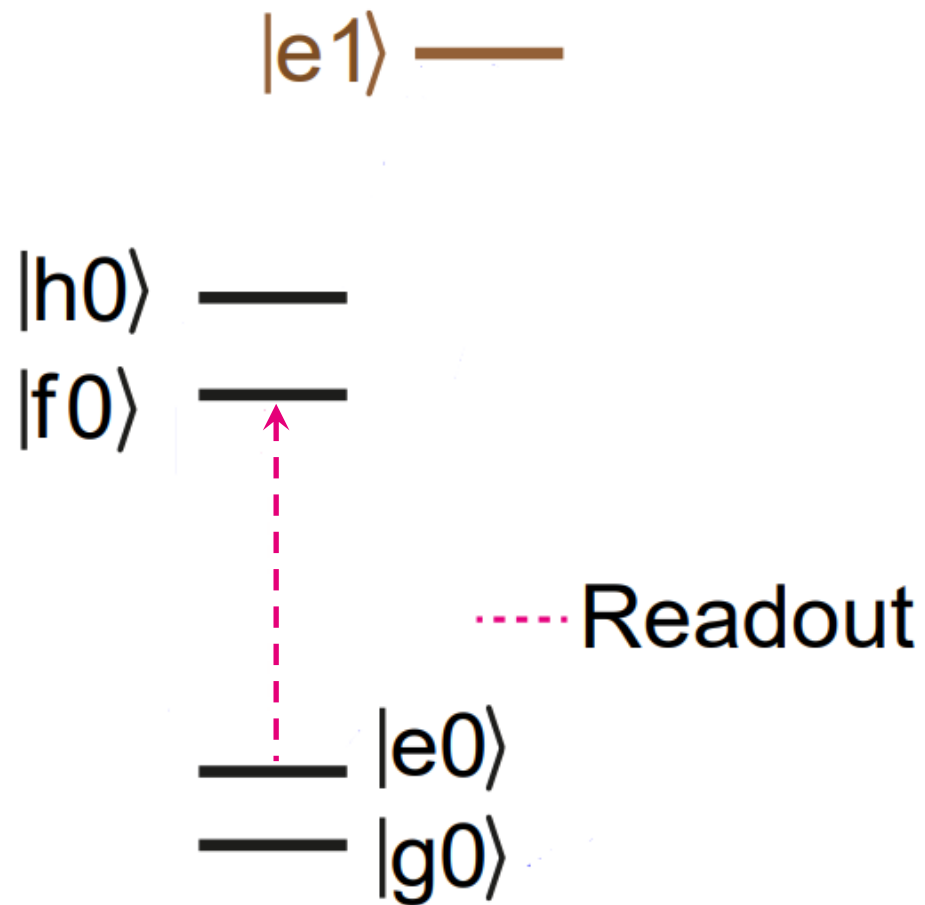


Qubit initialization and readout



- Reset fidelity limited by cavity & $|f\rangle$ state population

Qubit initialization and readout

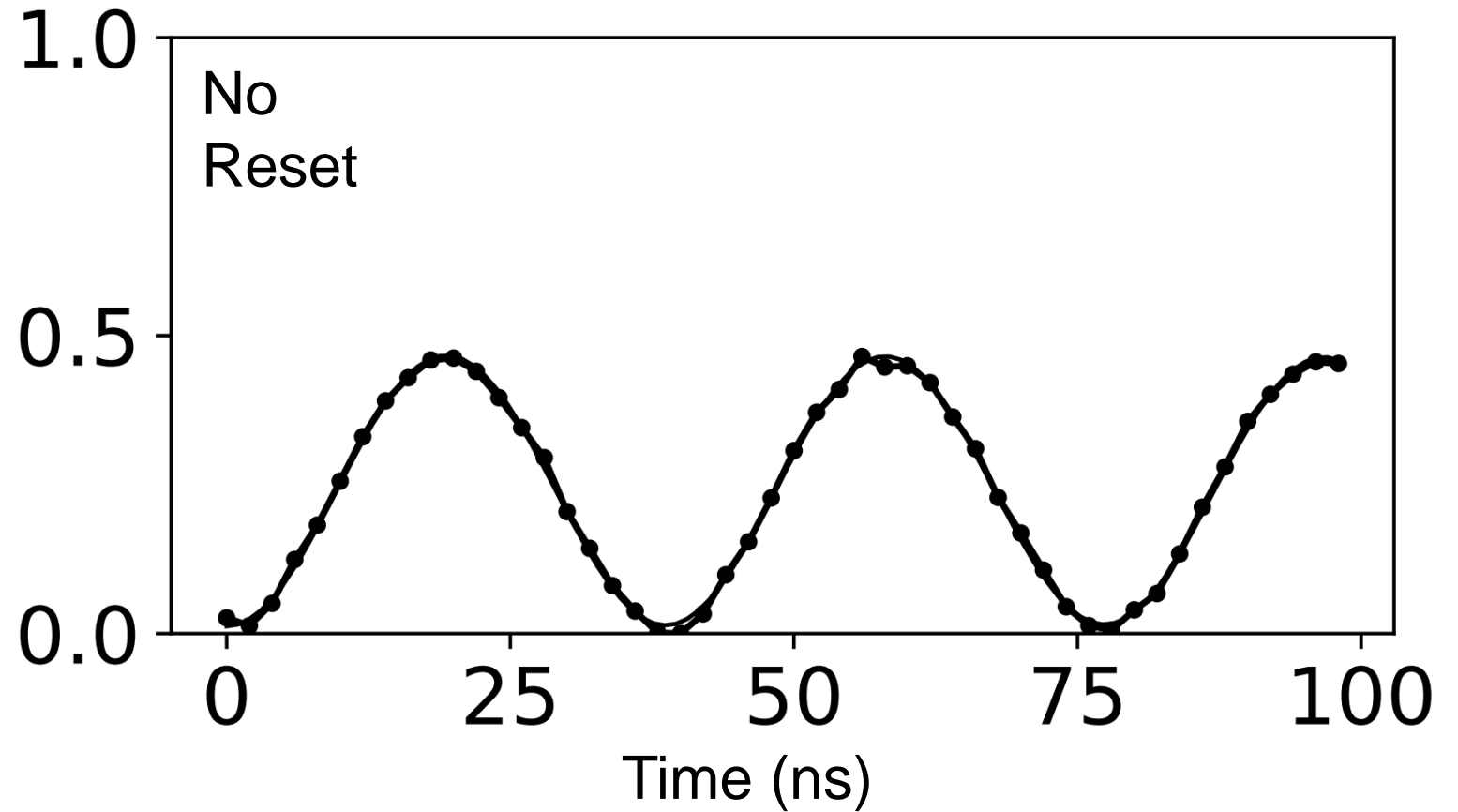
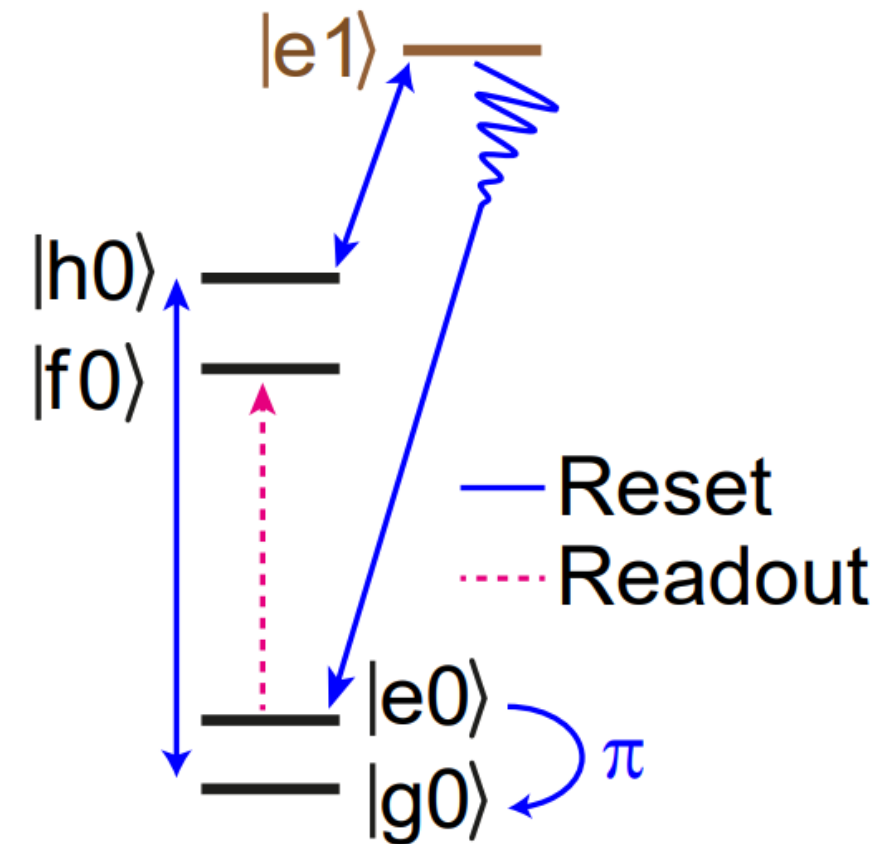


- Reset fidelity limited by cavity & $|f\rangle$ state population
- Perform a $|e\rangle \rightarrow |f\rangle \pi$ pulse before readout
- 50% readout fidelity
- In newer samples no longer necessary to excite for readout

Qubit initialization and readout



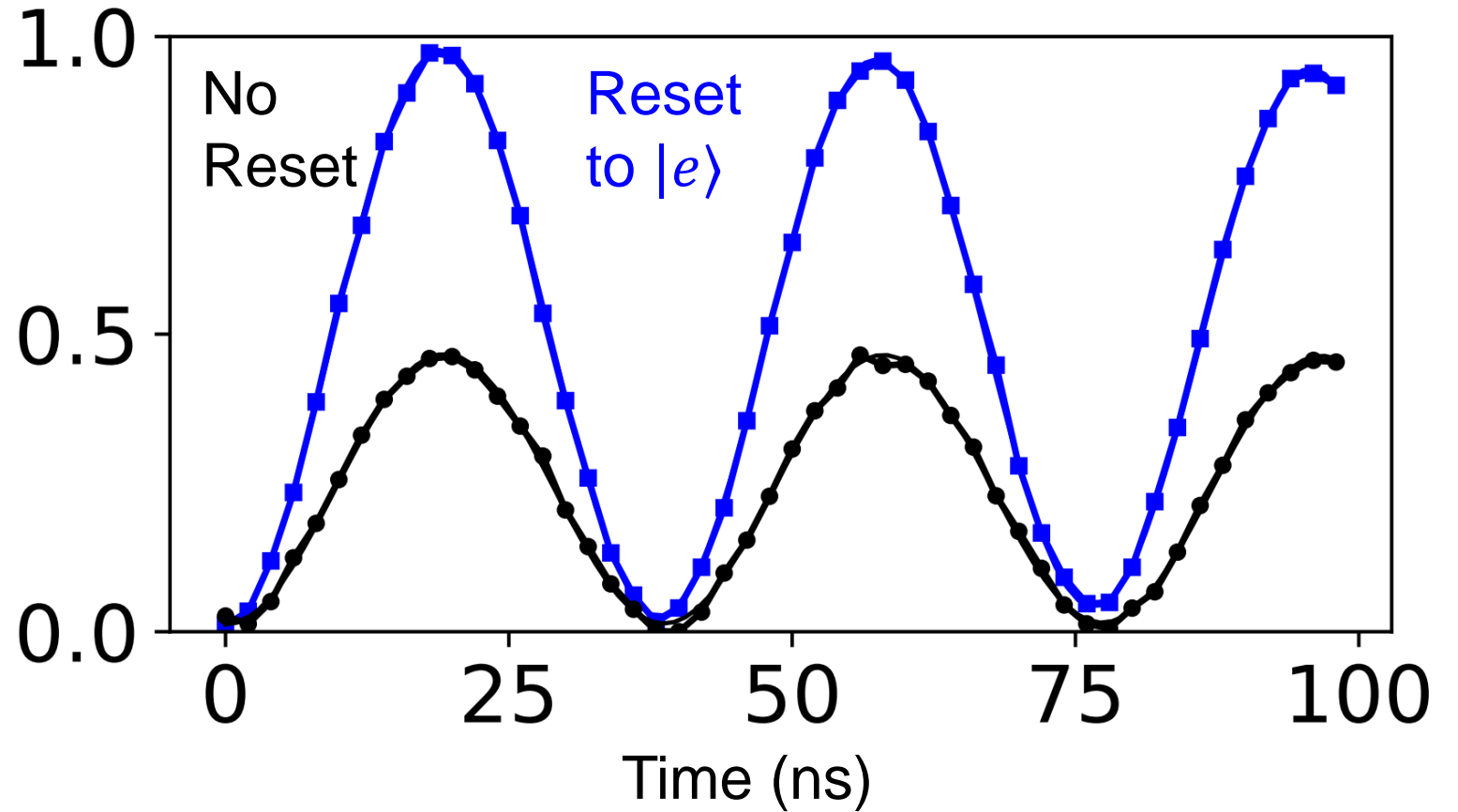
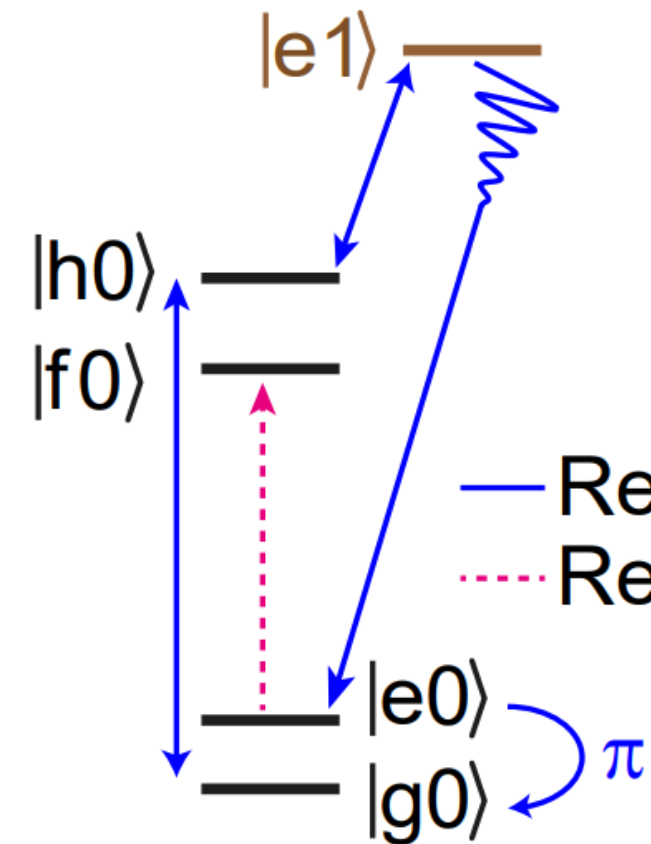
$|e\rangle \leftrightarrow |f\rangle$ Rabi Contrast



Qubit initialization and readout



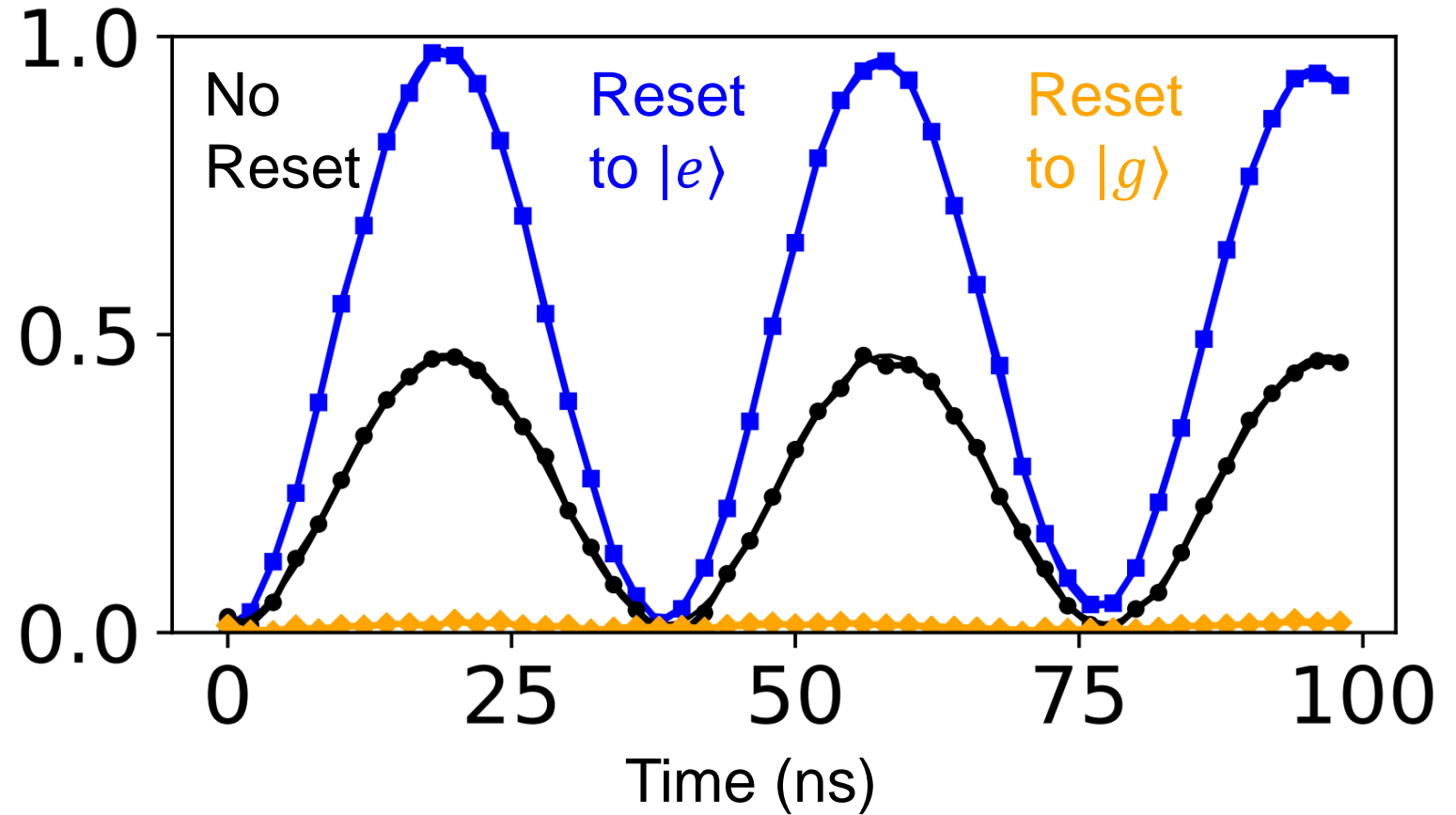
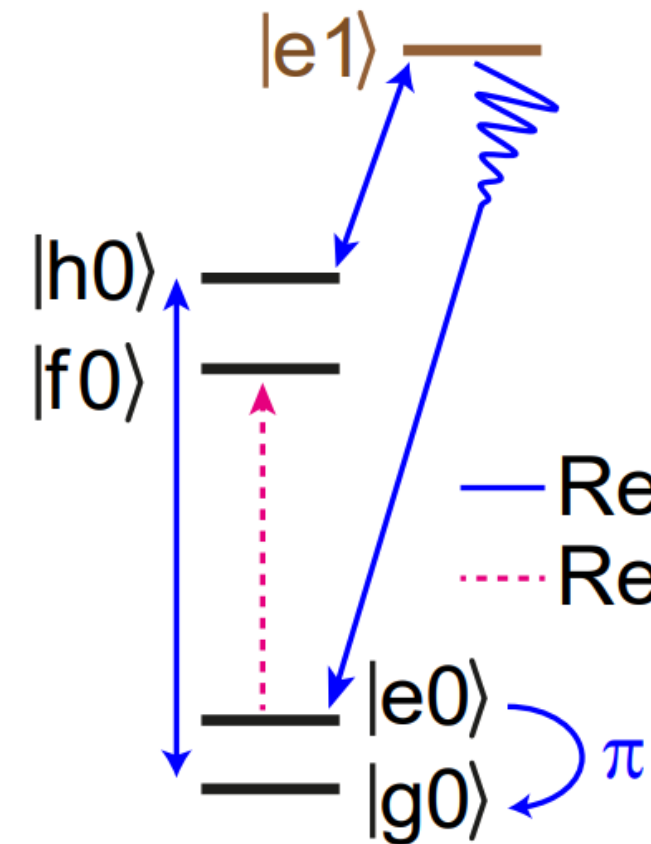
$|e\rangle \leftrightarrow |f\rangle$ Rabi Contrast



Qubit initialization and readout



$|e\rangle \leftrightarrow |f\rangle$ Rabi Contrast

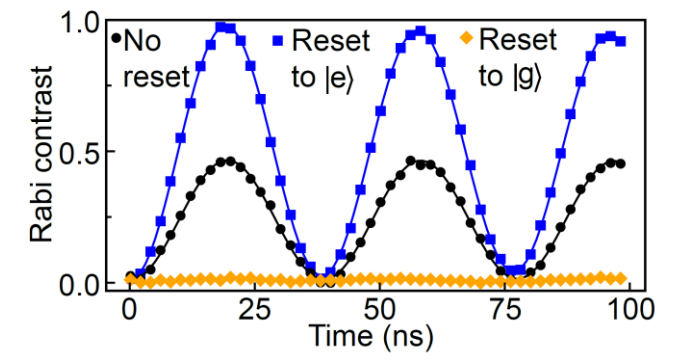
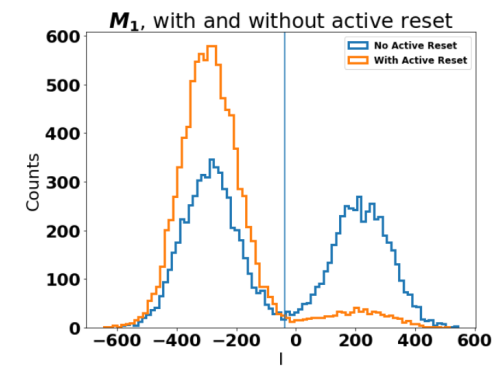


97% initial state fidelity, $T_q = 190 \mu K$

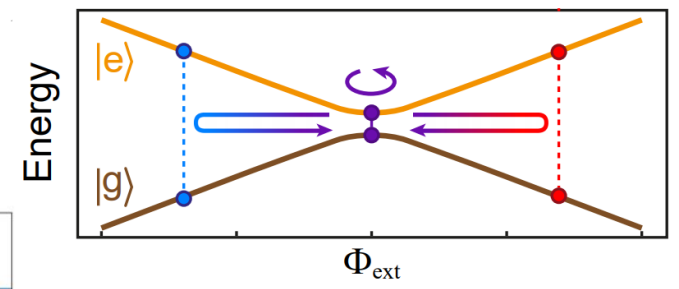
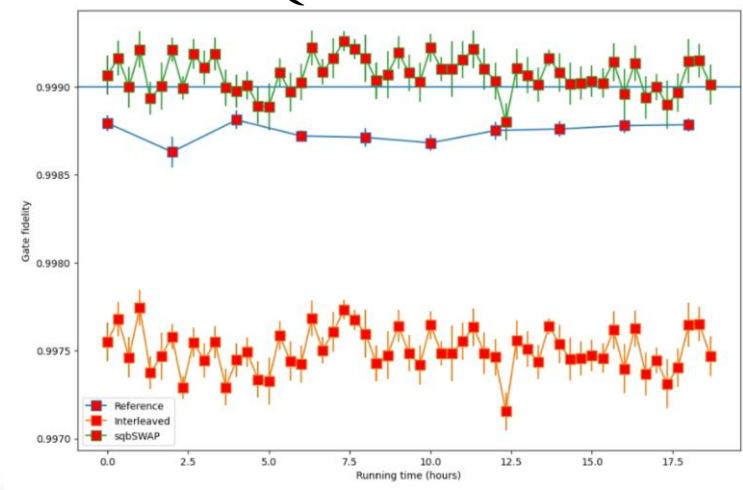
Rethinking paradigms



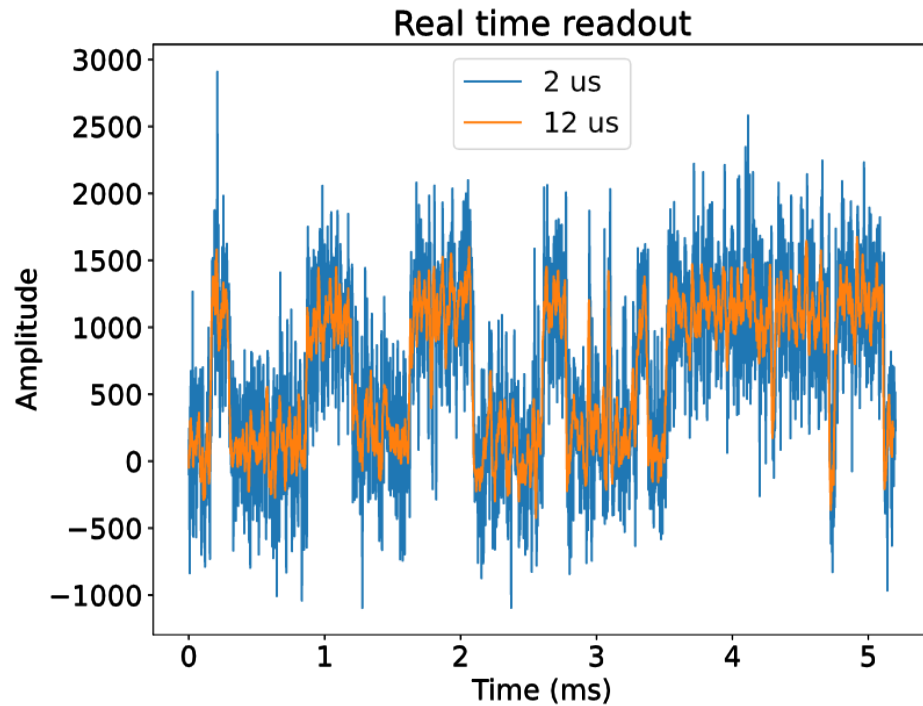
- Initialization via laser cooling
- **Initialization via active reset**
- Ultrafast single qubit gates
- High fidelity galvanically-coupled two qubit gates



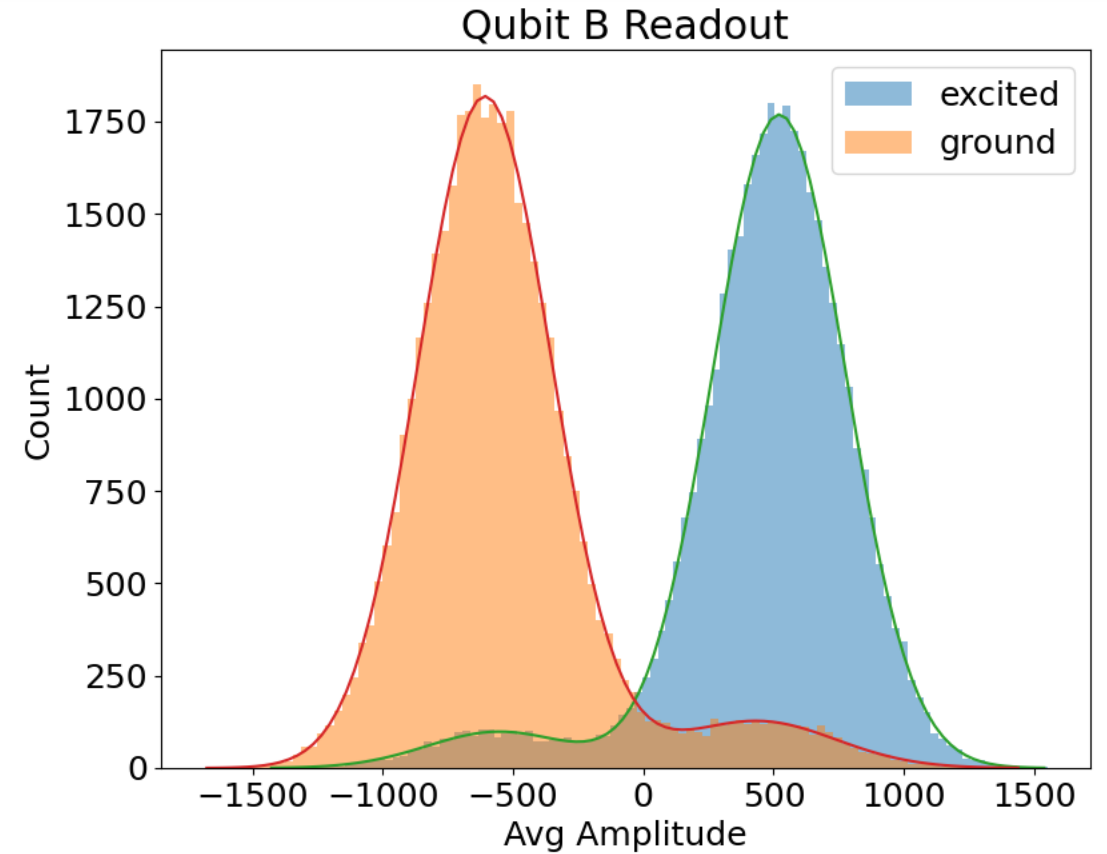
SQBSWAP XEB



Real time dispersive readout

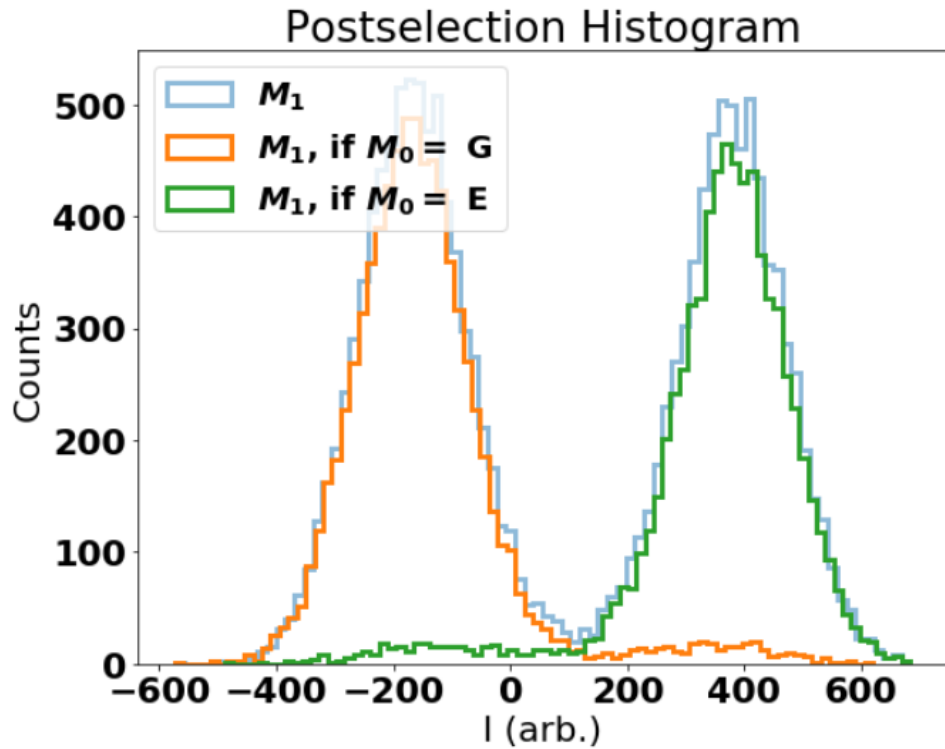
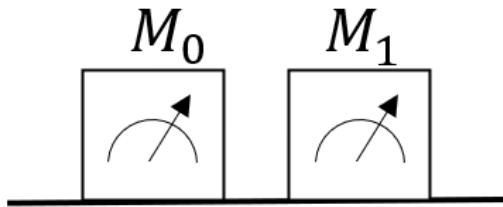


- Better readout design results in high fidelity readout
- Paramp would still speed up (and improve) measurement, state preparation

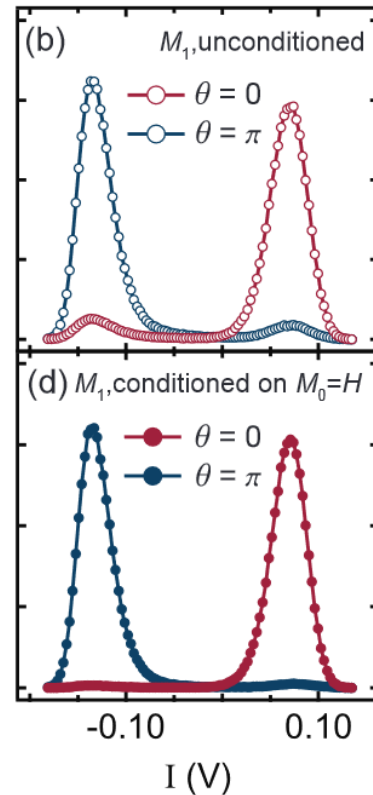
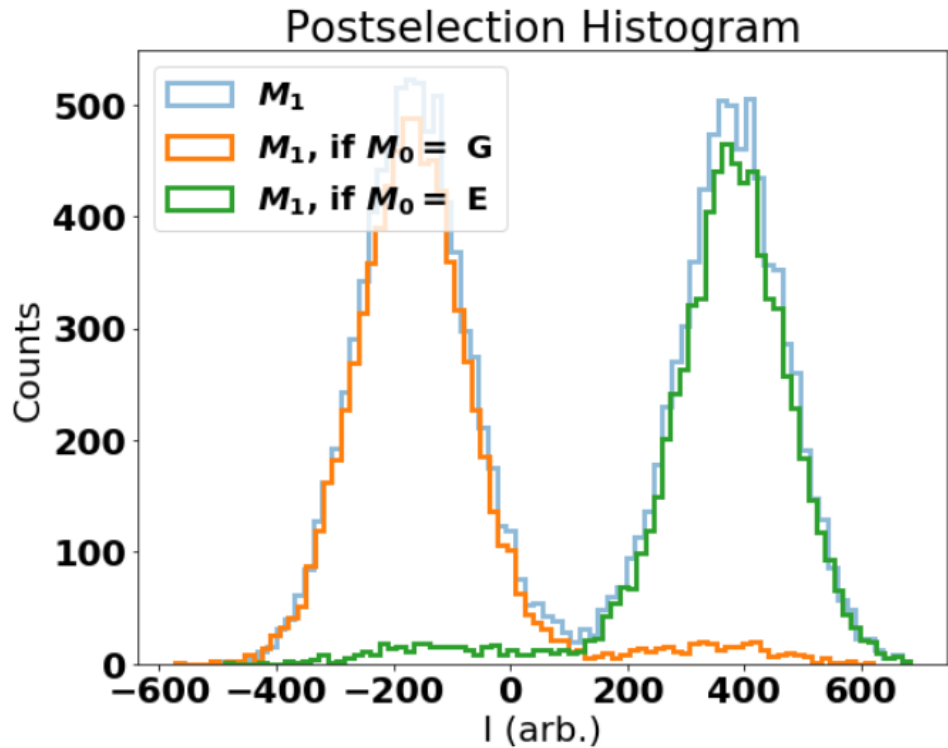
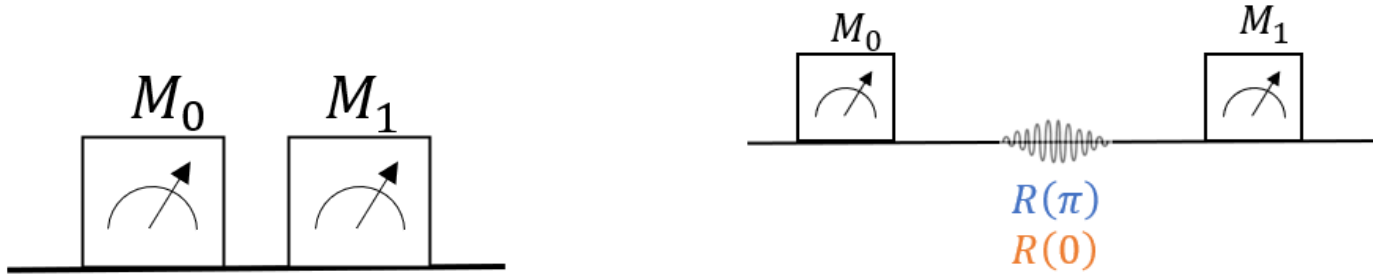


- Qubit prepared by measurement w/ feedback
- Statistical infidelity $\sim 2\%$
- Infidelity of prepared states $\sim 18\%$ ($\sim .98\text{mK}$)

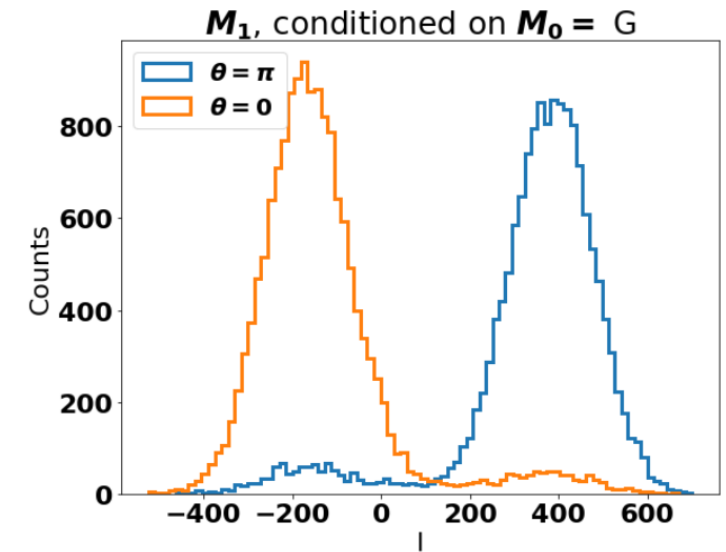
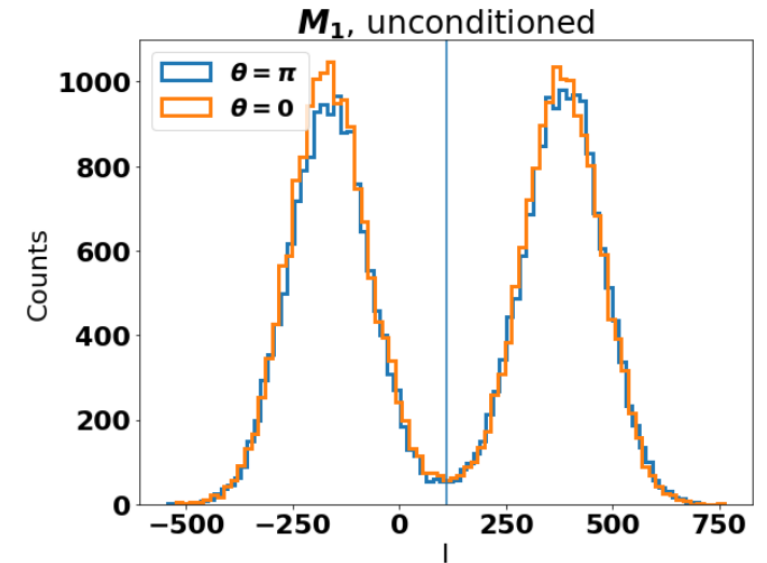
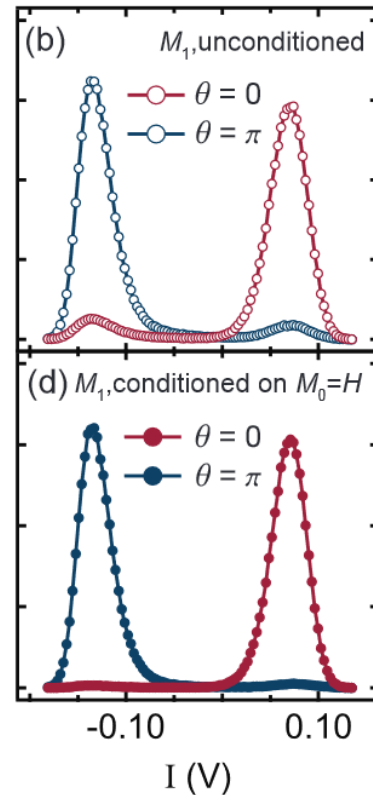
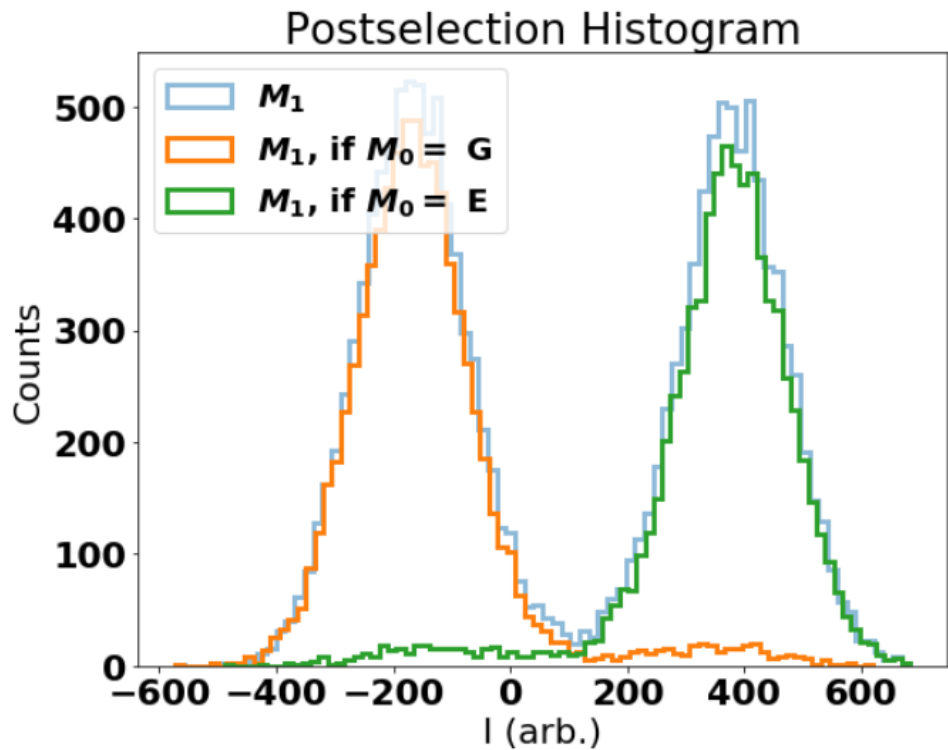
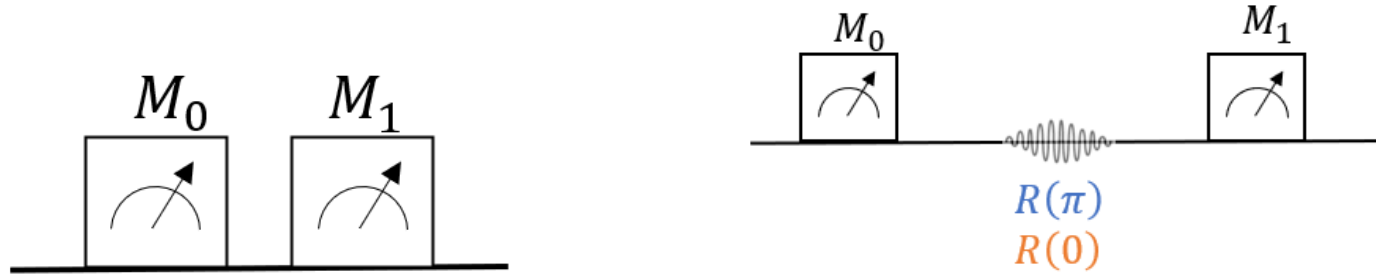
Demonstration of post-selection protocol



Demonstration of post-selection protocol



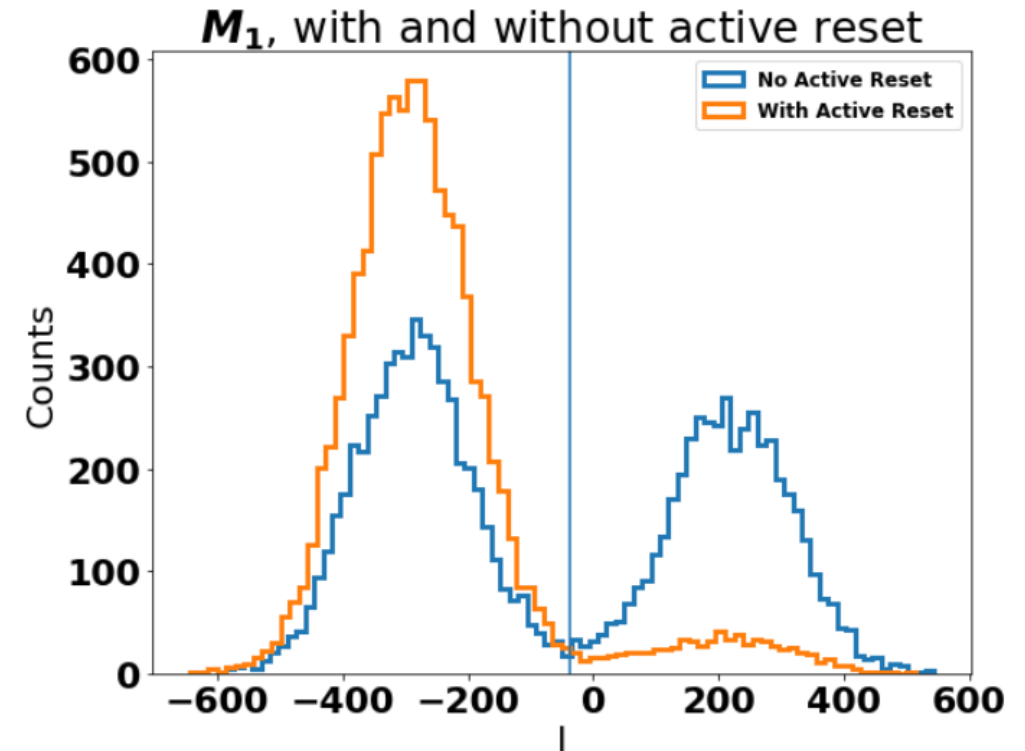
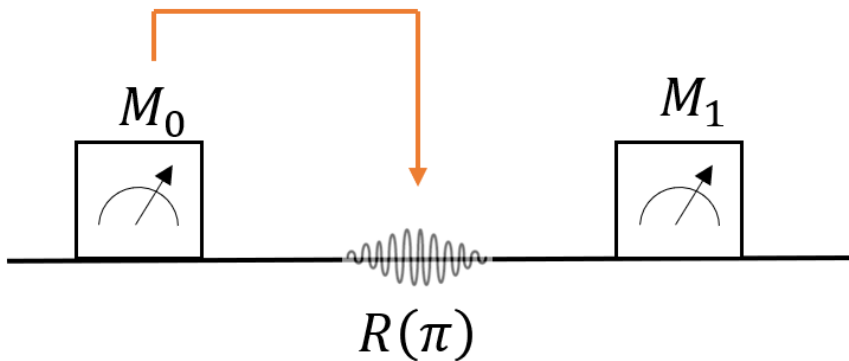
Demonstration of post-selection protocol



Executing the active reset protocol with the QICK



Stefanazzi 2021 arXiv: 2110.00557



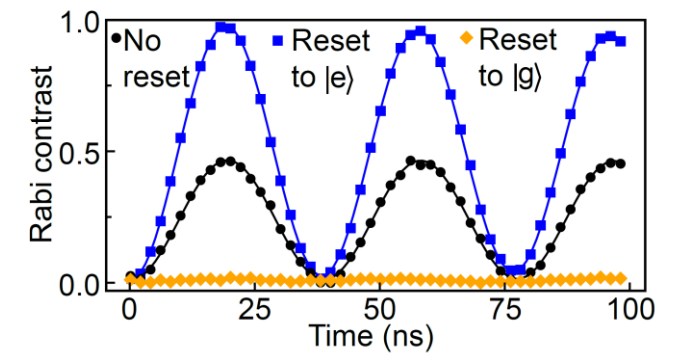
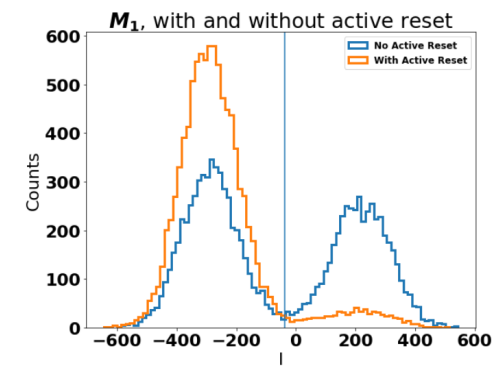
Firmware and software is published and **open-source**
<https://github.com/openquantumhardware/qick>

Developed with team at Fermi National Lab
Review of Scientific Instruments **93**, 044709 (2022)

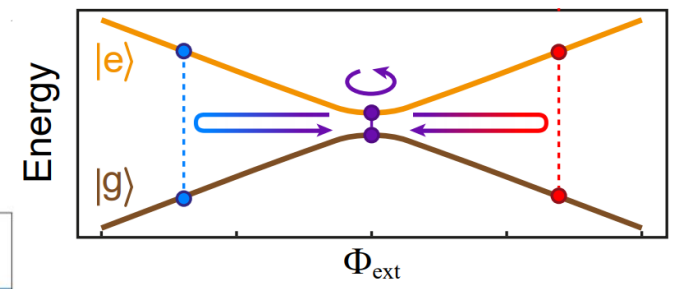
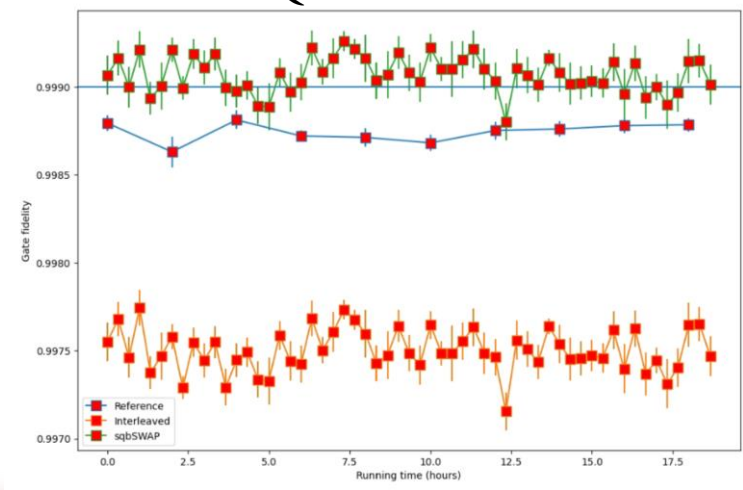
Rethinking paradigms



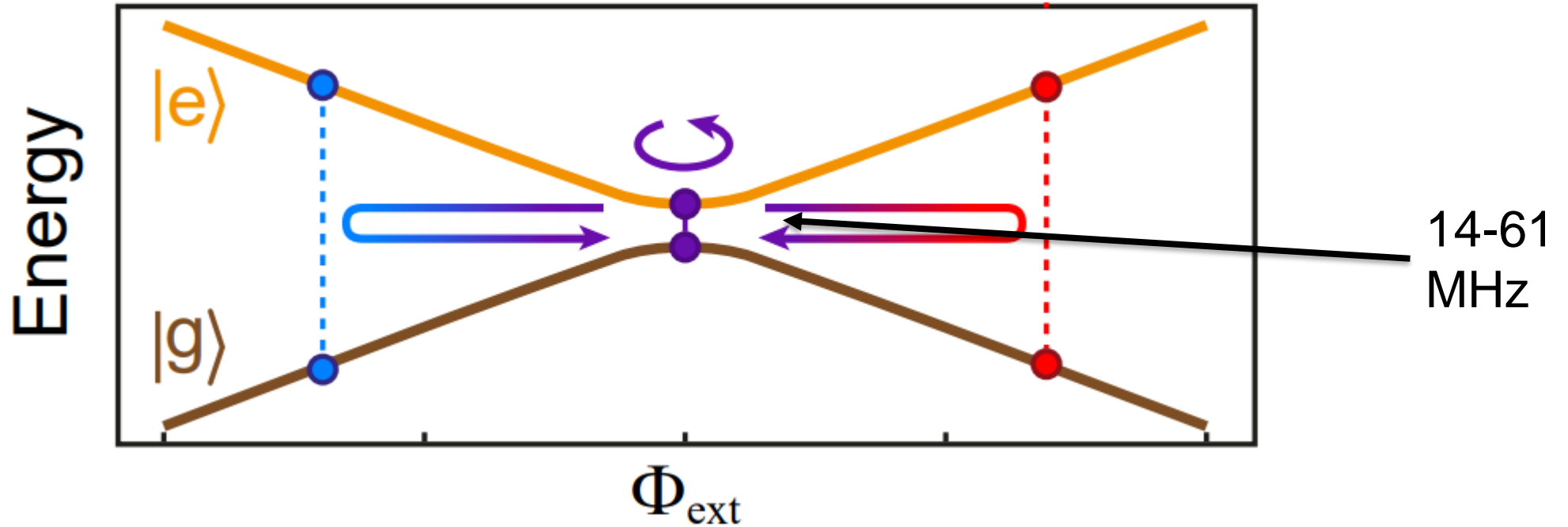
- Initialization via laser cooling
- Initialization via active reset
- **Ultrafast single qubit gates**
- High fidelity galvanically-coupled two qubit gates



SQBSWAP XEB



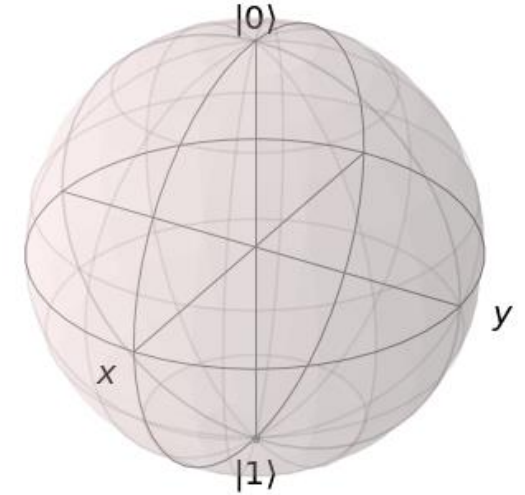
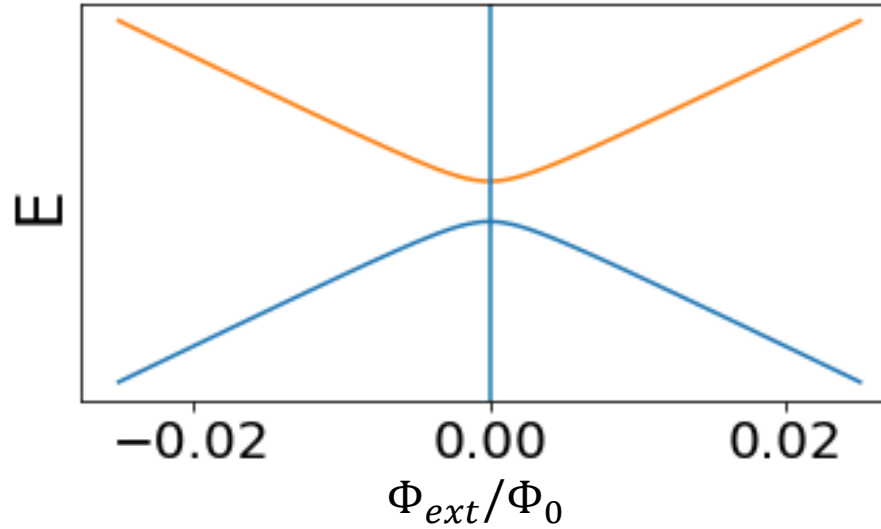
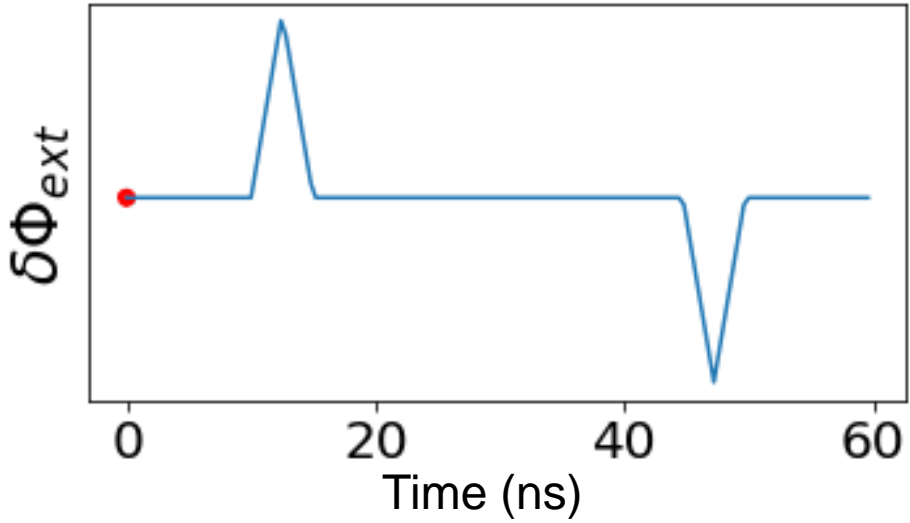
Fast single-cycle flux gates



$$H/h = A\delta\Phi_{ext}\sigma_x + \frac{\omega_q}{2}\sigma_z$$

It's easy to have $A\delta\Phi_{ext} \gg \frac{\omega_q}{2}$

Fast single-cycle flux gates

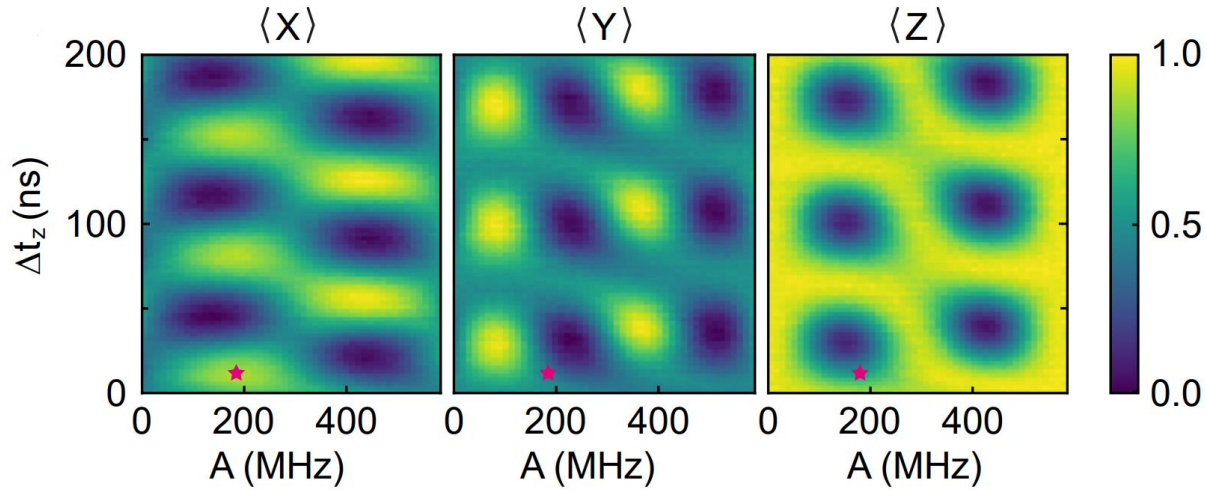


- Lab frame
- Finished within a qubit cycle
- zero total net flux, effective echo for low frequency noise

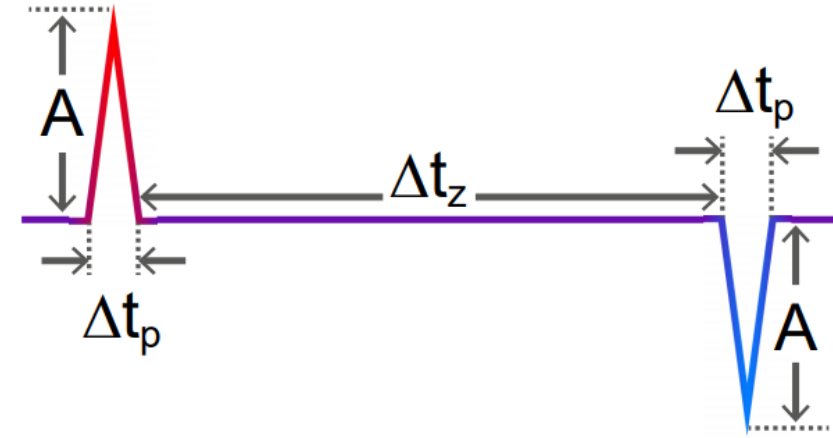
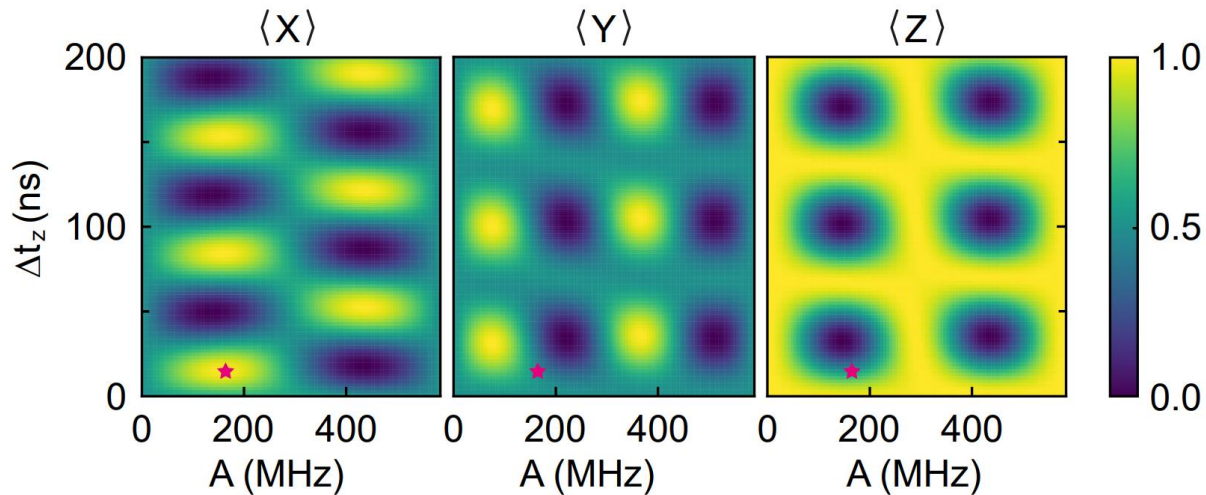
Fast single-cycle flux gates



Experiment



Simulation



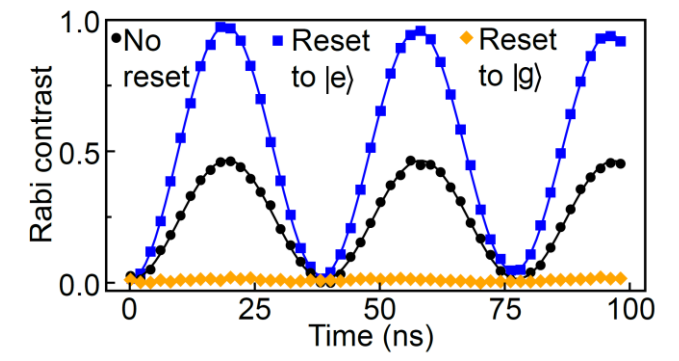
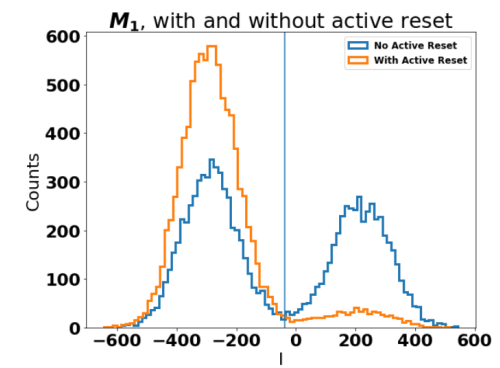
Duration (about 4x faster than previous gen)
Y/2: 5.5 ns **Z/2**: 6ns **Y**: 8ns (total pulse)

Qubit frequency **100x slower** than transmon,
gate speed **2x faster!**

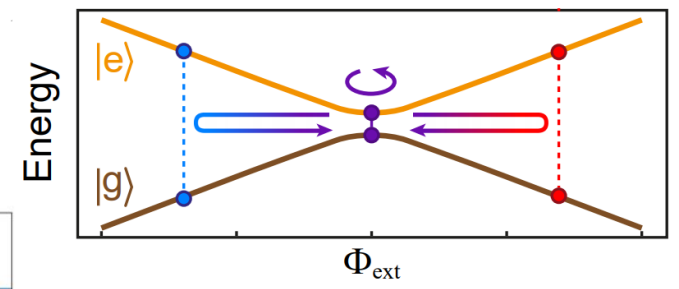
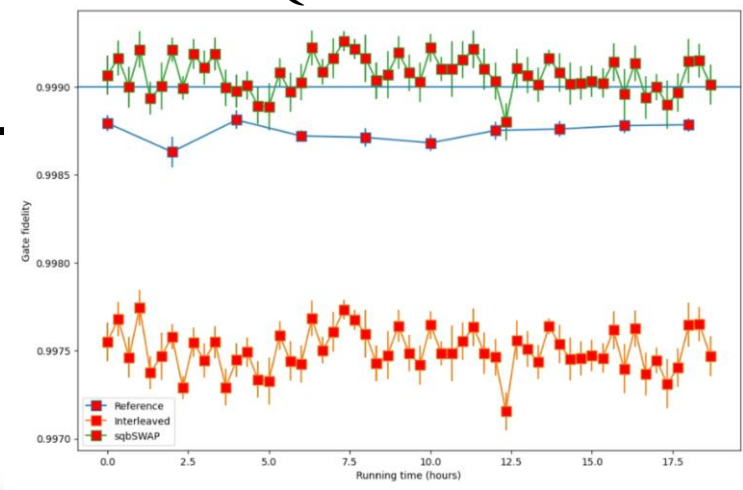
Rethinking paradigms



- Initialization via laser cooling
- Initialization via active reset
- Ultrafast single qubit gates
- **High fidelity galvanically-coupled two qubit gates**



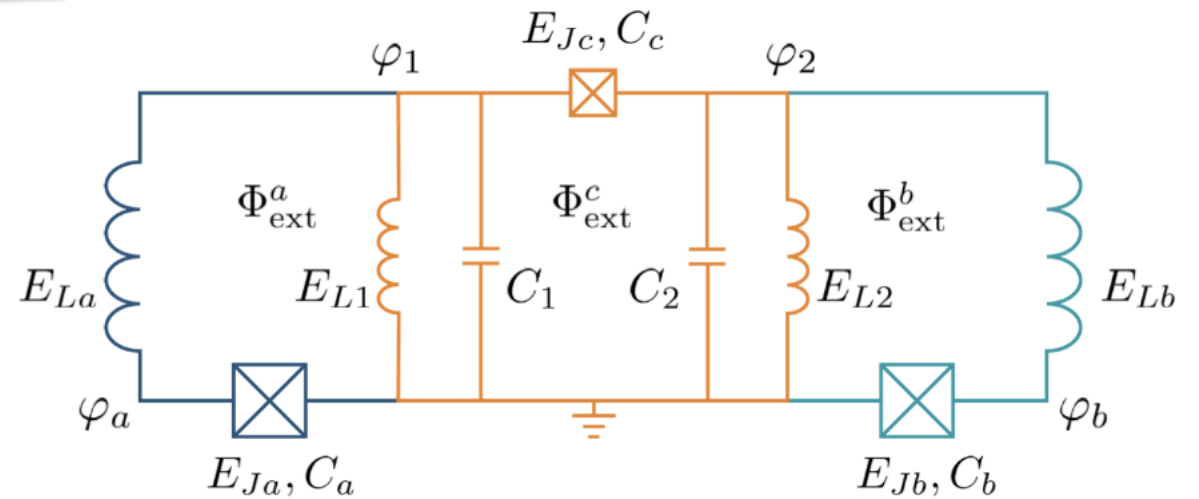
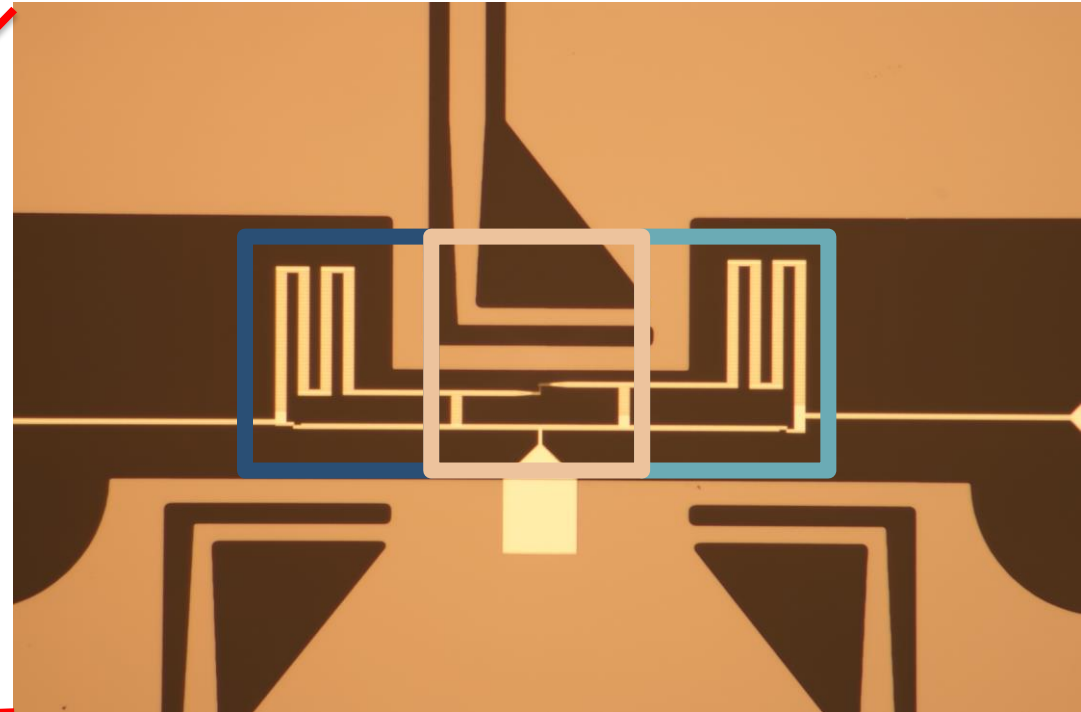
SQBSWAP XEB



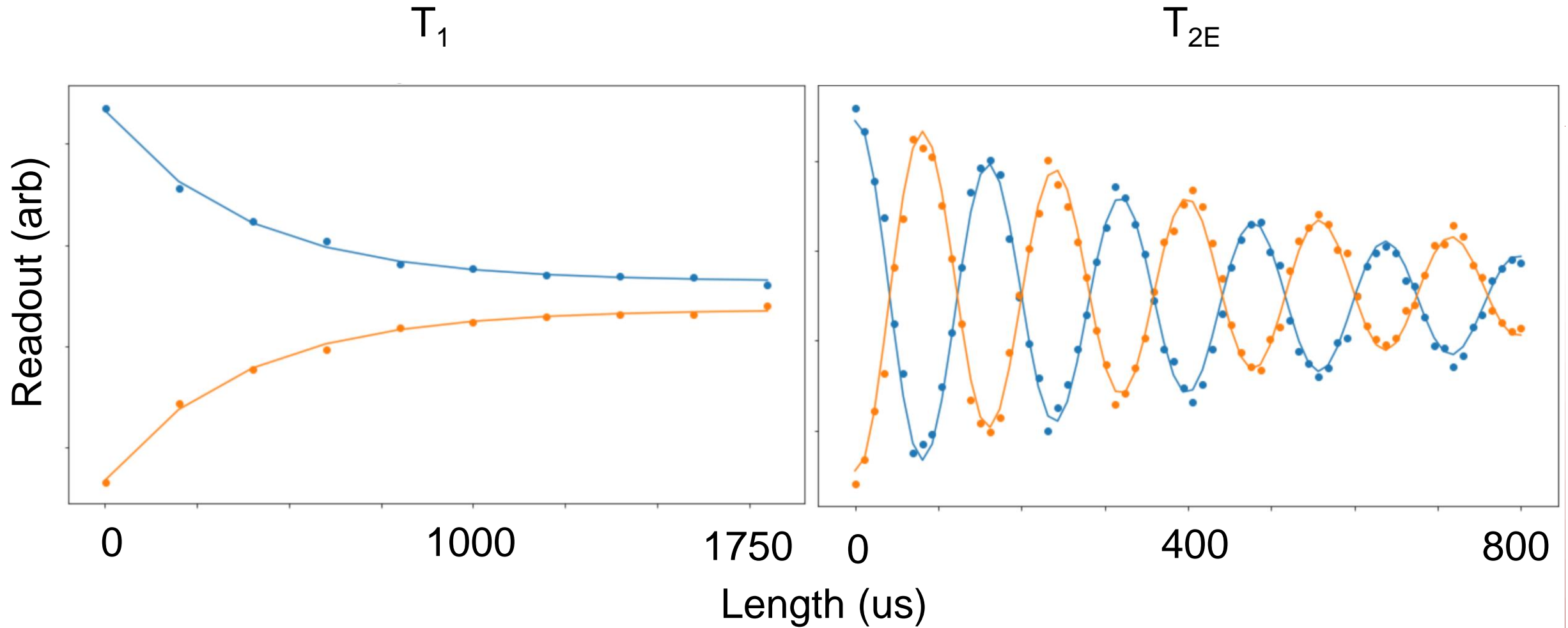
Optical Image of the device



Ta substrate



Coherences – best case (qB)

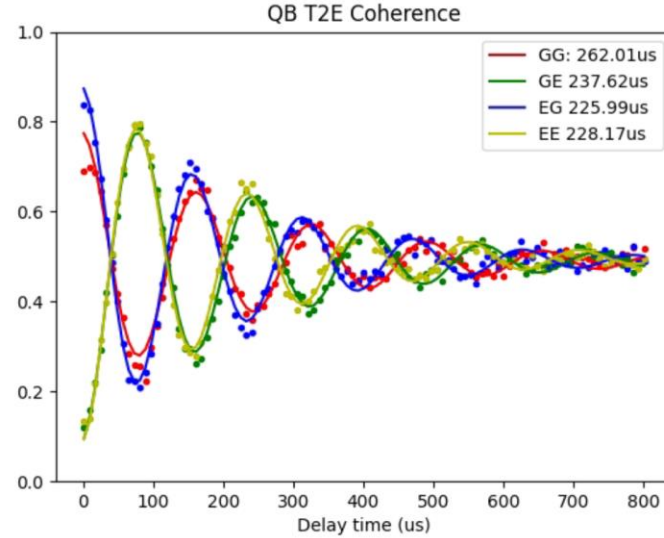
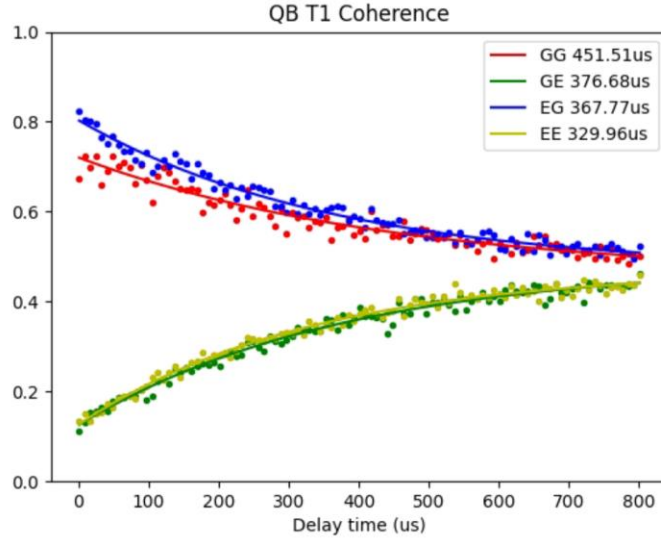
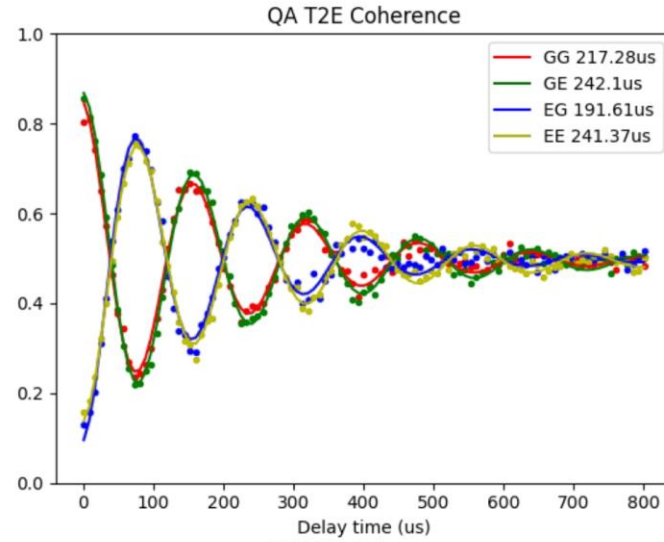
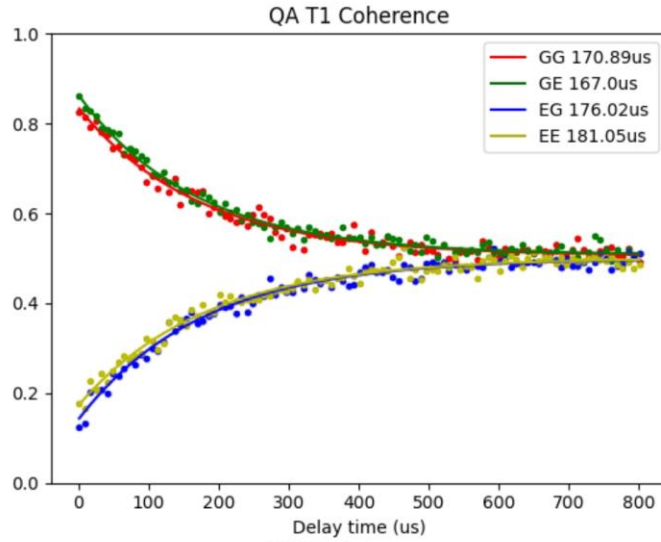


$T_1 \sim 370 \mu\text{s}$, $T_{2E} \sim 580 \mu\text{s}$

Coherences – working point

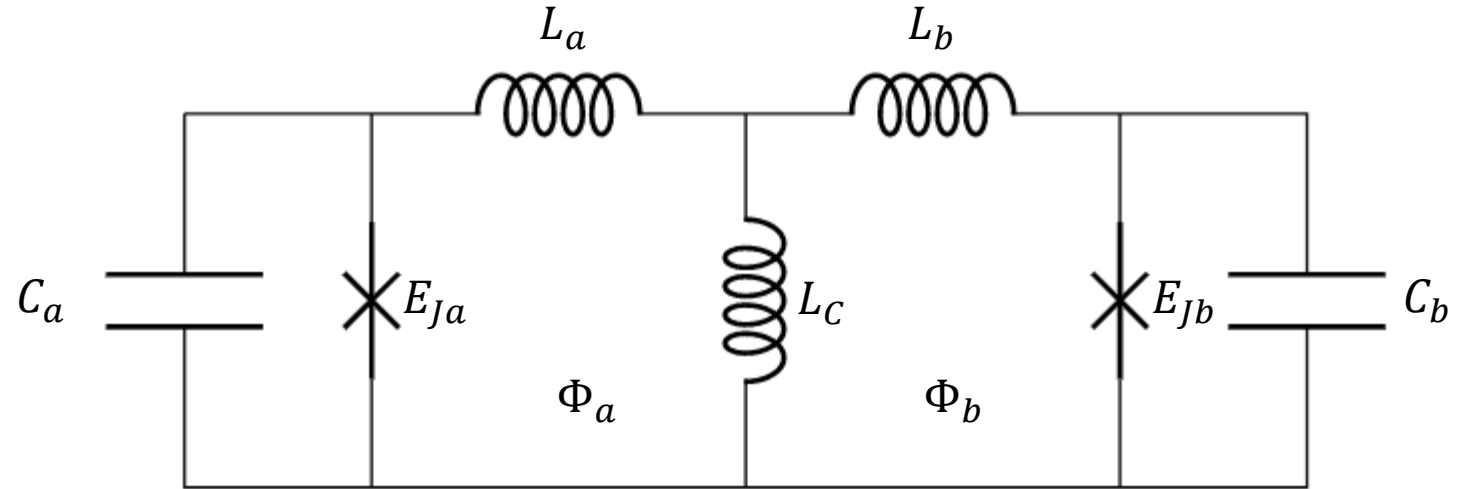
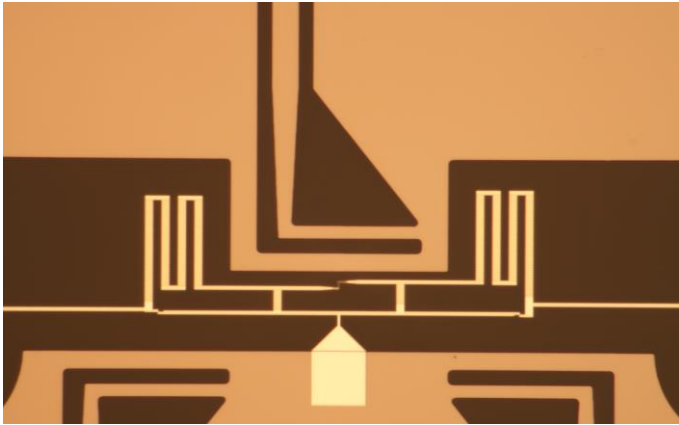


State population $P(g)$



	Freq (MHz)	$T_1(us)$	$T_{2E}(us)$
qA	48.45	173.74	223.09
qB	61.76	381.48	238.45

Two inductively coupled fluxoniums



Two qubits are coupled via a shared inductor

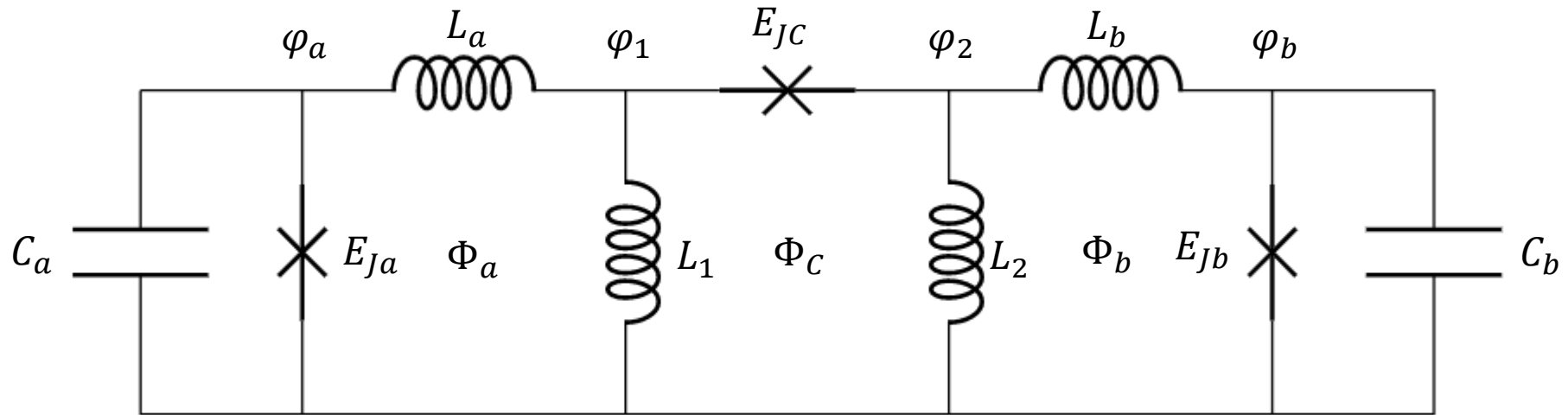
$$H \approx H_a + H_b + \frac{L_c}{L_a} E_L \varphi_a \varphi_b$$

In computational basis looks like

$$H_c \approx 4\pi^2 \frac{L_c}{L_a} E_L \sigma_x \sigma_x$$

Problem: It's always on!

Tunable Coupler – Circuit Analysis

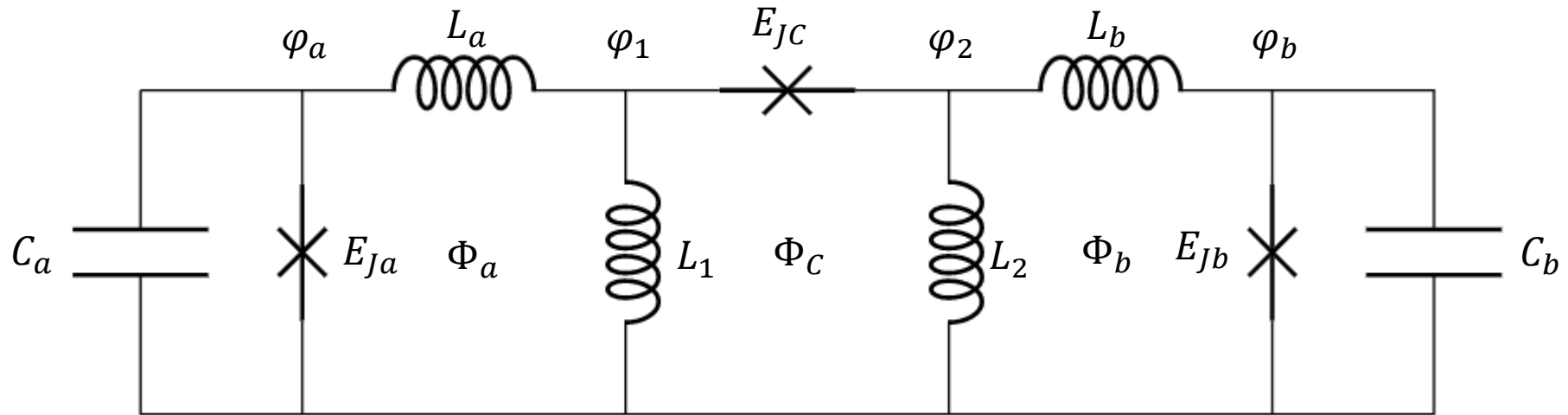


Now the shared inductance is tunable with coupler flux Φ_C

To the second order, it gives a tunable coupling term from 0 to

$$\sim E_{JC} \frac{E_{La} E_{Lb}}{E_{L1} E_{L2}} \varphi_a \varphi_b$$

Tunable Coupler – Circuit Analysis



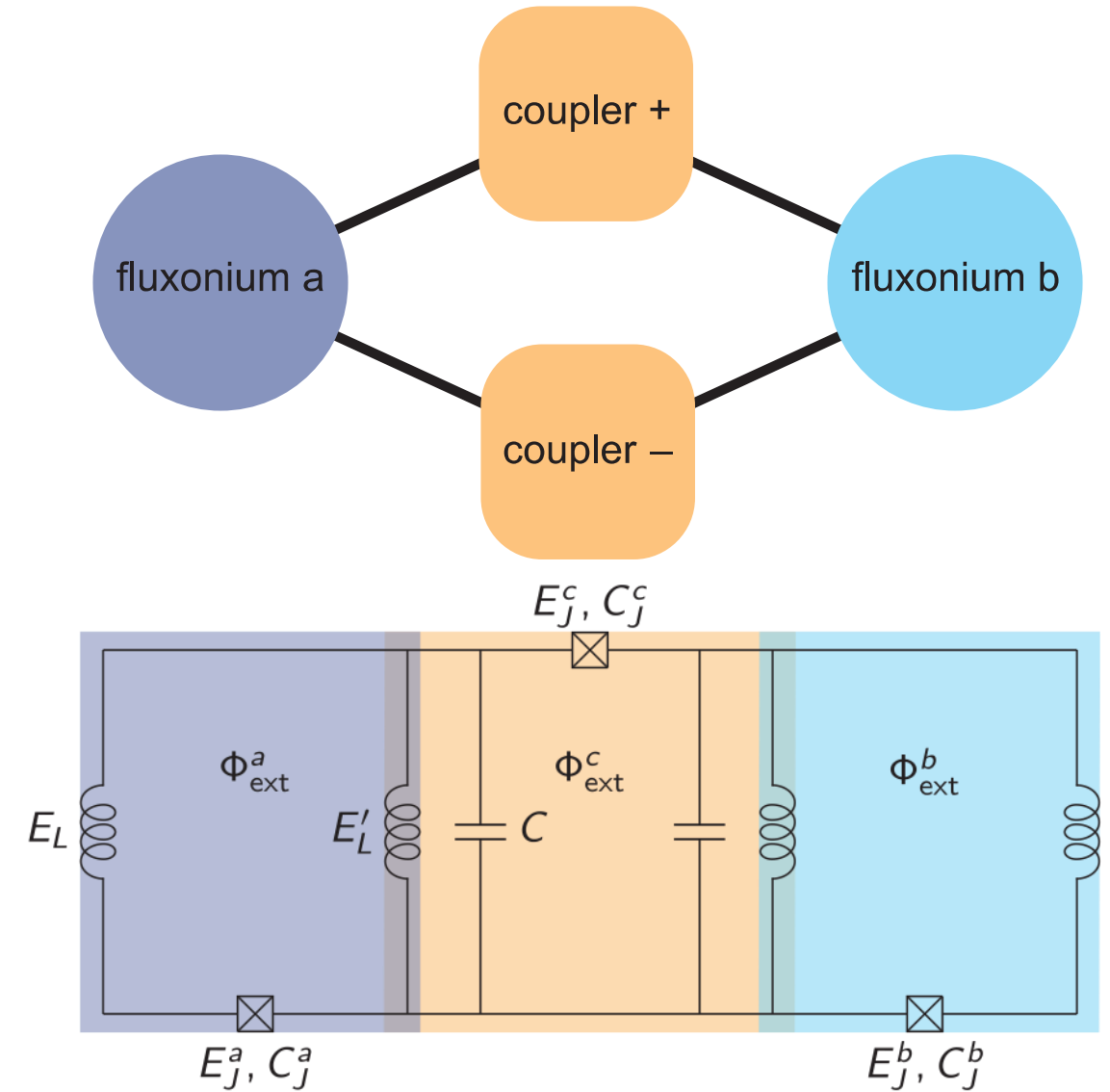
Decomposing the two qubit Hamiltonian into single qubit bases:

$$H = \omega_1 \sigma_{z1} + \omega_2 \sigma_{z2} + A_1(\Phi_a, \Phi_C) \sigma_{x1} + A_2(\Phi_b, \Phi_C) \sigma_{x2} + J(\Phi_C, \Phi_a, \Phi_b) \sigma_{x1} \sigma_{x2}$$

Tunable-coupled fluxonium



- Couple two fluxonium qubits galvanically/inductively:
 - Tunable coupling
 - Exactly cancel coupling
 - Coupling strength can rival single qubit energy



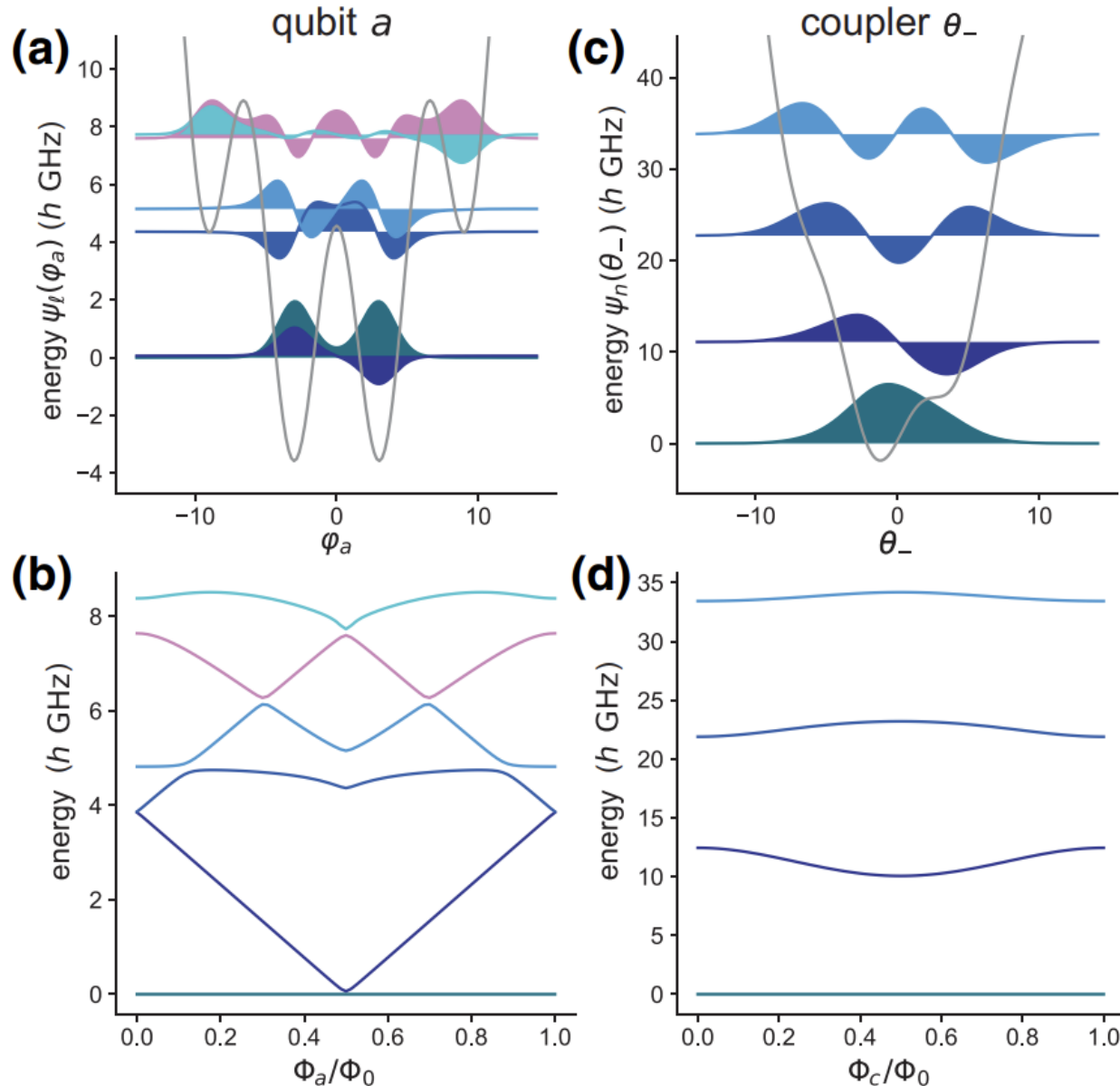
experiment: capacitive coupling

- F. Bao *et al.* arXiv:2111.13504 (2021)
- Q. Ficheux *et al.* PRX **11**, 21026 (2021)
- H. Xiong *et al.* arXiv:2103.04491 (2021)
- L. Ding *et al.* arXiv: 2304.06087 (2023)

theory: capacitive coupling

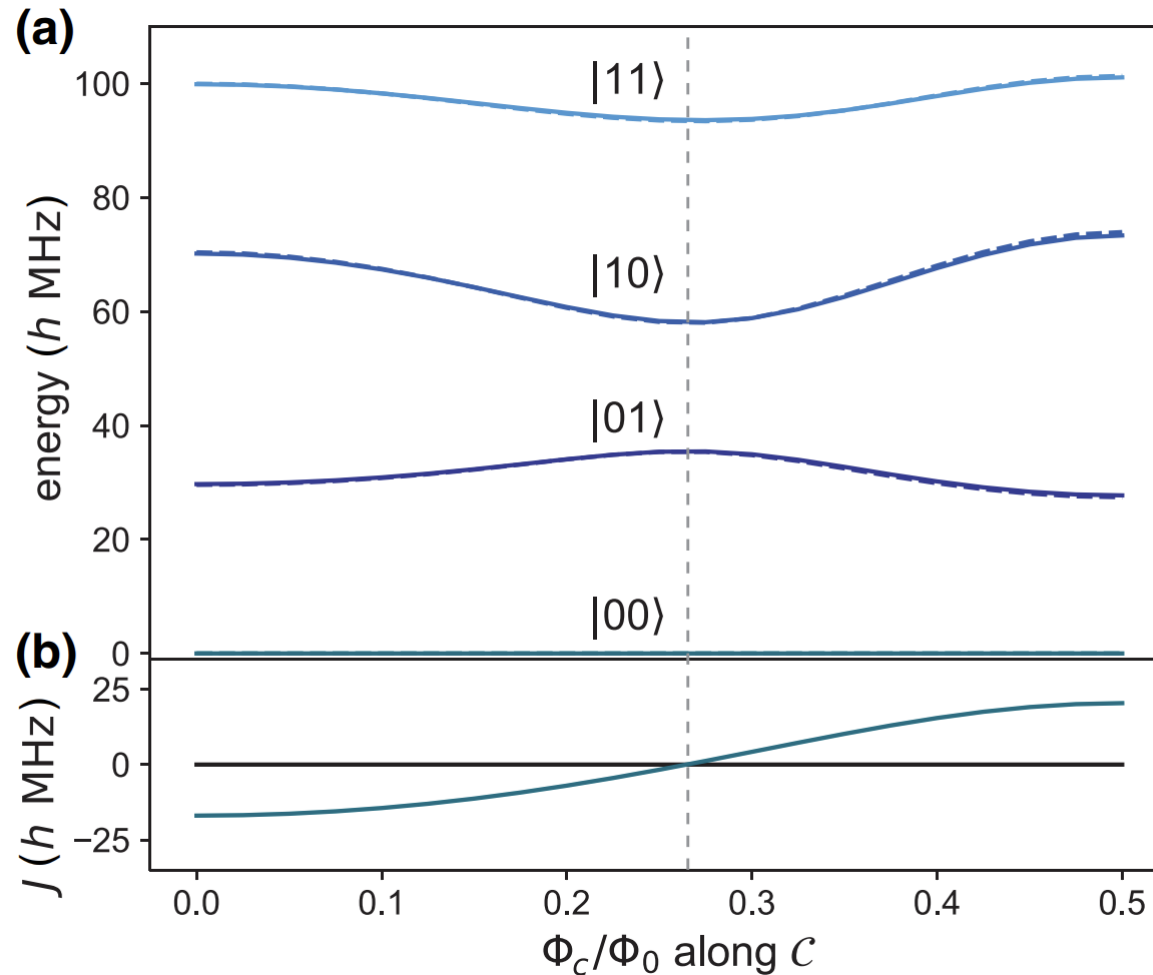
- Y. Chen *et al.* arXiv:2110.00632 (2021)
- K. N. Nesterov *et al.* PRX Quantum **2**, 020345 (2021)
- I. N. Moskalenko *et al.* arXiv:2107.11550 (2021)

Tunable Coupler – Circuit Analysis



- Coupler $E_L \gg$ Qubit E_L
- Lowest coupler mode frequency ~ 10 GHz
- $J = J_+ + J_-$, thus when $J_+ = -J_-$, we turn off all coupling

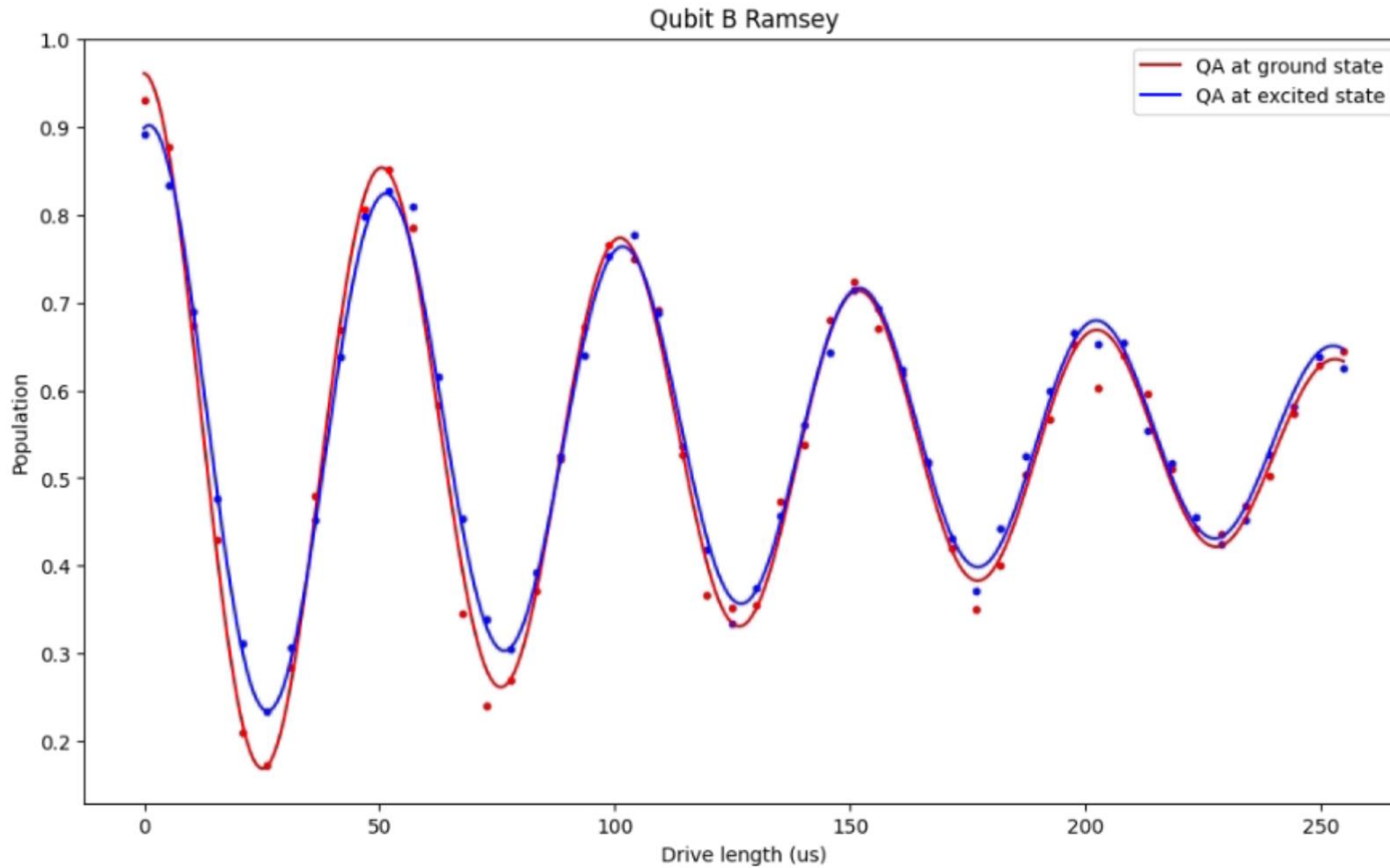
Tunable Coupler – Couplings vs coupler flux



Effective Hamiltonian

$$H = \omega_1 \sigma_{z1} + \omega_2 \sigma_{z2} + A_1(\Phi_a, \Phi_c) \sigma_{x1} + A_2(\Phi_b, \Phi_c) \sigma_{x2} + J(\Phi_c, \Phi_a, \Phi_b) \sigma_{x1} \sigma_{x2}$$

ZZ coupling measurement by Qubit B Ramsey

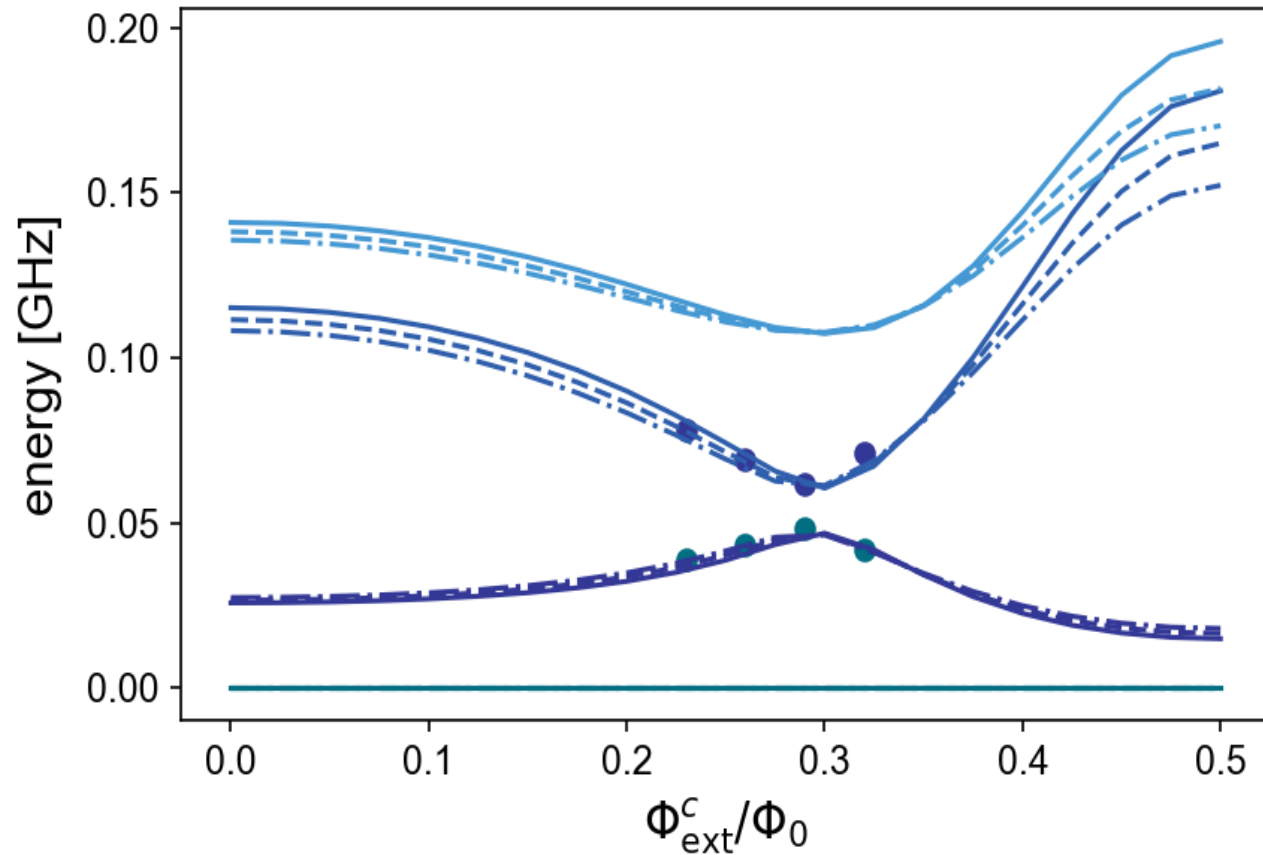


$$T_2^* \approx 180 \mu\text{s} \quad ZZ \approx 0.2 \text{ kHz}$$

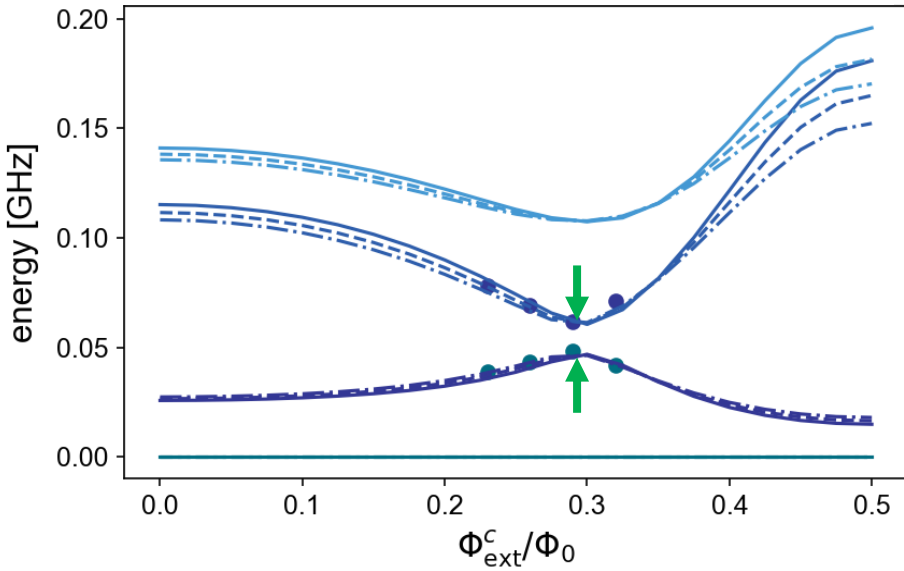
Coupler parameters



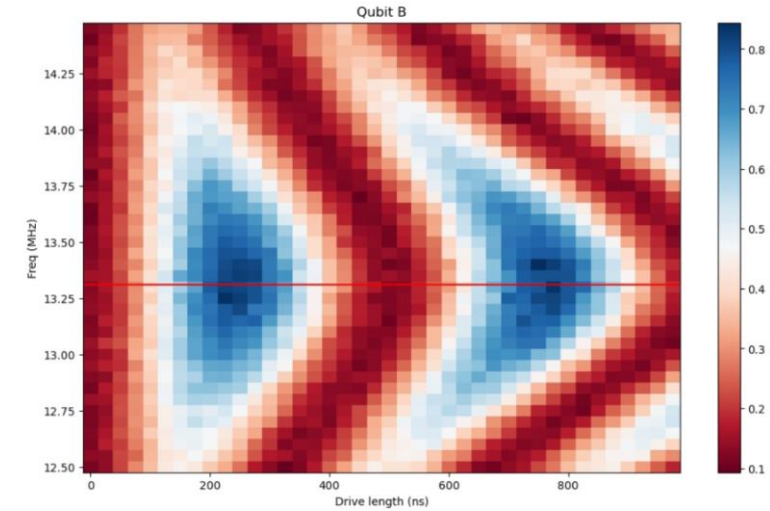
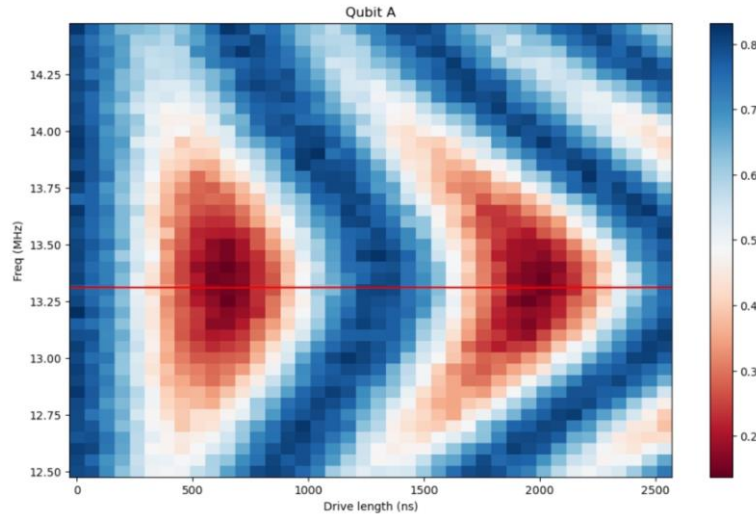
Checking coupler parameters by varying coupler E_c from -5% to 5%



Inducing a parametric interaction – i Swap

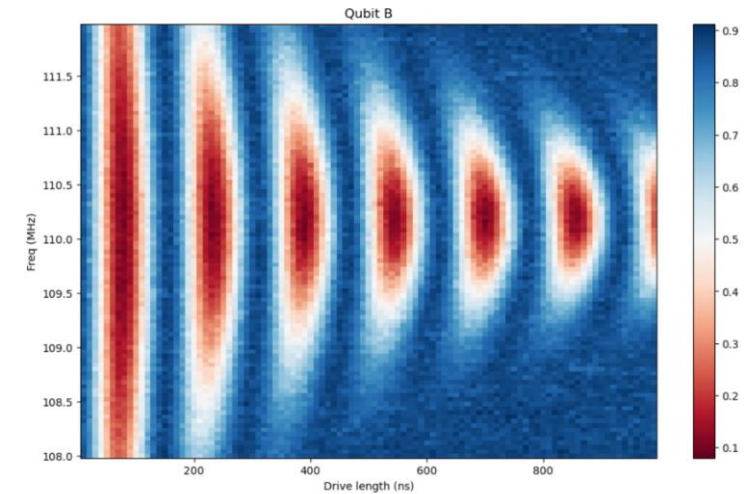
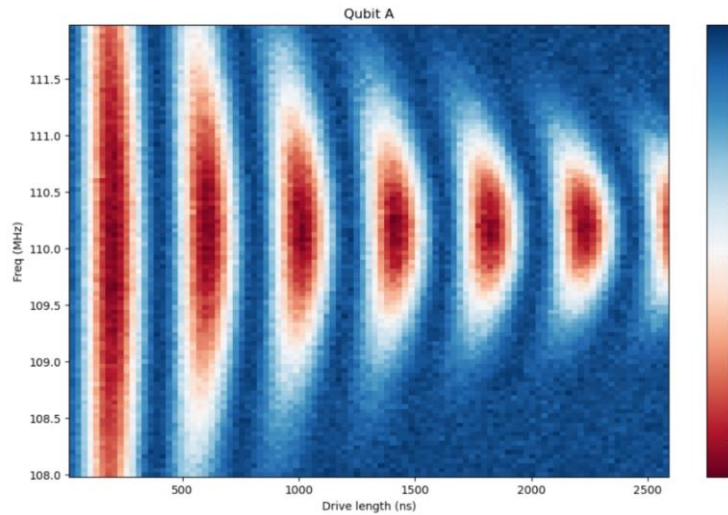
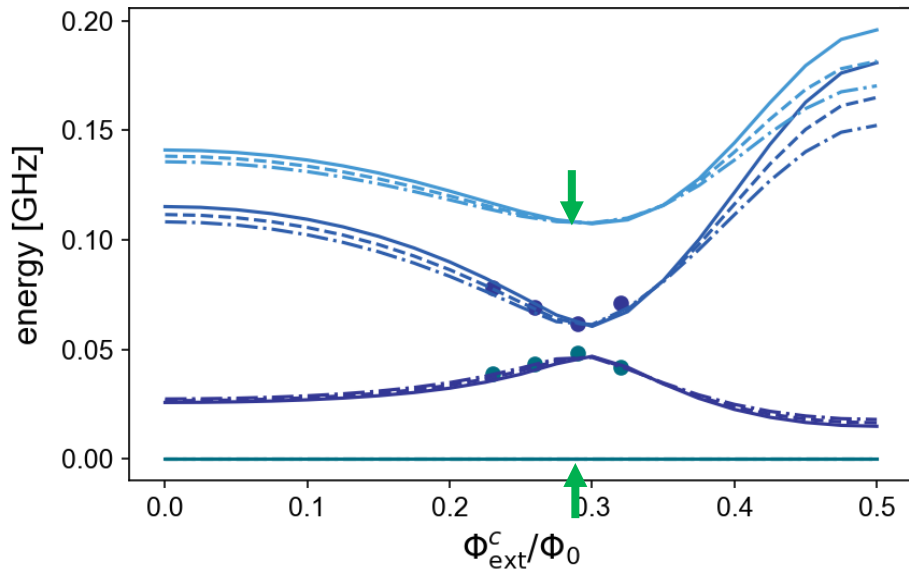


$$\omega_d = \omega_B - \omega_A$$



\sqrt{iSWAP} gate length = 170ns

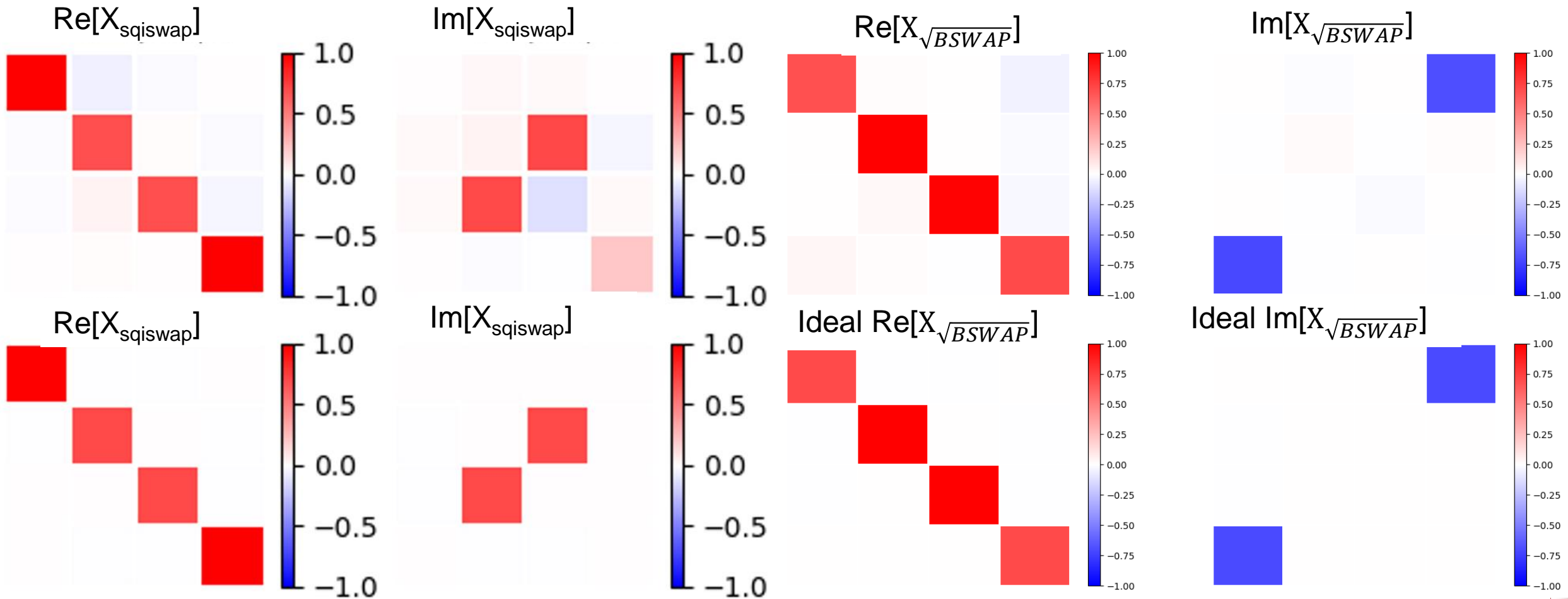
Inducing a parametric interaction - *b*Swap



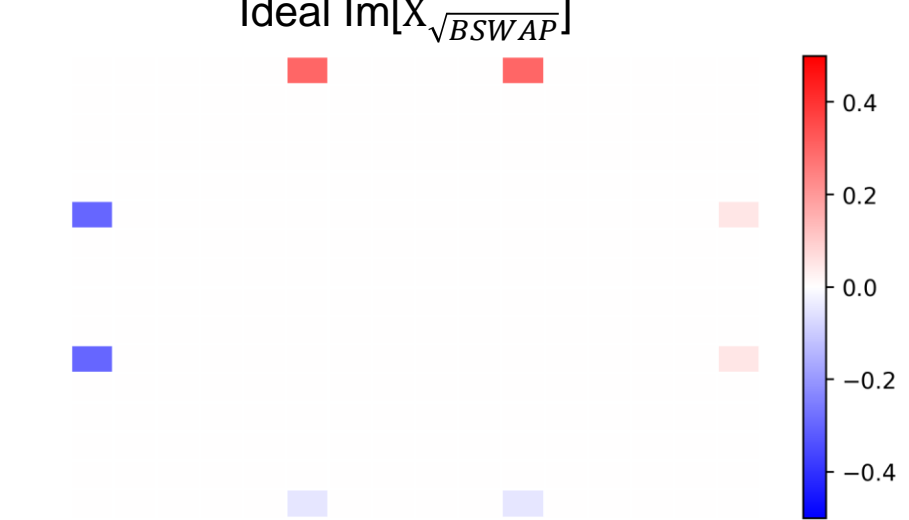
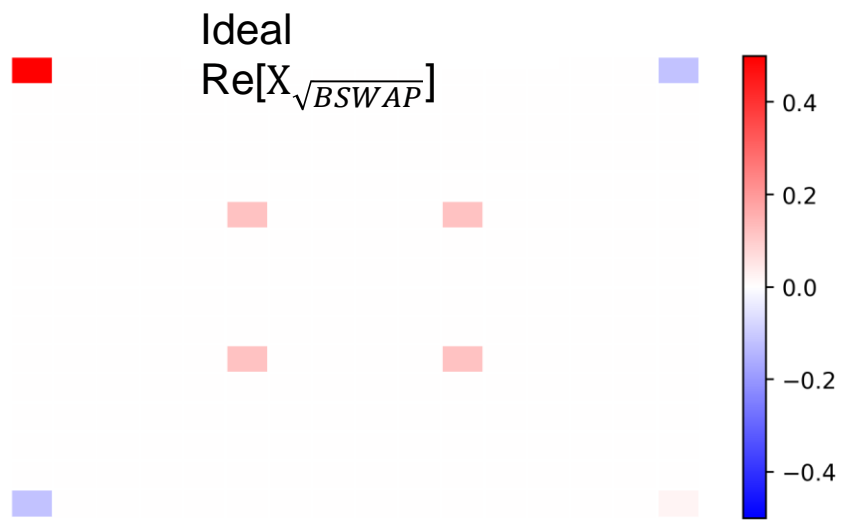
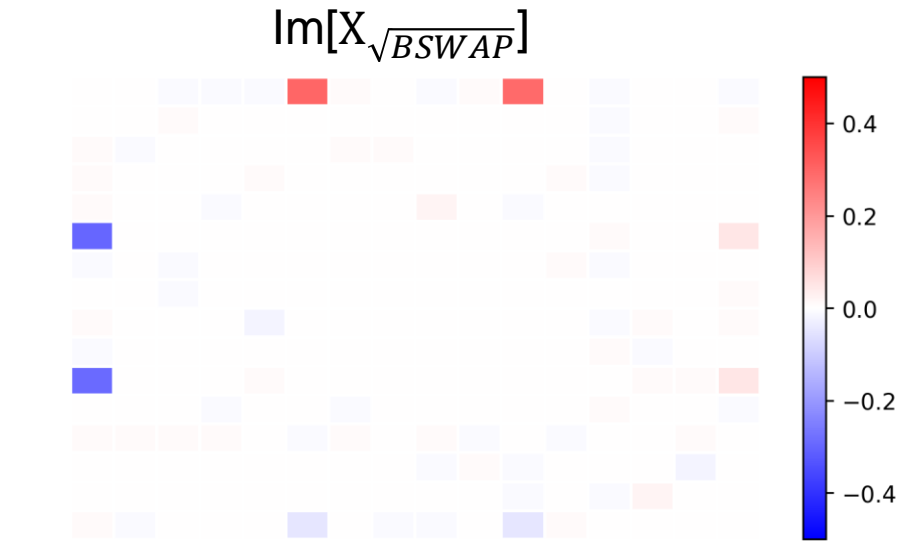
$$\omega_d = \omega_B + \omega_A$$

\sqrt{bSWAP} gate length = 101ns

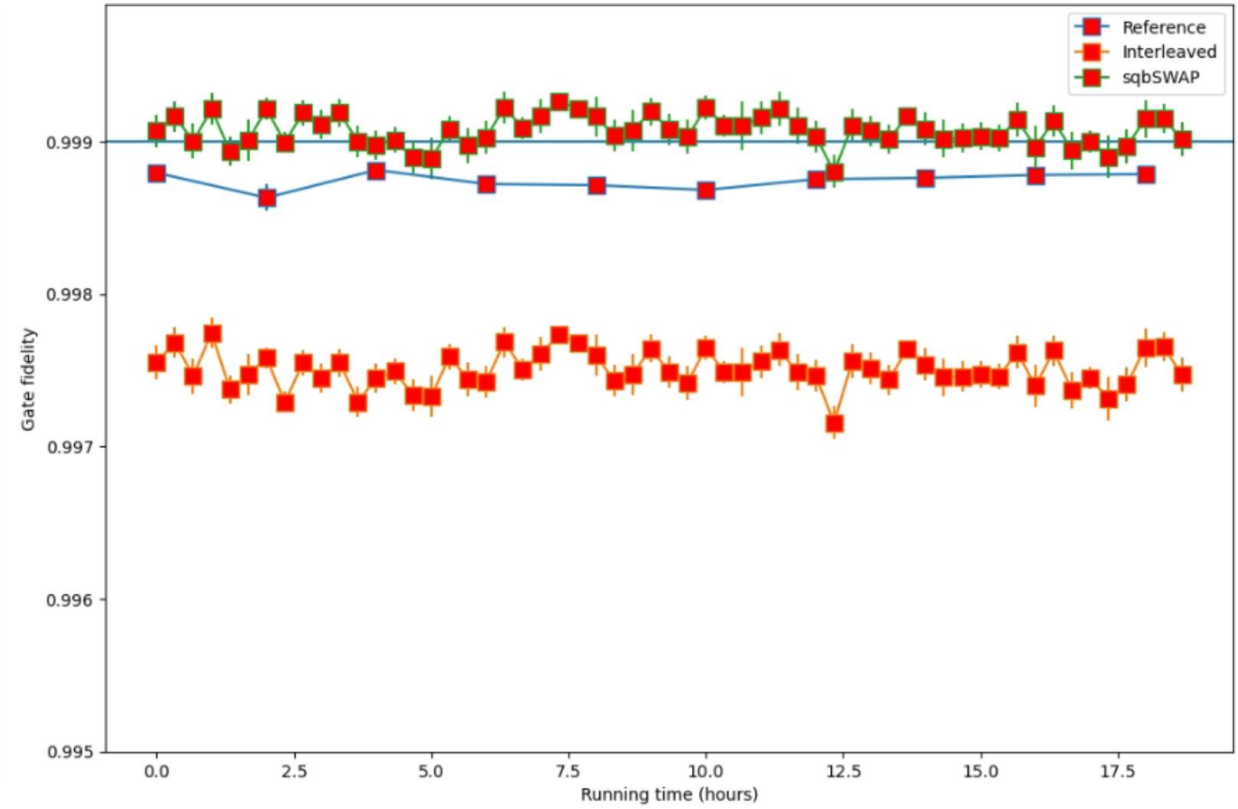
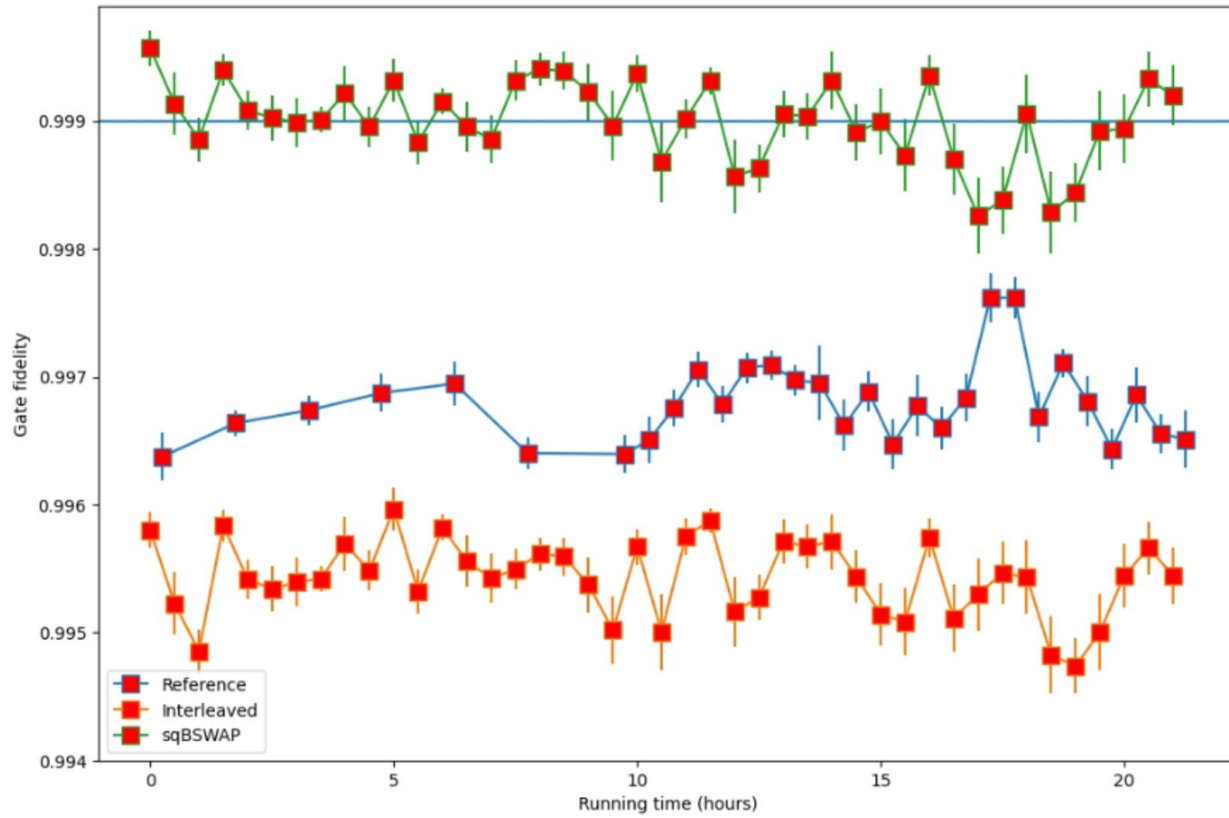
sqiSWAP and sqbSWAP Kraus matrix



SQBSWAP Process Matrix



XEB results - sqbswap



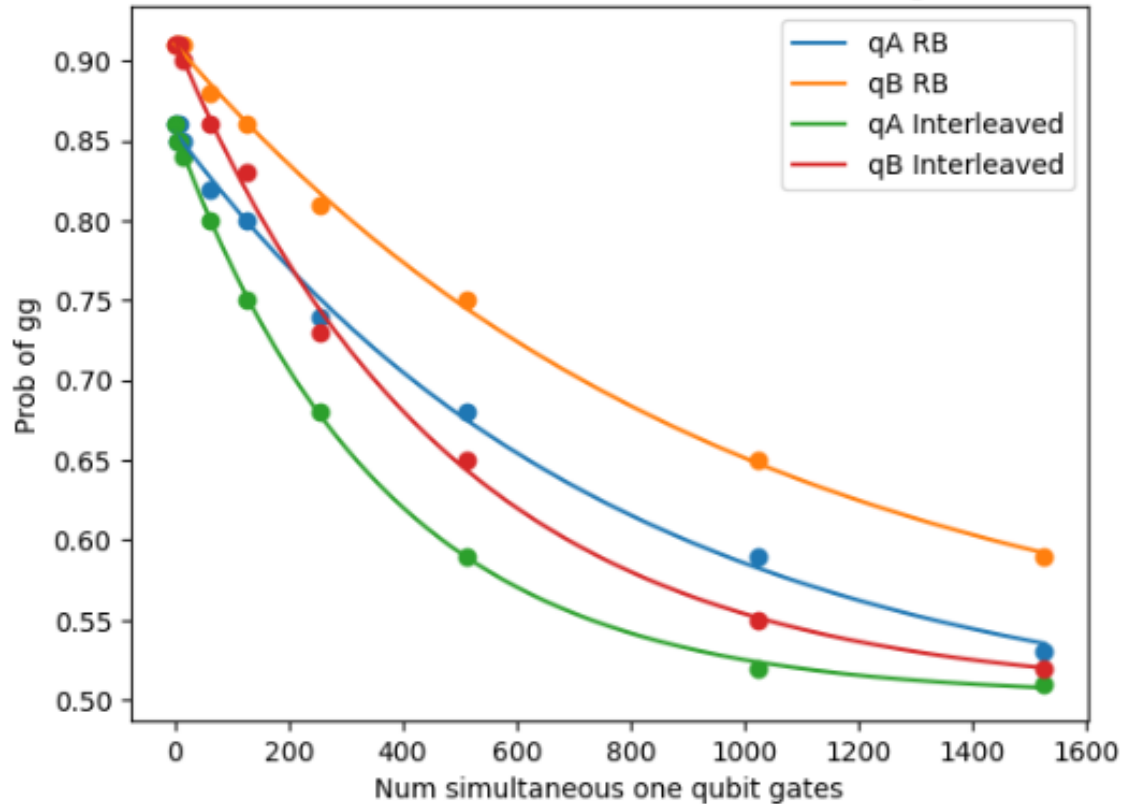
Interleaved Randomized Benchmarking



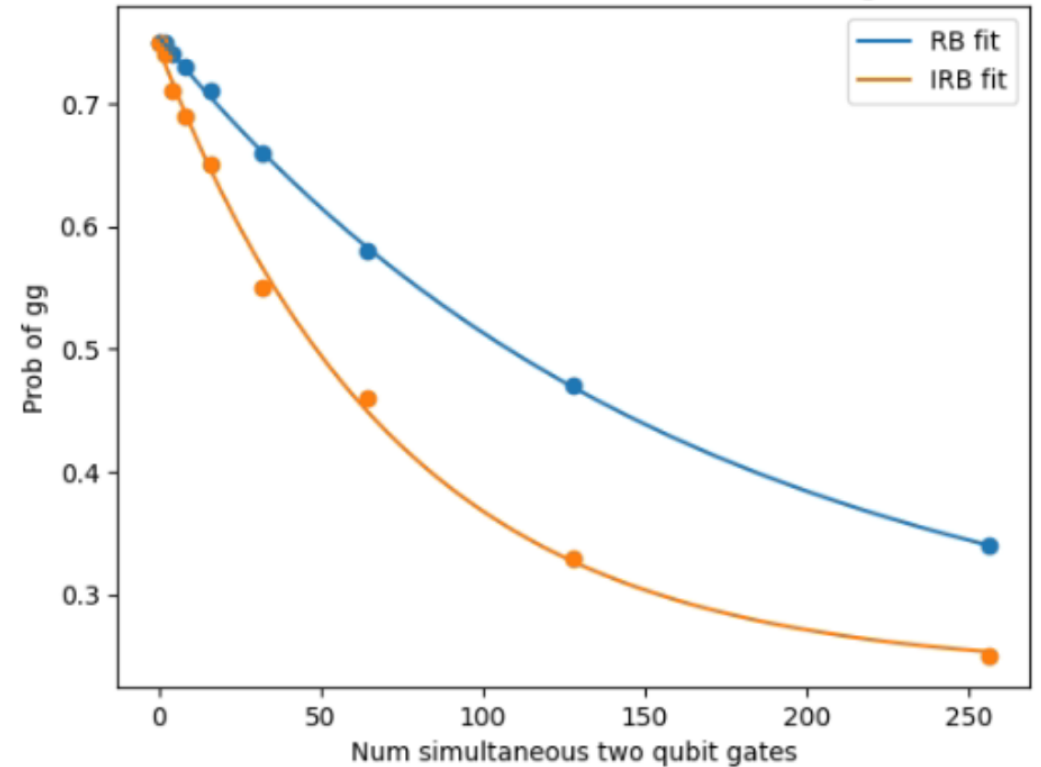
0.065% +/- 0.006% error per qA 1Q Clifford
0.052% +/- 0.004% error per qB 1Q Clifford
0.137% +/- 0.003% error per 1Q Clifford (Interleaved qA X/2)
0.104% +/- 0.006% error per 1Q Clifford (Interleaved qB X/2)
0.072% error per qA X2 gate (IRB extracted)
0.052% error per qB X2 gate (IRB extracted)

99.533 +/- 0.02% fidelity per average 2Q Clifford
98.986 +/- 0.064% fidelity per 2Q Clifford (Interleaved CNOT)
99.45% fidelity per CNOT gate (IRB)

1Q Cliffords Randomized Benchmarking



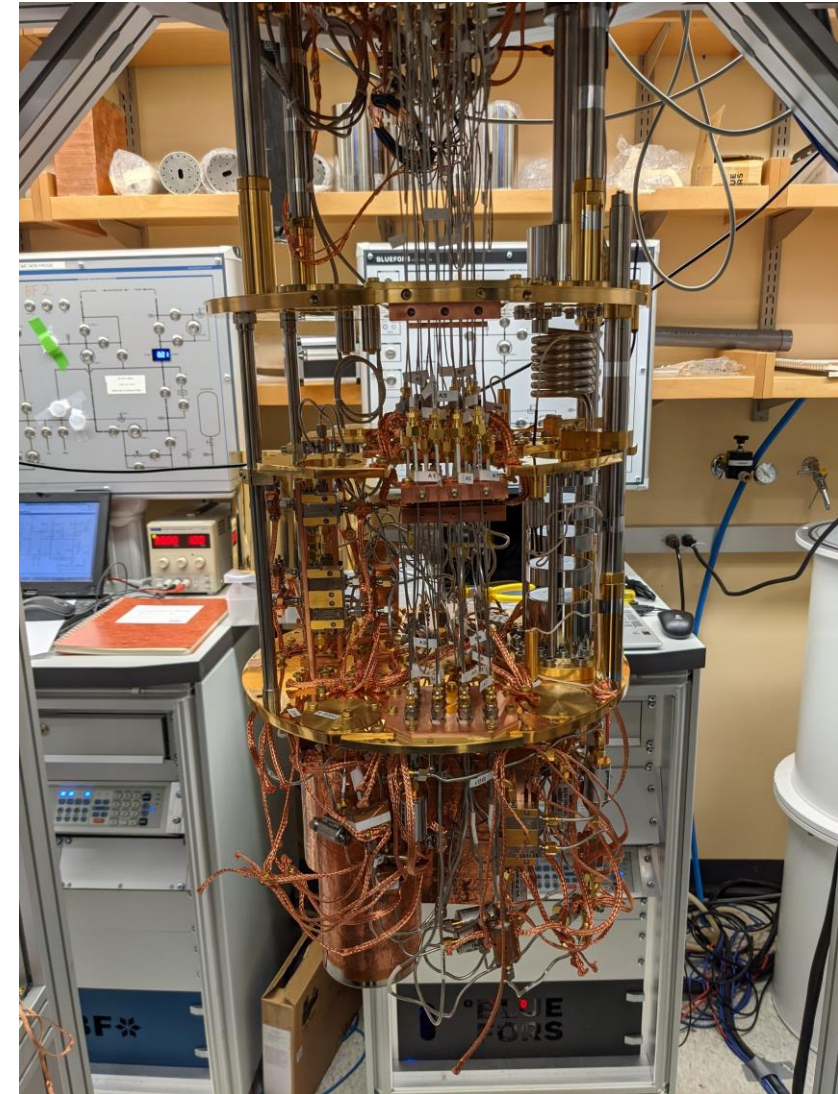
2Q Cliffords Randomized Benchmarking



Major challenges?



- Very sensitive to flux bias noise (sweet spot $<1e-4 \Phi_0$), and thus requires a very clean fridge ground
- Crosstalk between gates is very significant, requires the use of cancellation pulses, careful tuning procedures, and prevents gates from being played in parallel
- Scaling the galvanically coupled design might be difficult for large 2D array of qubits



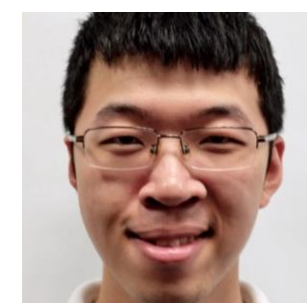
Thank you!



Yuwei Ma



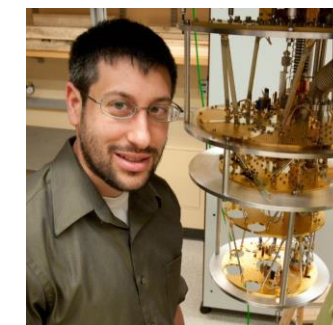
Daniel Weiss



Helin Zhang



Sai Paivan Chitta

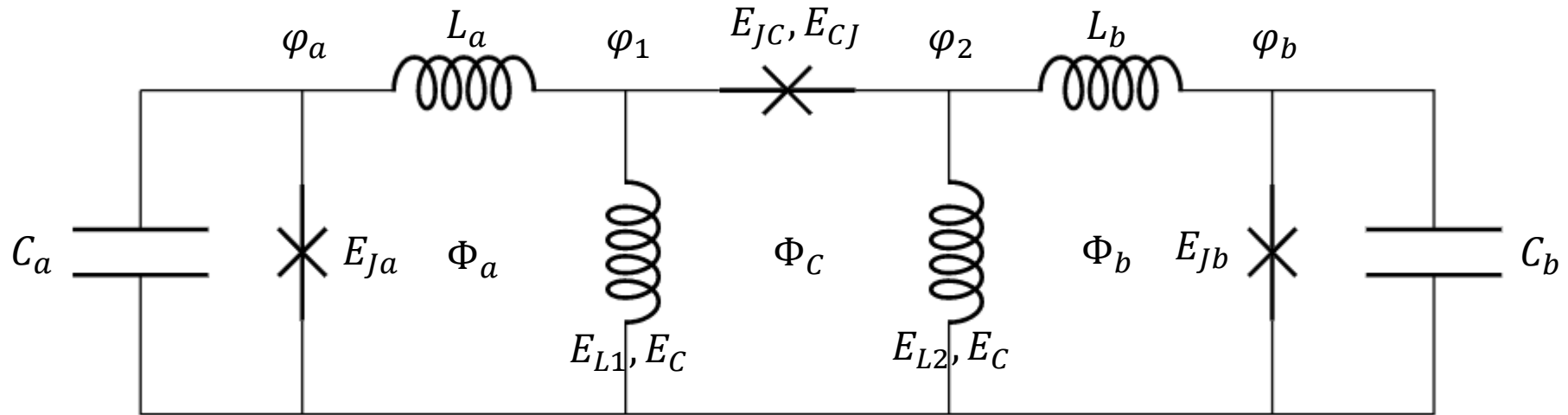


PI: David Schuster



PI: Jens Koch

Current sample parameters



QubitA

$$E_J \approx 5.837 \text{ GHz}$$

$$E_C \approx 0.892 \text{ GHz}$$

$$E_L \approx 0.271 \text{ GHz}$$

$$E_r \approx 7.241 \text{ GHz}$$

Coupler

$$E_J \sim 4.246 \text{ GHz}$$

$$E_C \sim 11 \text{ GHz}$$

$$E_L \sim 3.52 \text{ GHz}$$

QubitB

$$E_J \approx 4.930 \text{ GHz}$$

$$E_C \approx 0.8655 \text{ GHz}$$

$$E_L \approx 0.266 \text{ GHz}$$

$$E_r \approx 6.9857 \text{ GHz}$$

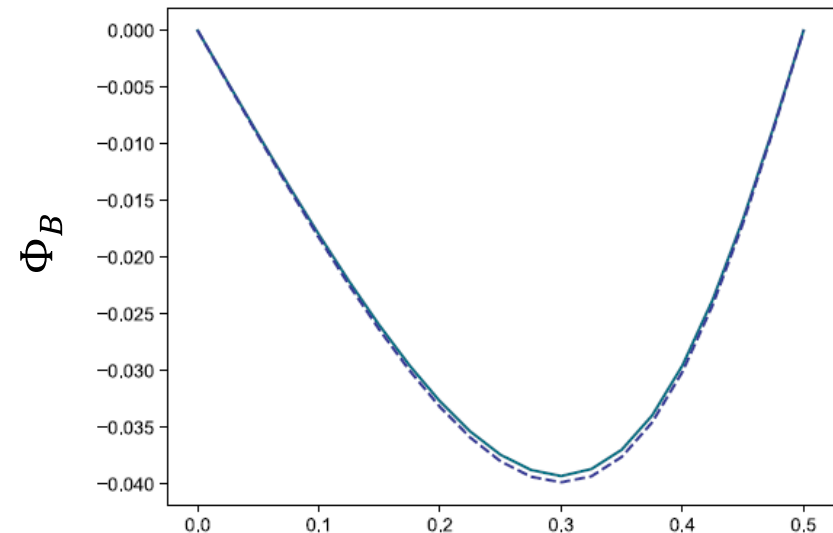
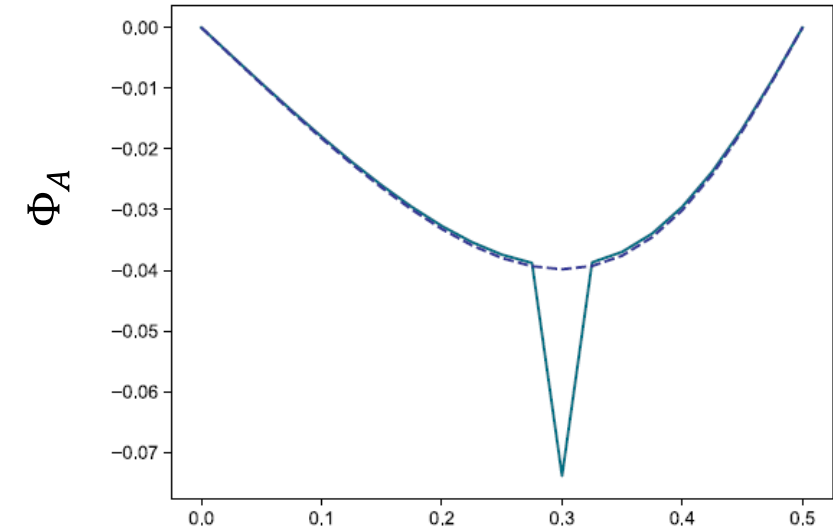
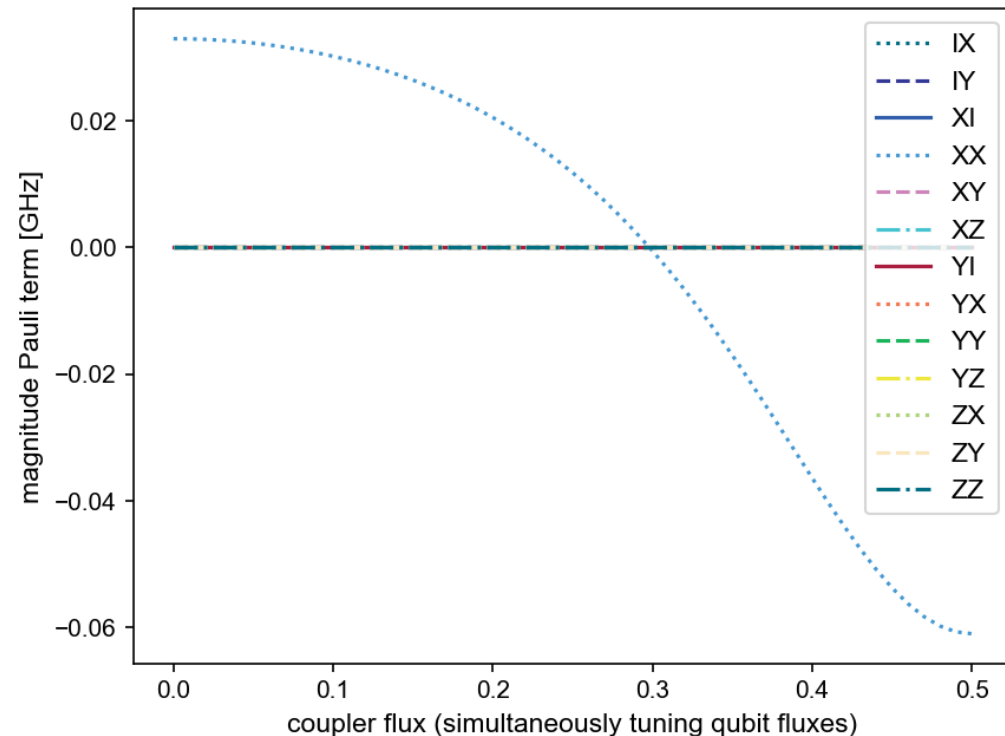
$$\omega_{qa} \approx 48 \text{ MHz}$$

$$\omega_{qb} \approx 61 \text{ MHz}$$

Tunable coupler performance

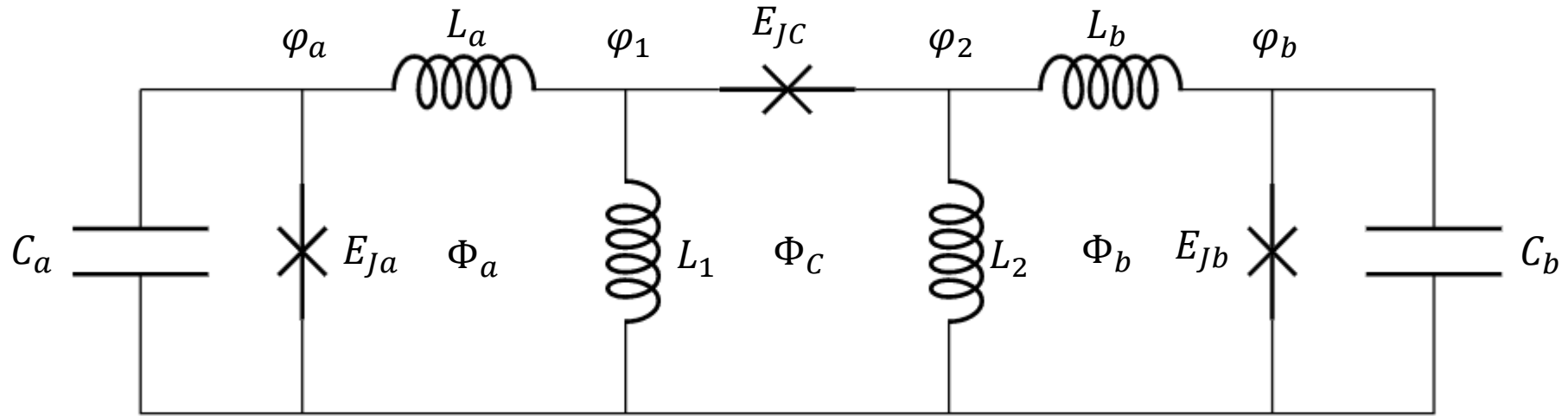


- Tunable from 0 – 60MHz using coupler flux, results in gate speeds of 20ns
- Requires adjustments to single qubit flux to remain on sweet spot



$\Phi_{Coupler}$

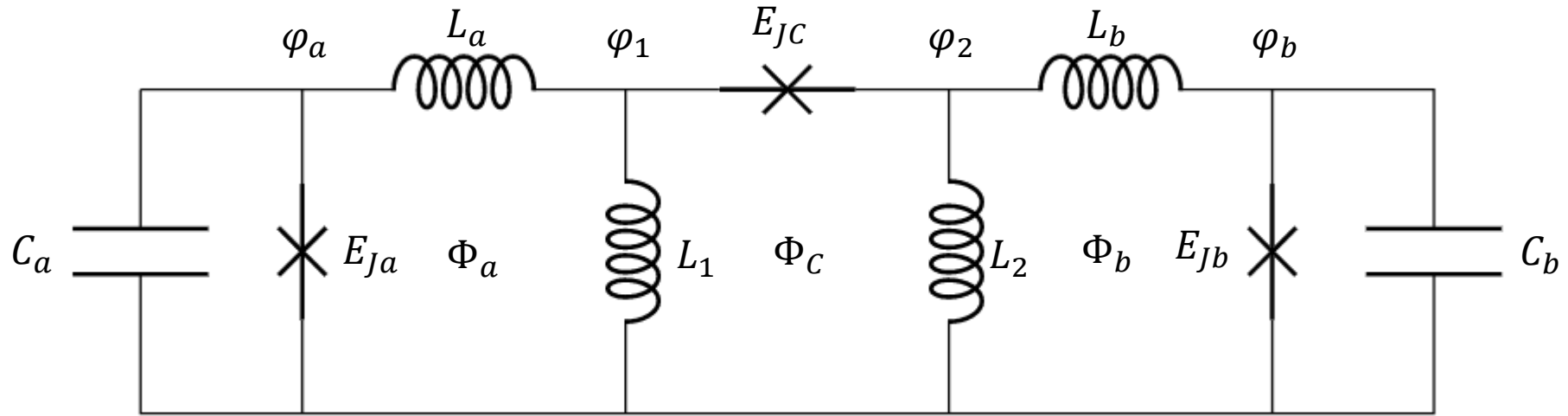
Tunable Coupler – Circuit Analysis



$$\varphi_+ = (\varphi_1 + \varphi_2)/2$$

$$\varphi_- = \varphi_1 - \varphi_2$$

Tunable Coupler – Circuit Analysis



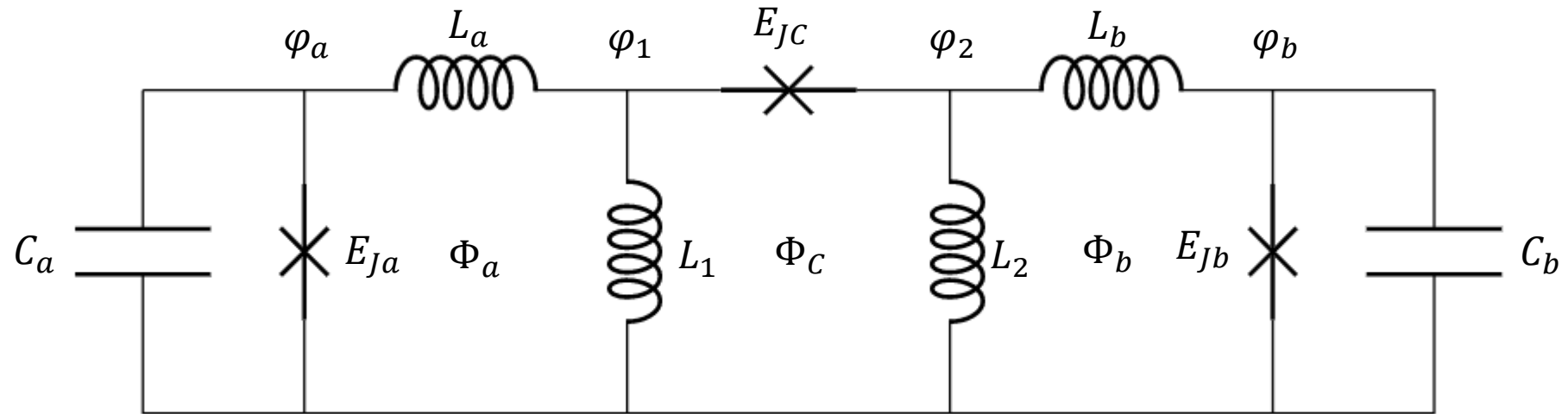
$$H = 4E_{C_a} n_a^2 + 4E_{C_b} n_b^2 + \frac{E_{L_a}}{2} \varphi_a^2 + \frac{E_{L_b}}{2} \varphi_b^2 - E_{J_a} \cos(\varphi_a + \varphi_{e,a}) - E_{J_b} \cos(\varphi_b + \varphi_{e,b})$$

$H_a + H_b$

$$H_+ + H_- = 2E_C n_+^2 + \frac{E_{\tilde{C}} n_-^2}{2} + \frac{\tilde{E}_{L_1} + \tilde{E}_{L_2}}{2} (\varphi_+^2 + \frac{\varphi_-^2}{4}) - E_J \cos(\varphi_- + \varphi_e) + (\tilde{E}_{L_1} - \tilde{E}_{L_2}) \varphi_+ \varphi_- / 2 - E_{L_a} \varphi_a (\varphi_+ + \varphi_- / 2) - E_{L_b} \varphi_b (\varphi_+ - \varphi_- / 2).$$

H'

Tunable Coupler – Circuit Analysis



Calculate H' with perturbation theory

First order:
$$\varphi_a \rightarrow \varphi_a + \frac{\langle \psi_0^{(0)} | \varphi_- | \psi_0^{(0)} \rangle}{2}, \quad \varphi_b \rightarrow \varphi_b - \frac{\langle \psi_0^{(0)} | \varphi_- | \psi_0^{(0)} \rangle}{2}$$

Second order gives us a coupling term:
$$E_{La} E_{Lb} (\chi^- / 2 - 2\chi^+) \varphi_a \varphi_b$$

where
$$\chi^- \equiv \sum_{m \neq 0} \frac{|\langle \psi_0^- | \varphi_- | \psi_m^- \rangle|^2}{e_m^- - e_0^-} \quad \chi^+ \equiv \frac{|\langle \psi_0^+ | \varphi_- | \psi_1^+ \rangle|^2}{\omega_+}$$

XEB results - sqiswap

