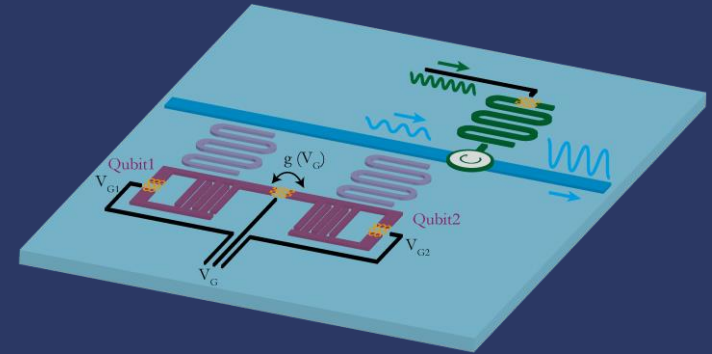
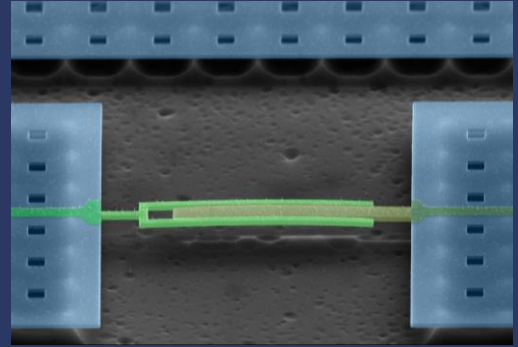


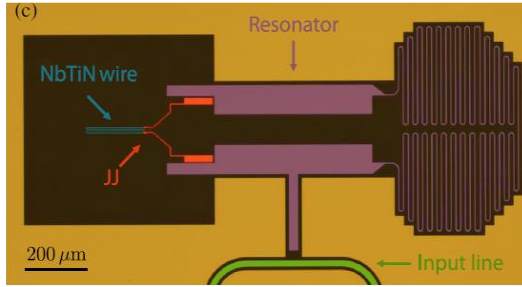
TiN high impedance superconducting circuits

(+ second part about graphene
superconducting circuits)

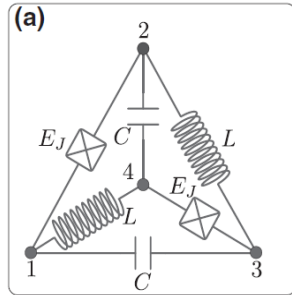
Julien Renard



Protected Qubits

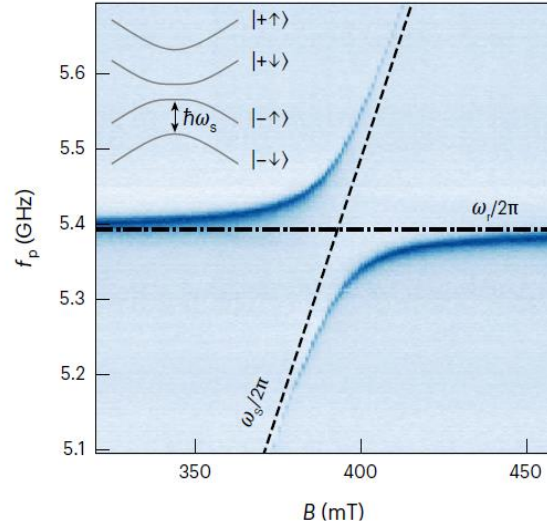


T.M. Hazard et al Phys. Rev. Lett. 2019
 V. Manucharyan et Science 2009



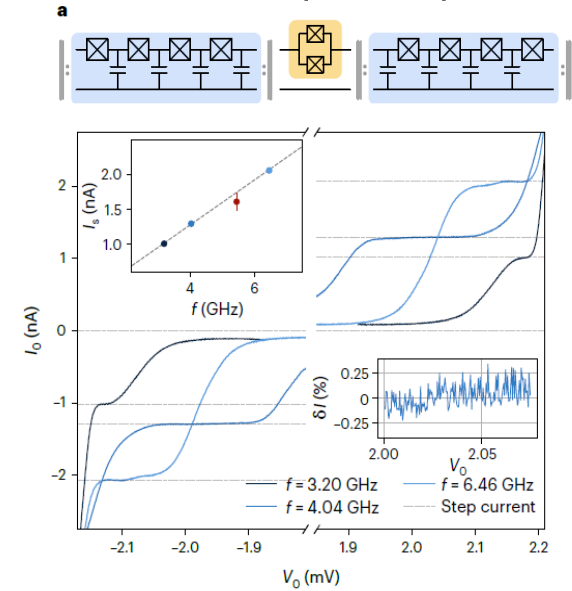
A. Gyenis et al Phys. Rev. Quantum 2021

Spin-photon coupling



C. Yu et al Nature Nano 2023

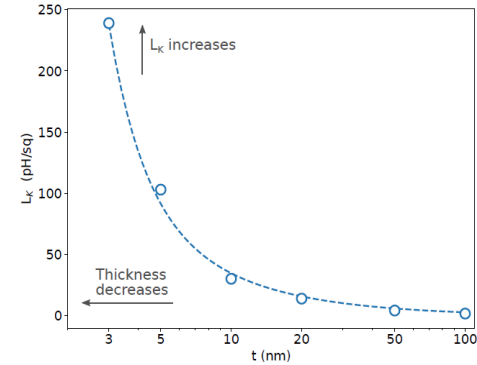
Dual Shapiro steps



N. Crescini et al Nature Phys 2023
 R. S. Shaikhaidarov et al Nature 2022

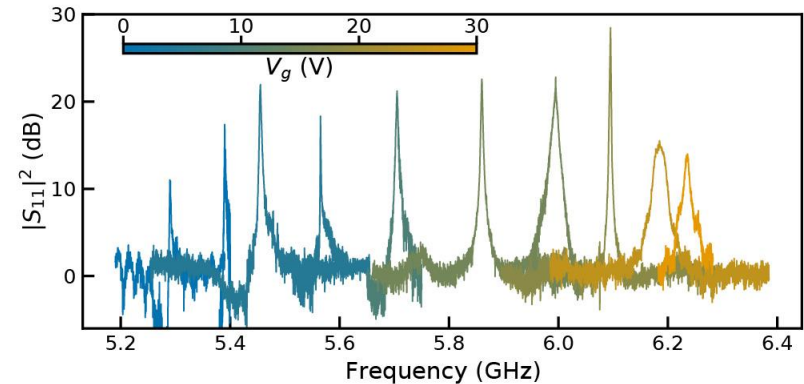
TiN high impedance superconducting circuits

K. Rafsanjani Amin et al Appl. Phys. Lett (2022)



Graphene based superconducting quantum circuits

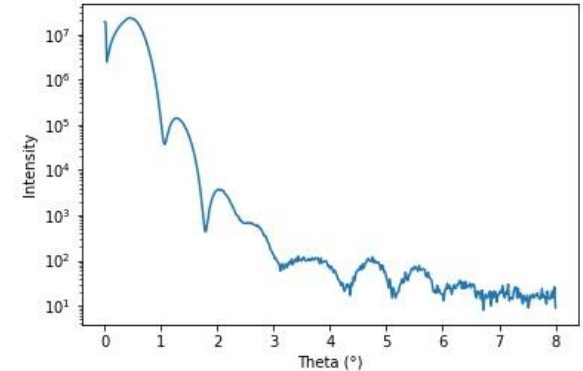
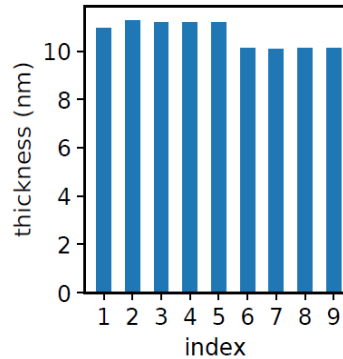
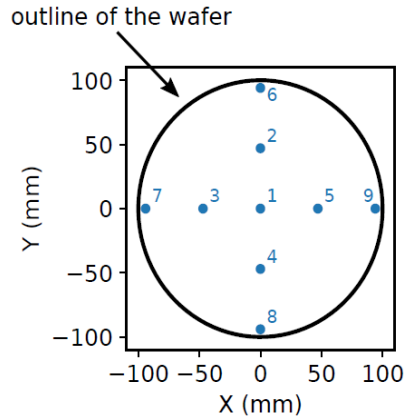
G. Butseraen et al Nature Nano (2022)

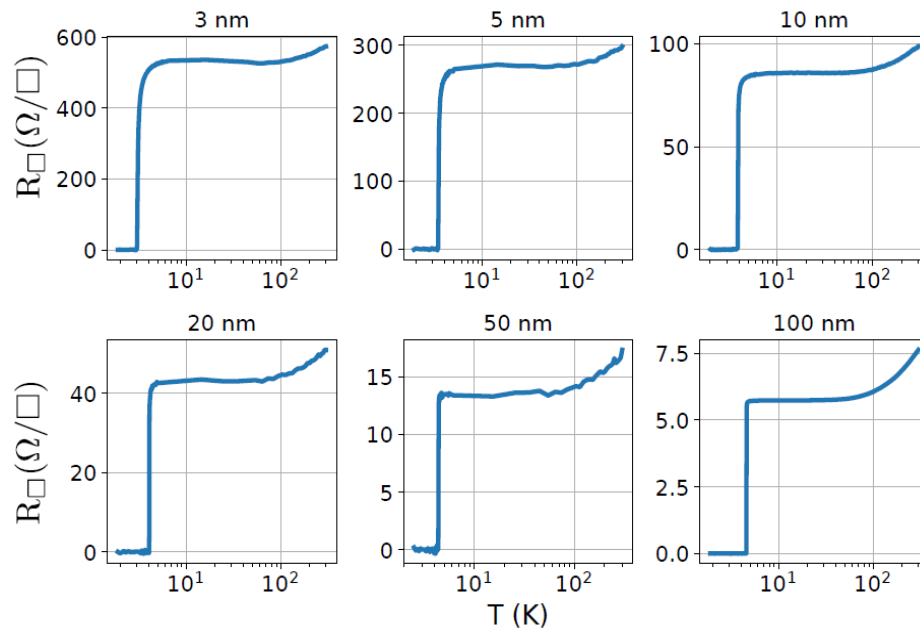


Sputtering @ 350°C on highly resistive Si (200 mm wafers)

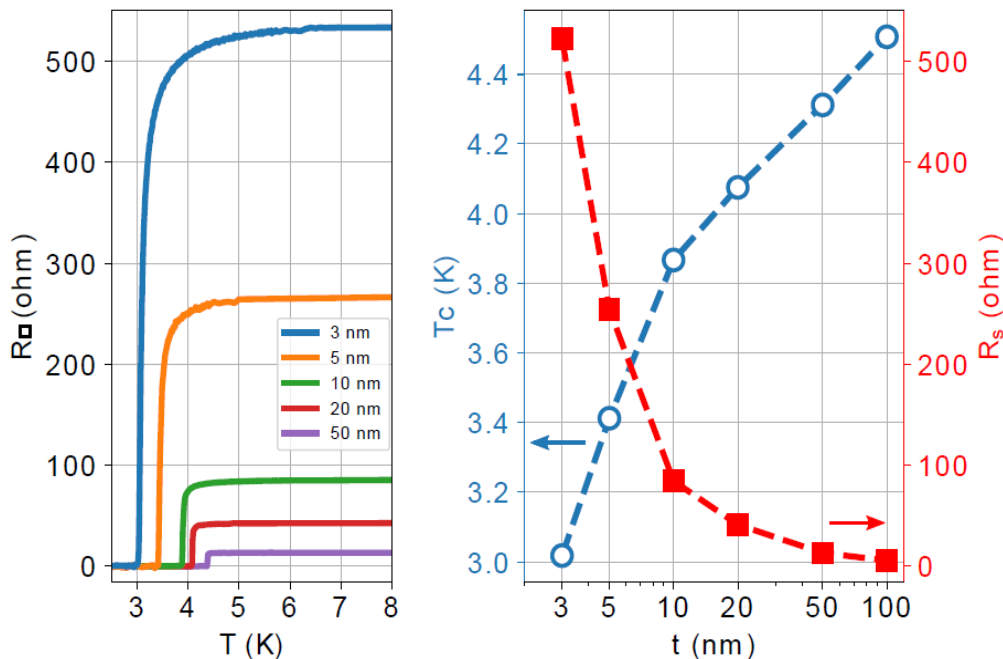
Structural characterization: ellipsometry, XRR, XRD

→ excellent thickness uniformity

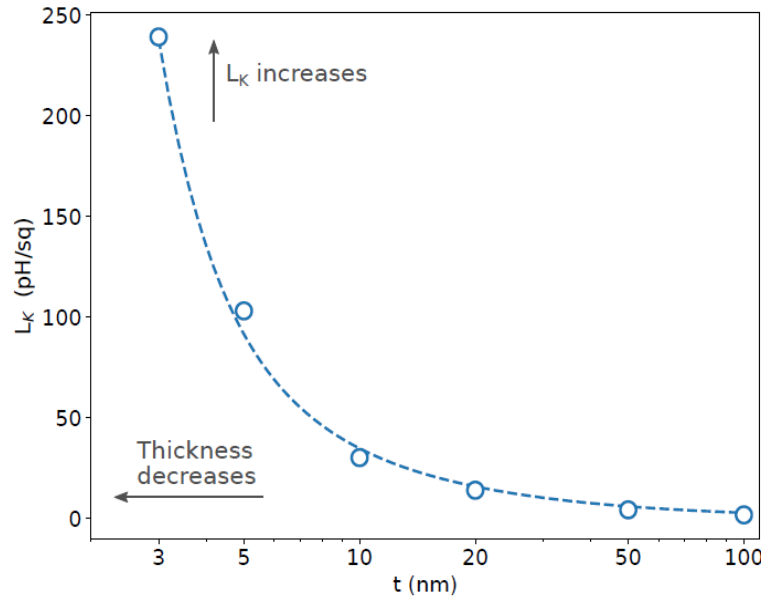




Superconductivity down to 3 nm thickness



Low disorder in ultrathin TiN films (usually closer to SIT @3nm)

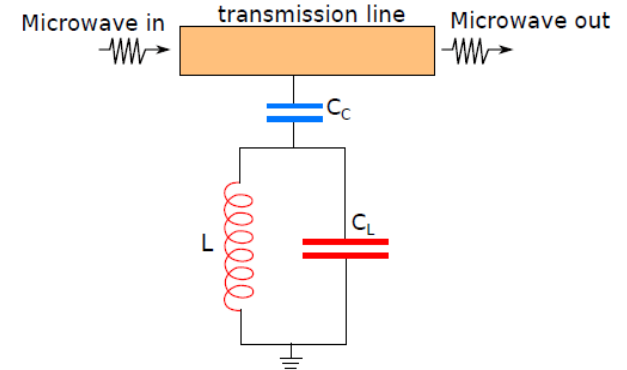
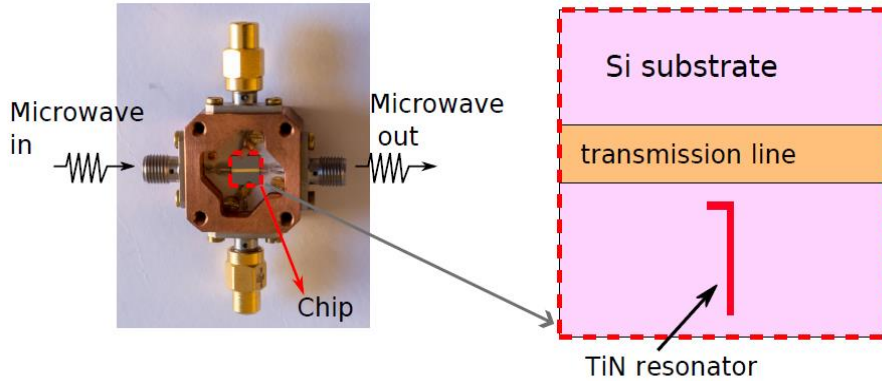


Use BCS formula: $L_K = \frac{\hbar R_{\square}}{\pi \Delta_0}$

Thickness (nm)	RRR	T_C (K)	R_S (Ω/\square)	L_K (pH/□)
3	1.10	3.0	522.0	239.0
5	1.17	3.4	254.4	103.0
10	1.17	3.9	93.5	33.4
20	1.23	4.1	41.2	14.0
50	1.25	4.3	13.3	4.2
100	1.33	4.5	5.7	1.7

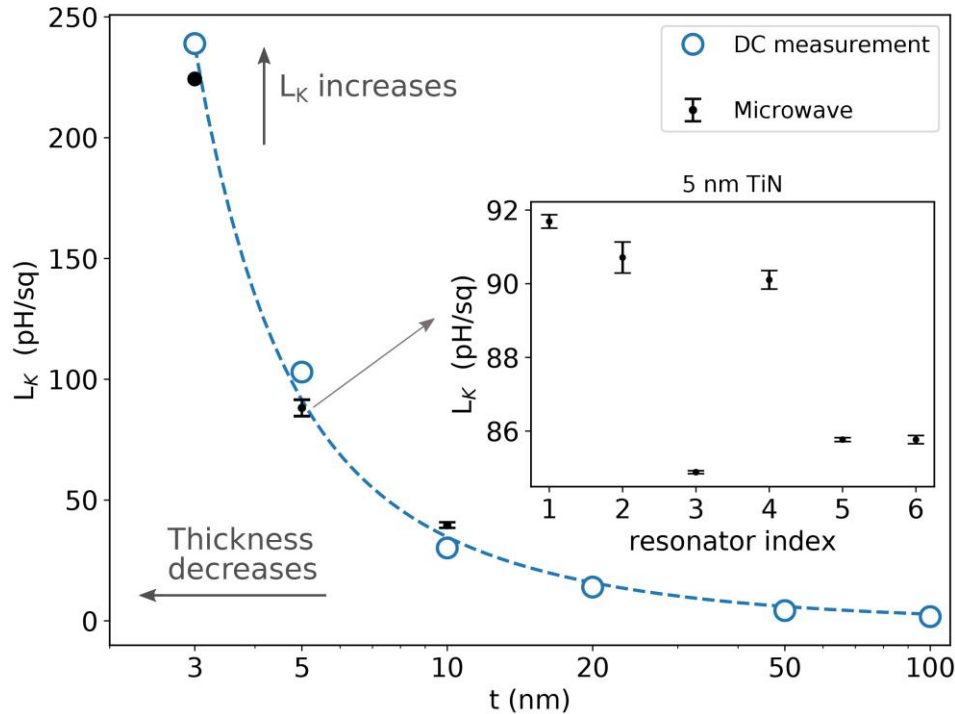
Similar values to Shearow et al APL 2018 (but much thinner films in our case)

Up to ~ 240 pH/□ for 3nm thick film

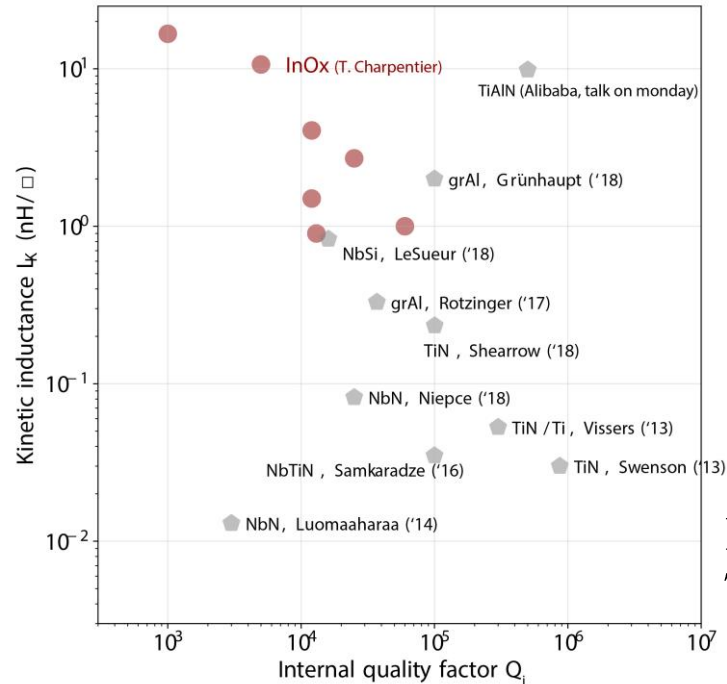


$$\omega = \frac{1}{\sqrt{LC}}, \text{ with } L \sim L_K$$

Directly determine L_K from resonance frequency

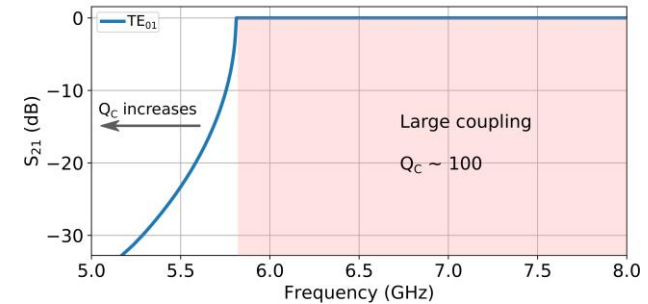
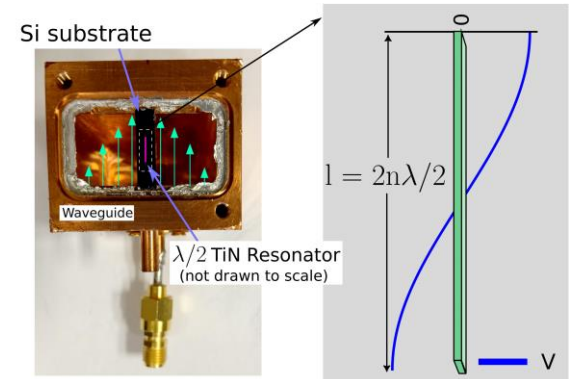
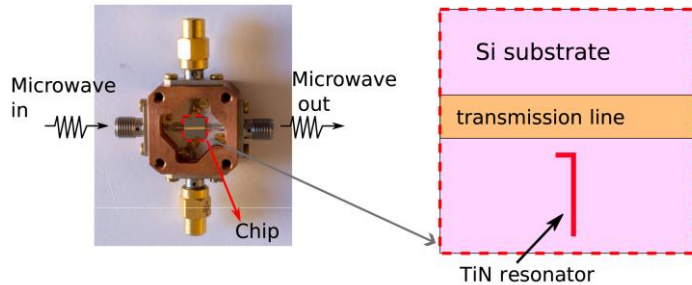
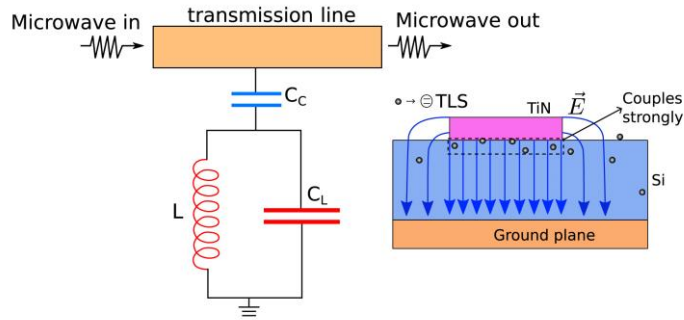


DC and microwave L_K fully consistent and wafer scale homogeneity

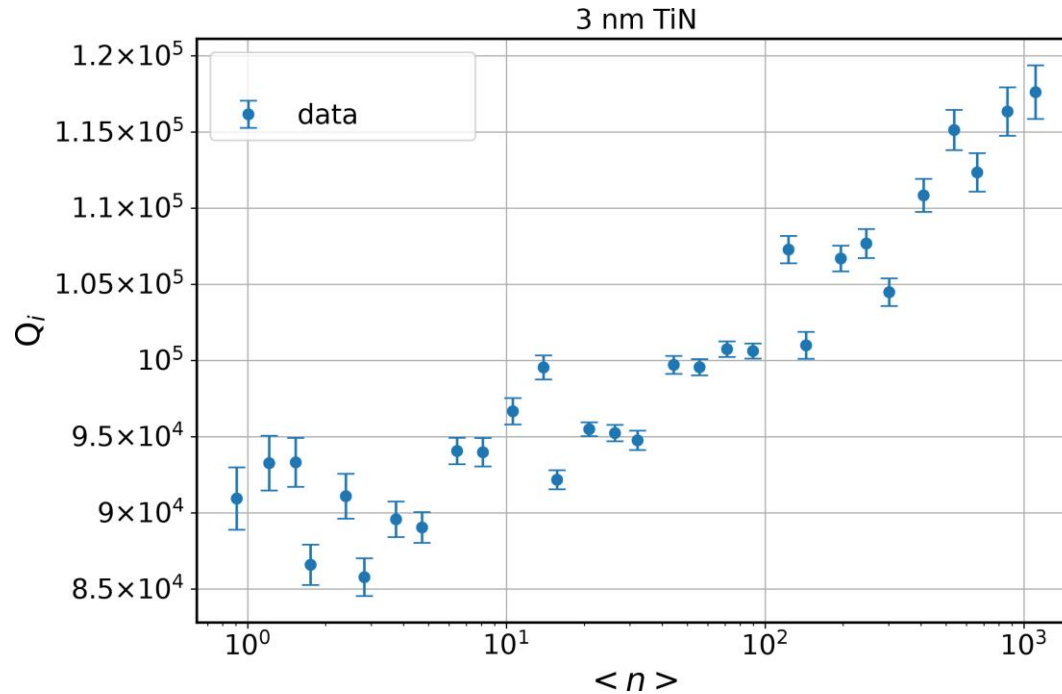


Plot courtesy of
T. Charpentier

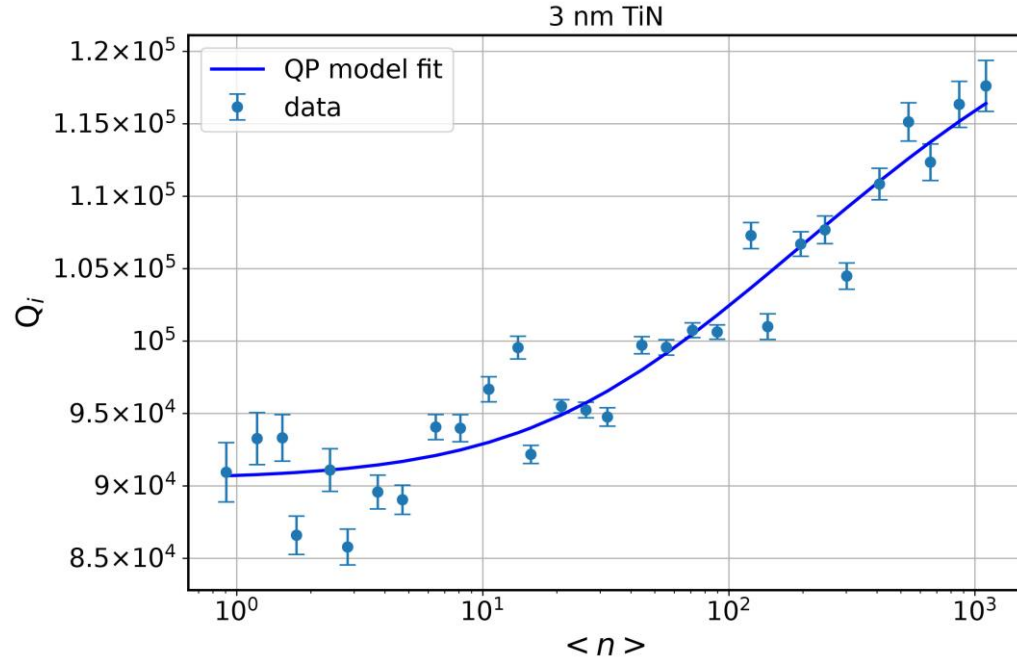
→ Study of losses (and mechanisms) in TiN thin films



→ Study of losses (and mechanisms) in TiN thin films



→ Modest increase with photon number



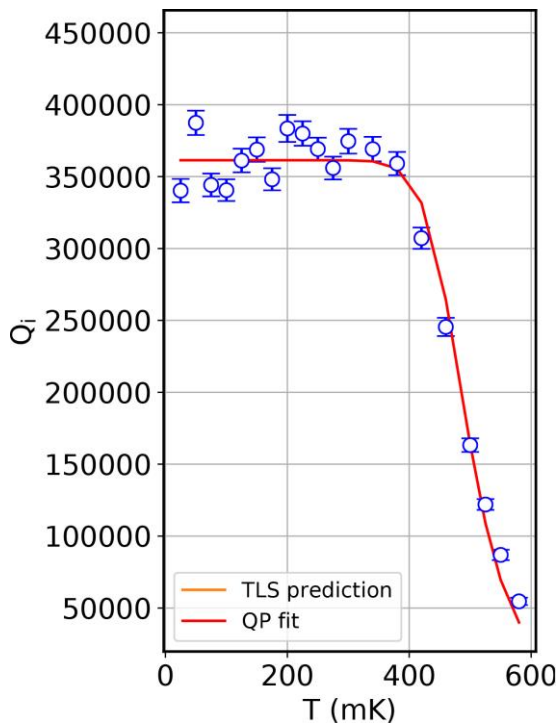
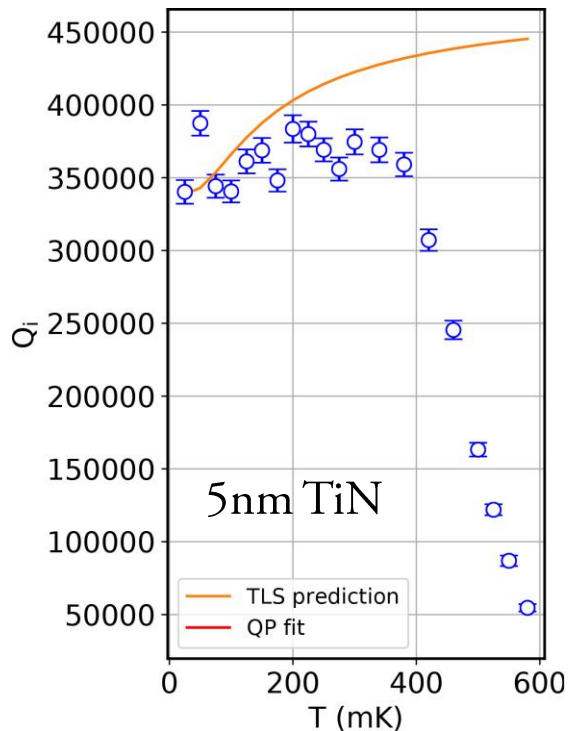
$$\frac{1}{Q_i} = \frac{1}{Q_0} + \beta \left(\frac{1}{1 + \frac{1}{1 + \frac{1}{2}(\sqrt{1+4\gamma n} - 1)}} - 1 \right)$$

Follow L. Grünhaupt et al, Phys Rev Lett 2018

Fitting parameters:

β , γ (linked to QP dynamics and interaction with photons)

Q_i possibly limited by quasiparticles

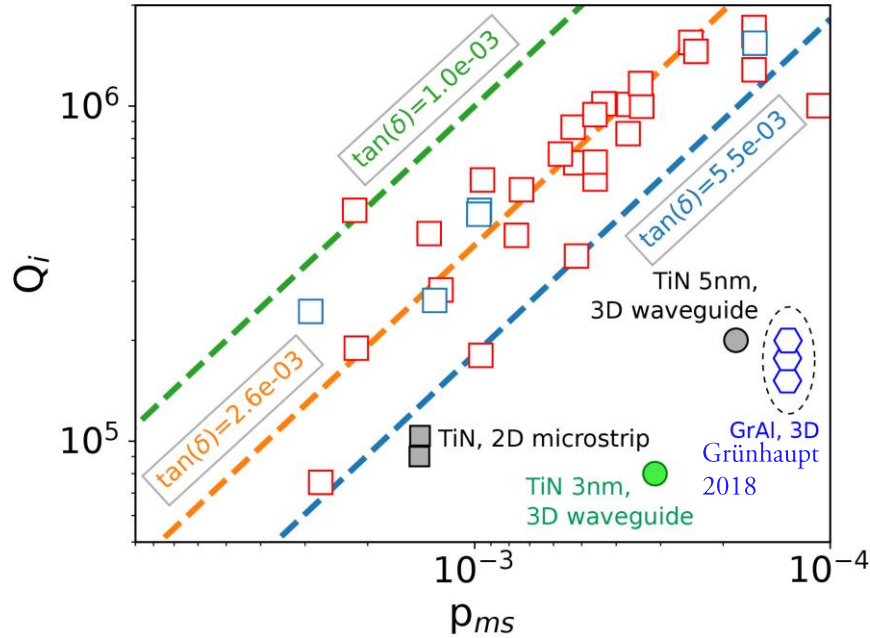


Thermally induced QP

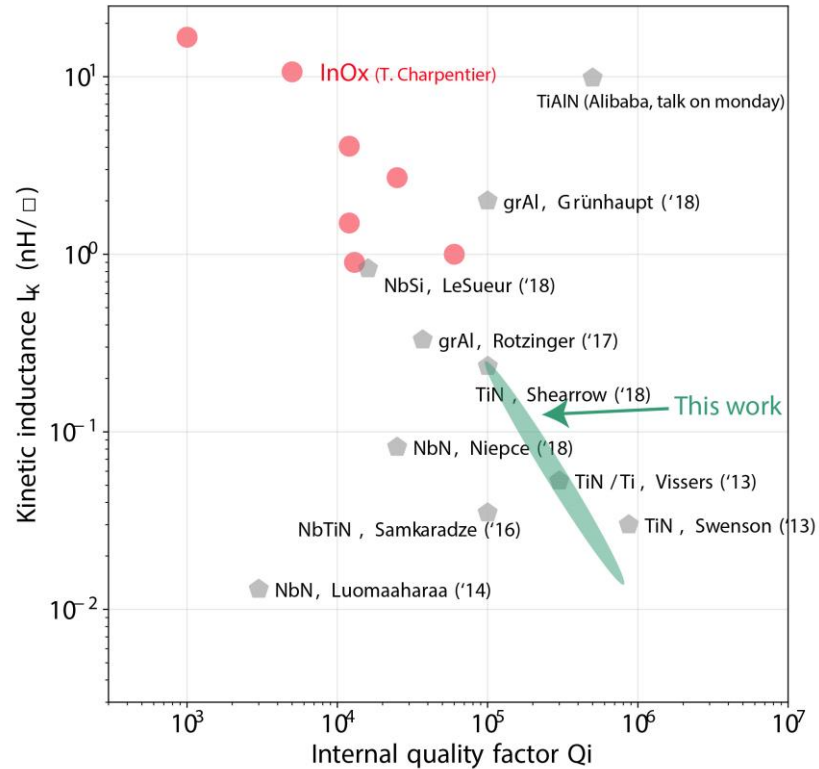
$$\frac{1}{Q_i} = 2\alpha \sqrt{\frac{k_B T}{\pi h f_R}} \exp\left(-\frac{\Delta}{k_B T}\right) + \frac{1}{Q_a}$$

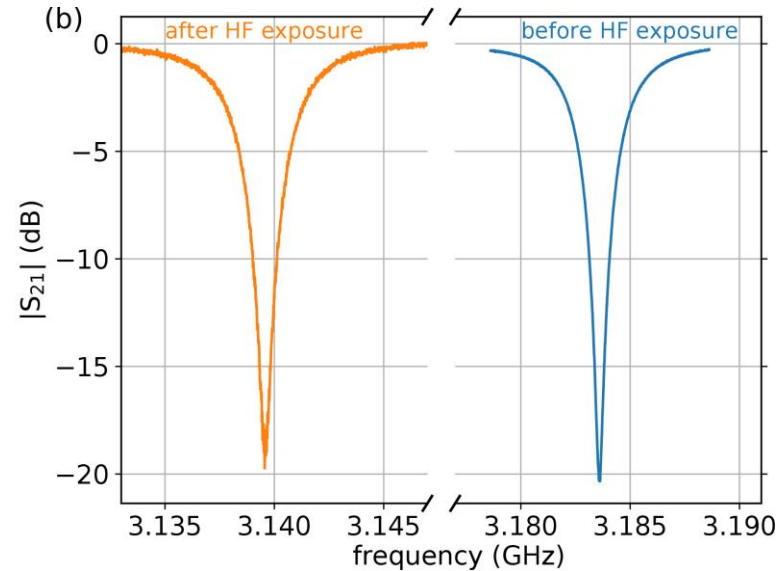
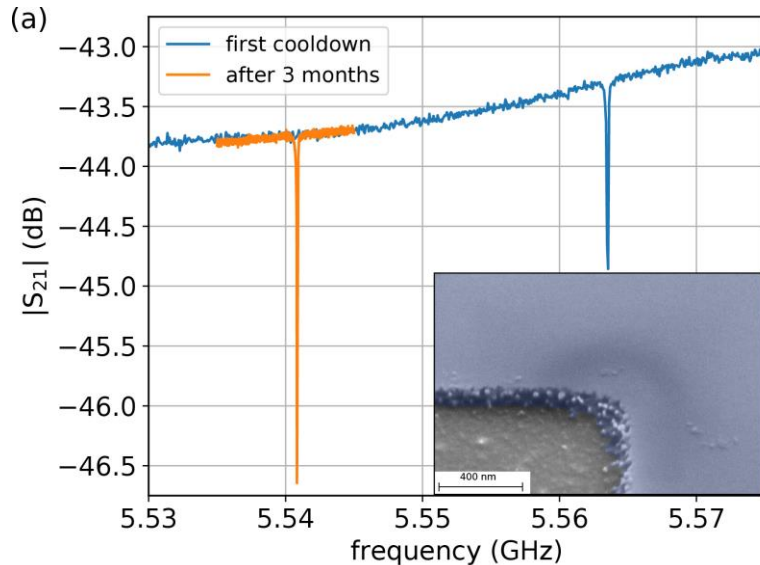
α : kinetic inductance fraction
 Q_a : other losses than thermally induced QP

No sign of TLS limited internal Q



Importance of QP in high L_K materials

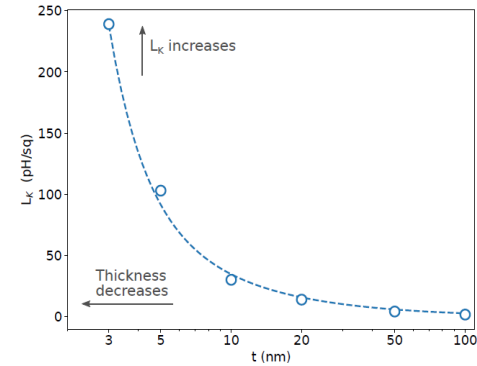




Very minimal aging (3nm thin film!) and resilient to HF

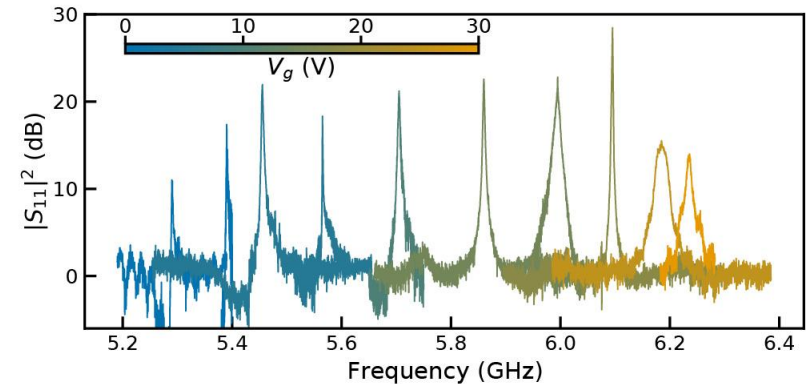
TiN high impedance superconducting circuits

K. Rafsanjani Amin et al Appl. Phys. Lett (2022)

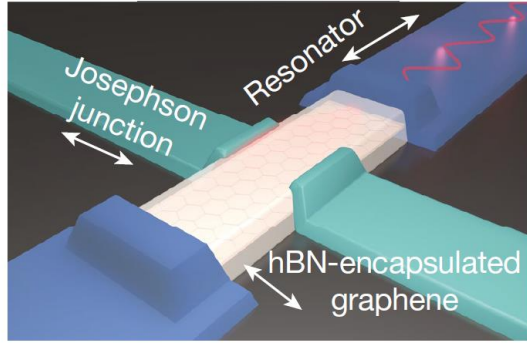


Graphene based superconducting quantum circuits

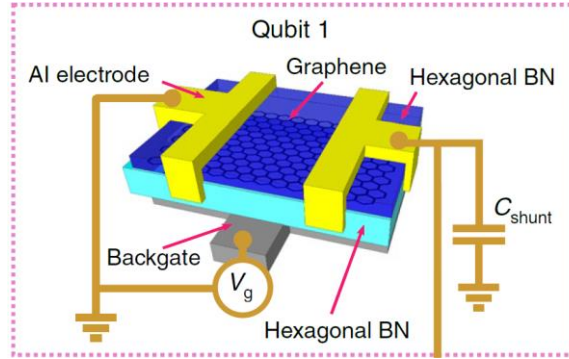
G. Butseraen et al Nature Nano (2022)



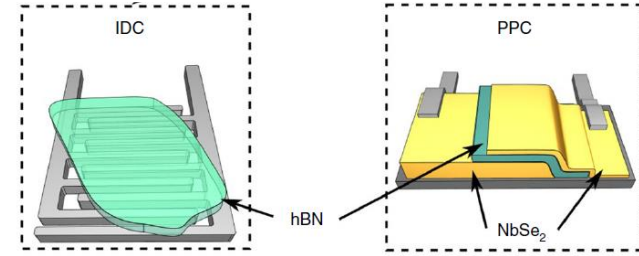
Microwave bolometers



Quantum bit



h-BN (low losses)



G.H. Lee et al Nature (2020)

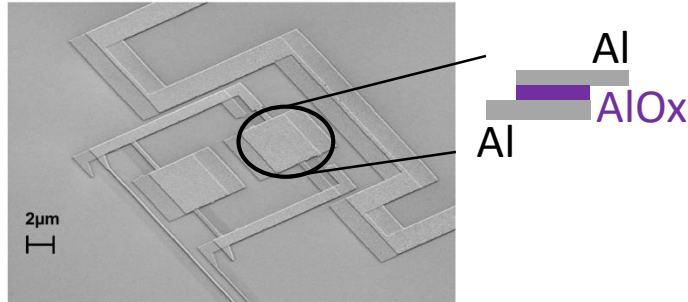
J.I.J. Wang et al Nature Nano (2019)

J.I.J. Wang et al Nature Mater (2022)

R. Kokkonen et al Nature (2020)

Large variety of applications in superconducting quantum technologies

« Standard » tunnel JJ

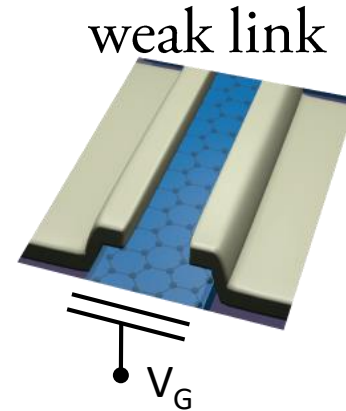


F. Lecocq PhD thesis, I. Néel

- Large number of channels with low transmission
- Josephson energy tunable with a magnetic field

$$E_J(\phi)$$

Graphene (semiconductor)

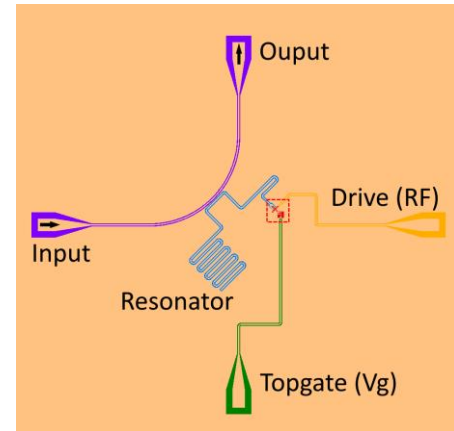
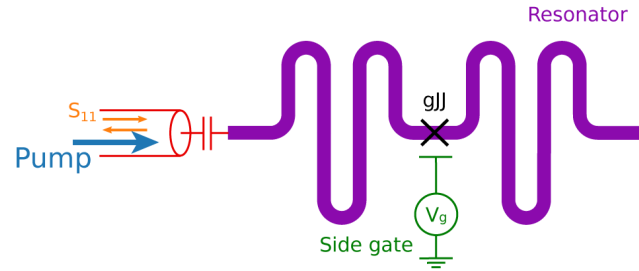


- Small number of channels with high transmission
- Josephson energy tunable with an electric field

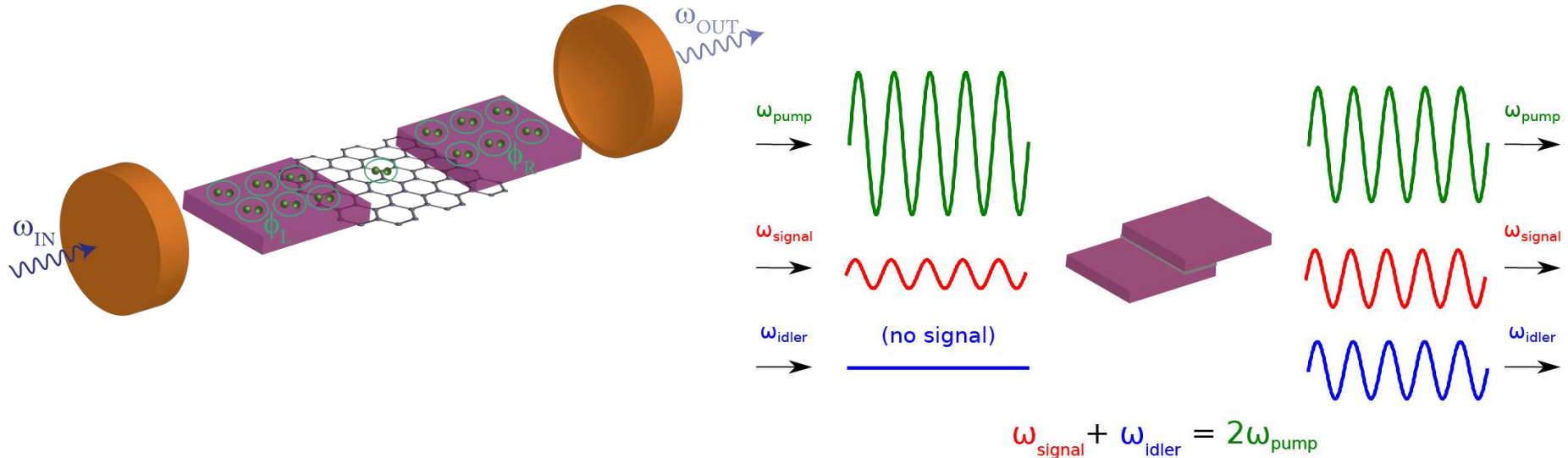
$$E_J(V_G)$$

Gate-tunable Josephson parametric amplifier

Qubit with a gate-tunable anharmonicity

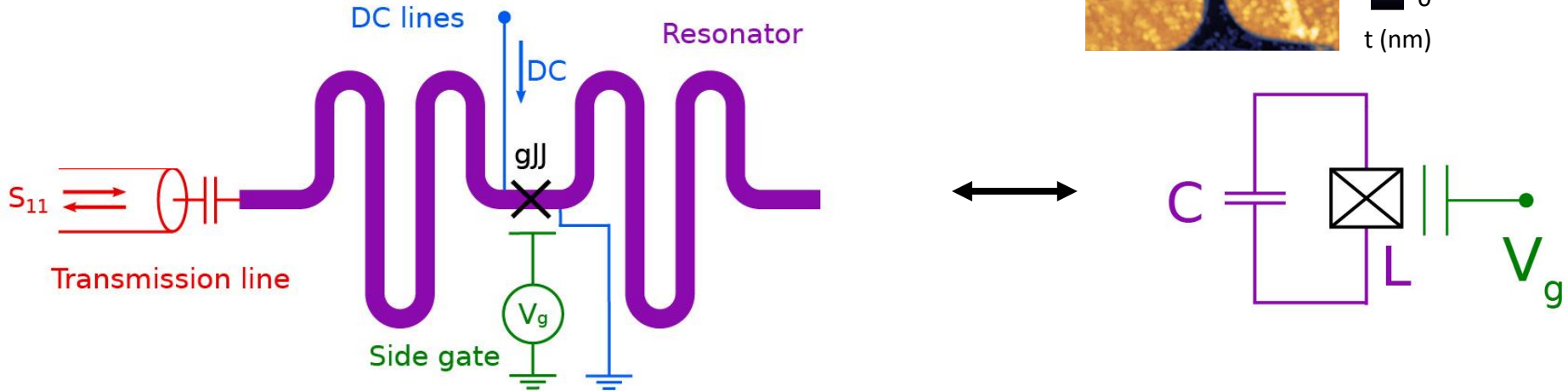
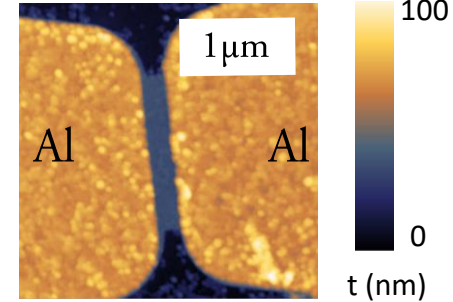
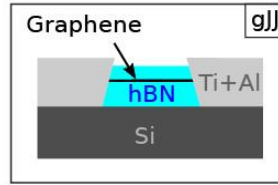
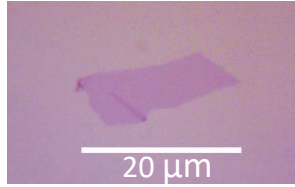


(Resonant) parametric amplifier

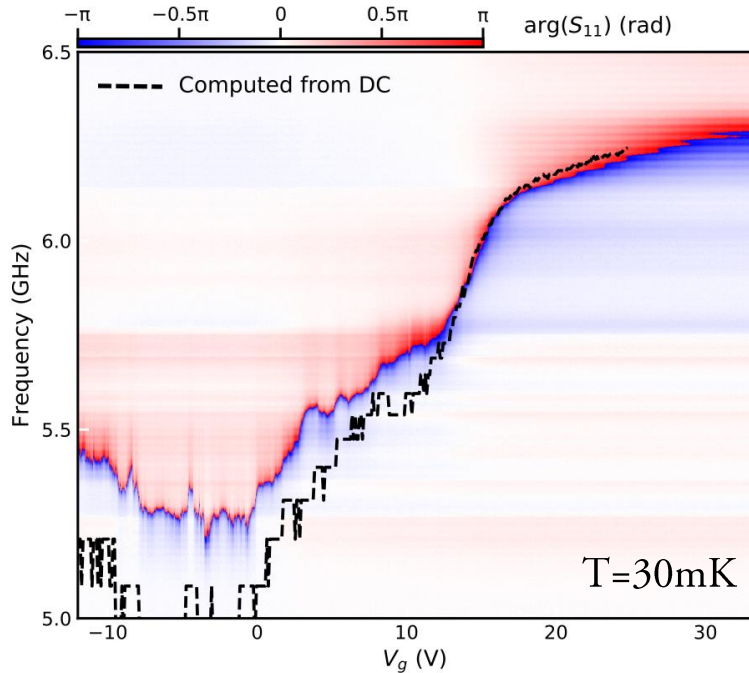


$$H = \hbar\omega a^\dagger a + \hbar g a^\dagger a^\dagger a a$$

Non-linearity introduced by the Josephson junction

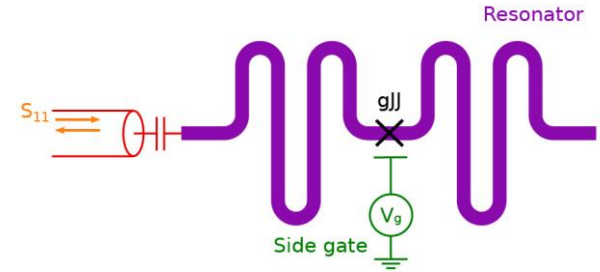
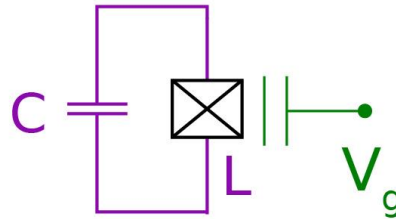


Microwave resonator with an embedded graphene Josephson junction

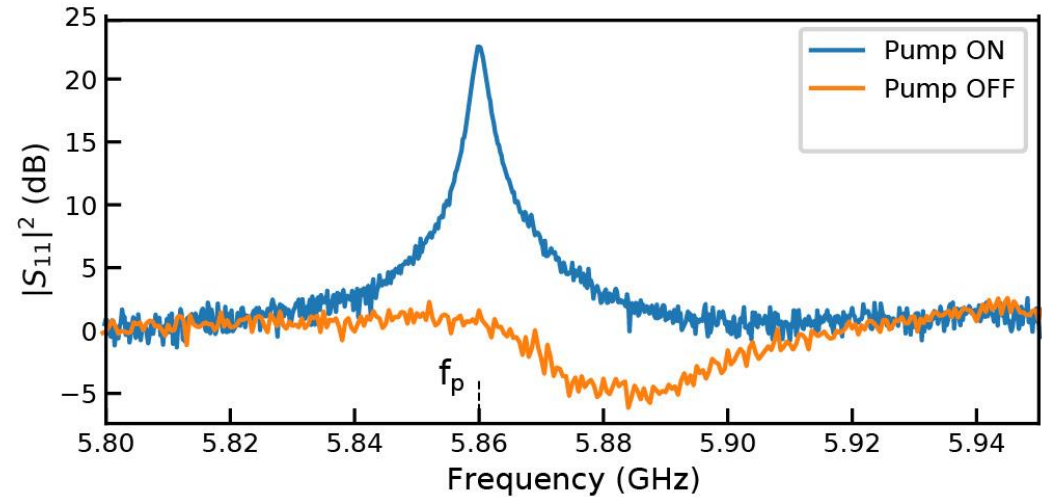
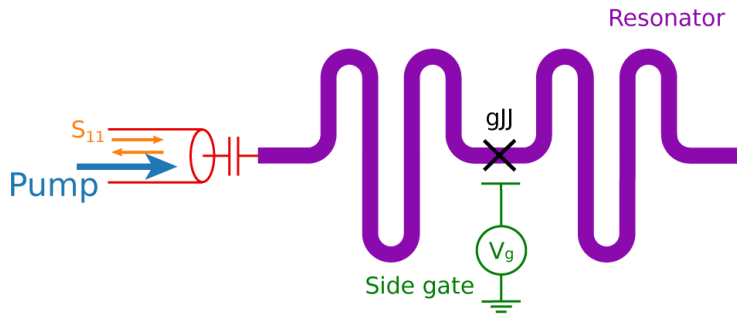


$$L_J = \frac{\Phi_0}{2\pi I_c(V_g)}$$

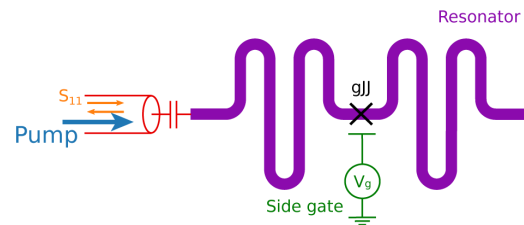
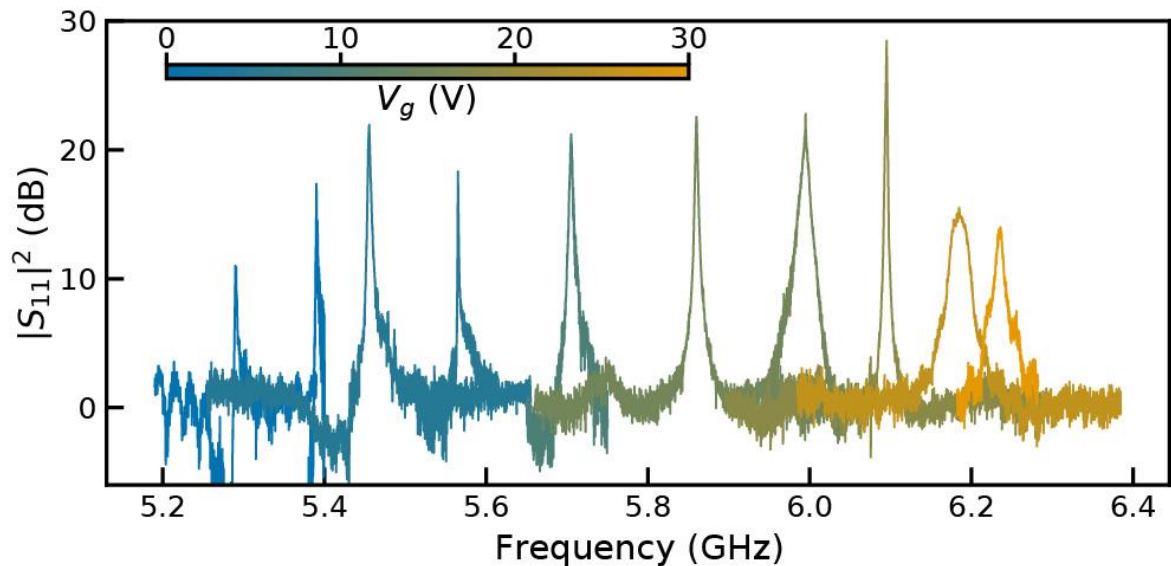
$$\omega_0 = \frac{1}{\sqrt{(L + L_J(V_g))C}}$$



Gate-tunable resonance frequency (>1GHz)



Gain >20 dB thanks to the nonlinearity of the graphene junction

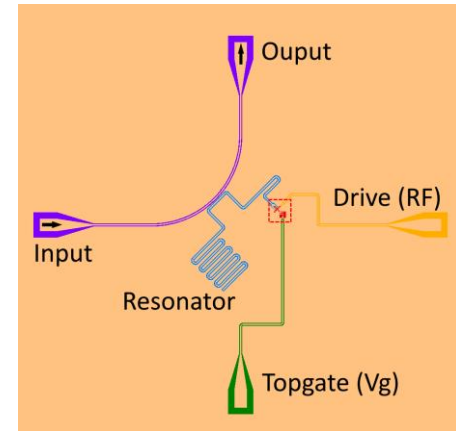
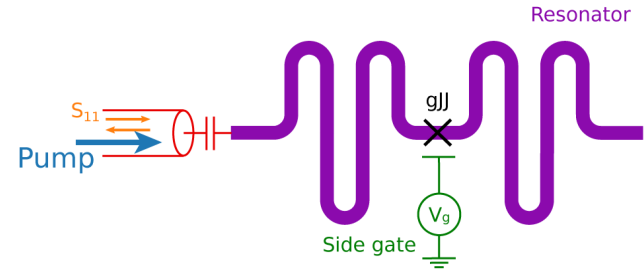


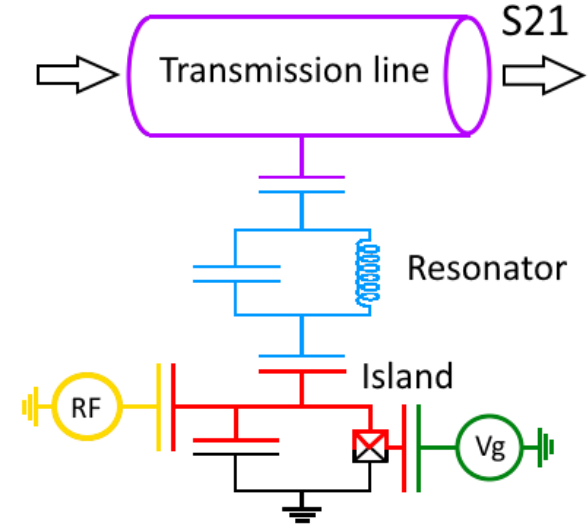
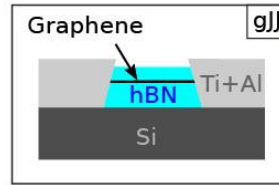
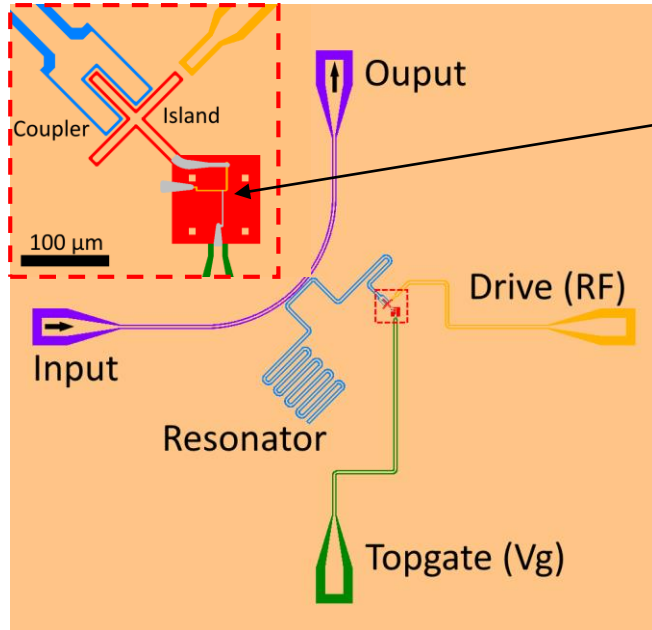
G. Butseraen et al Nature Nano. 2022
 (see also J. Sarkar et al Nature Nano. 2022)

First gate-tunable Josephson parametric amplifier

Gate-tunable Josephson
parametric amplifier

Qubit with a gate-tunable
anharmonicity





Gate-tunable superconducting Qubit

Tunnel junction ($\tau \rightarrow 0$)



$$H = 4E_c n^2 - E_J(B) \cos \phi$$

E_c : charging energy (geometry) ~ 400 MHz

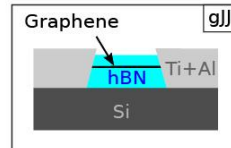
n : charge

E_J : Josephson energy

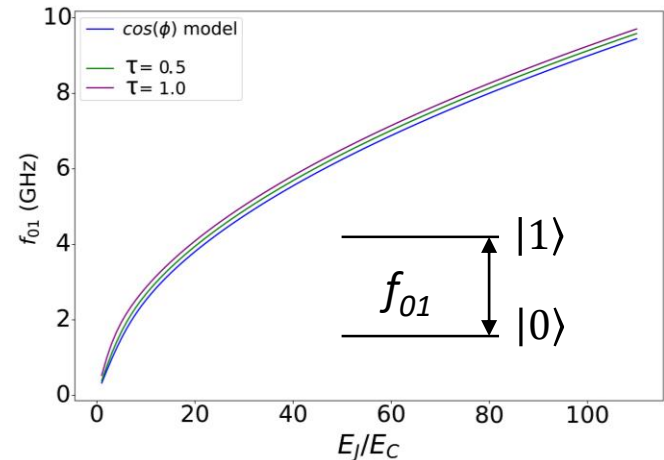
ϕ : phase

τ : transmission

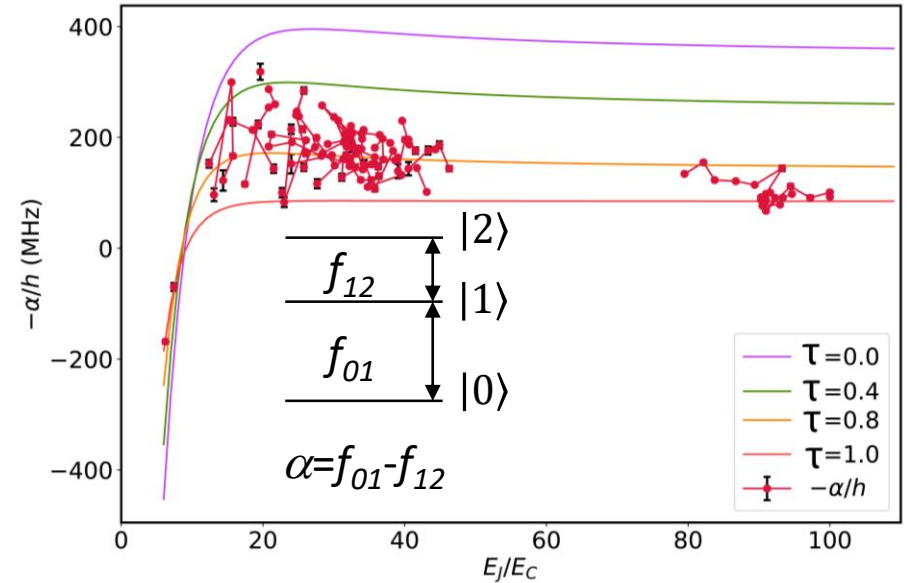
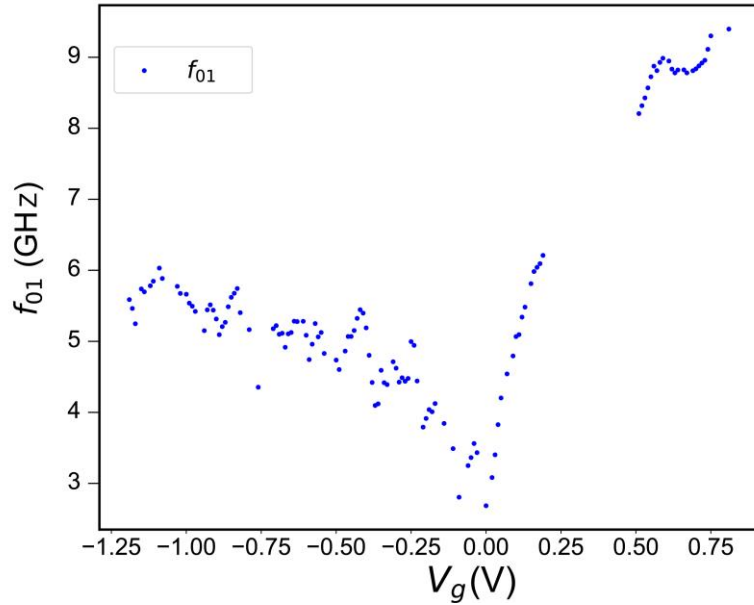
Graphene junction



$$H = 4E_c n^2 - \frac{4E_J(V_g)}{\tau} \sqrt{1 - \tau \sin^2 \frac{\phi}{2}}$$



Gate-tunable Qubit spectrum



Gate-tunable frequency (2-10 GHz) and anharmonicity (α)

- High kinetic inductance in TiN thin films

K. Rafsanjani Amin, N. Roch (Univ. Grenoble Alpes, CNRS, Grenoble INP, Institut Néel)
C. Ladner, G. Jourdan, S. Hentz (Univ. Grenoble Alpes, CEA, LETI)



- Graphene superconducting quantum circuits

G. Butseraen, A. Ranadive, N. Aparicio, K. Rafsanjani Amin, A. Juyal, M. Esposito, E. De Seze, N. Zhurbina, S. Messelot, J. Coraux, N. Roch (Univ. Grenoble Alpes, CNRS, Grenoble INP, Institut Néel)
K. Watanabe, T. Taniguchi (NIMS, Tsukuba, Japan)
F. Lefloch (Univ. Grenoble Alpes, CEA, Grenoble INP, IRIG)

